# Crystal Chemistry of Boron-Bearing Mullite-Type Compounds

Inauguraldissertation der Philosophisch-naturwissenschaftlichen Fakultät der Universität Bern

vorgelegt von

Martin Fisch von Egnach TG

Leiter der Arbeit: Prof. Dr. T. Armbruster Institut für Geologie

Von der Philosophisch-naturwissenschaftlichen Fakultät angenommen.

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Bern, 13. April 2011

Der Dekan Prof. Dr. S. Decurtins

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## 1 Summary

Compounds with mullite-type crystal structures are characterized by unique thermal properties making them important raw materials for industrial applications. Within a constrained stability field in the system  $Al_2O_3-B_2O_3-SiO_2$ , mullite-type boron compounds exist. Boronbearing alumosilicates with mullite-type structures are rare in nature and restricted to specific formation conditions, whereas mullite-type aluminoborates have not been observed in nature. With regard on industrial applications,  $Al_{18}B_4O_{33}$  is the most cited boron-mullite.

Difficulties in precise quantitative boron analyses may be the reason that this compound is reported with two slightly different compositions:  $Al_{18}B_4O_{33}$  (= $Al_{4.91}B_{1.09}O_9$ ,  $9Al_2O_3$ :2 $B_2O_3$ ) and  $Al_5BO_9$  (5 $Al_2O_3$ : $B_2O_3$ ). The crystal structure of this compound is in agreement with the  $Al_5BO_9$  stoichiometry and the discrepancy to the boron-richer composition was explained by a small amount of boron substituting aluminum.

In this study, single-crystal and powder samples were synthesized by flux, solid-state and sol-gel methods from various Al- and B-containing starting materials representing both stoichiometries. Position and coordination of boron in the crystal structures of the samples was precisely determined by single-crystal and powder X-ray diffraction (XRD), as well as by <sup>11</sup>B and <sup>27</sup>Al solid-state MAS nuclear magnetic resonance (NMR), Fourier-transformed infrared (FTIR) and Raman spectroscopy. Results indicated that none of the products crystallized as  $Al_{18}B_4O_{33}$  but rather as  $Al_{5-x}B_{1+x}O_9$  with  $x \approx 0.03(2)$ , which is close to  $Al_5BO_9$ . This result was subsequently confirmed from quantitative  $B_2O_3$  analyses by laser-ablation inductively coupled plasma mass spectrometry.

Compounds with boron in the crystal structure are known for low thermal expansion due to rigid behavior of the boron polyhedra. From in-situ measured high-temperature X-ray diffraction data, thermal expansion parameters  $\alpha$  were obtained for synthetic Al<sub>5</sub>BO<sub>9</sub>, Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> and natural grandidierite ((Fe,Mg)Al<sub>3</sub>BSiO<sub>9</sub>). Despite its different structure type, F-rich jeremejevite (Al<sub>6</sub>B<sub>5</sub>O<sub>15</sub>(F,OH)<sub>3</sub>) was included because it contains the same cations as the aluminoborates.

Results showed that anisotropic thermal expansion of the compounds is governed by changes of M-O-M angles and anomalous anisotropic behavior of Al- and Mg-polyhedra, whereas angular changes and expansion of boron polyhedra are not significant. Expansion of  $Al_5BO_9$  is lowest parallel to the chains of Al-octahedra.

During high-temperature data collection of jeremejevite  $(Al_6B_5O_{15}F_3)$ , the sample decomposed above ca. 800°C, yielding a phase identified as mullite-type  $Al_4B_2O_9$ . At ca. 1050°C,  $Al_4B_2O_9$  transforms to  $Al_5BO_9$ .

Elastic behavior and pressure-induced structural evolution of  $Al_5BO_9$  was analyzed by insitu high-pressure single-crystal X-ray diffraction up to 7.4 GPa at hydrostatic conditions. With increasing pressure, the unit cell volume decreases due to changes of M-O-M angles. The lowest compressibility has been found along the chains of octahedra.

Solid solutions in the "B-mullite" stability field have been suspected, however, only few systematic data are available. Thus, sol-gel derived samples between compositions of aluminoborates  $Al_5BO_9$  and  $Al_4B_2O_9$  and the alumosilicate sillimanite ( $Al_2SiO_5$ ) were synthesized at 950°C and 1250°. Decreasing lattice parameters of mullite-like phases (space group *Pbam*) may indicate boron incorporation in the 950°C samples. Samples with less than 20 mol-%  $SiO_2$  and more than 20 mol-%  $B_2O_3$  contain four-fold coordinated boron, as identified from FTIR spectra. A mullite-type phase in space group  $Cmc2_1$  was detected in the 1250°samples. Refining this phase with variable amounts of  $Al_5BO_9$ - and sillimanite-modules may indicate a trend to more Si towards sillimanite composition. However, excess phases and sluggish reaction rates are an unambiguous indicator for non-equilibrium conditions. These preliminary results may thus be regarded as lead for further studies.

## 2 Introduction

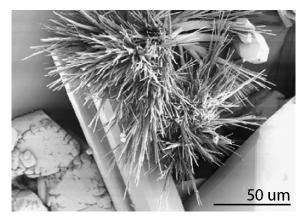
## 2.1 Mullite

Mullite as mineral occurs rarely in nature where it forms as prismatic needles from contact metamorphosed (high temperature, low pressure) clay-sandstones (Tröger, 1971) or in argillaceous inclusions in Tertiary eruptive rocks (Figure 2.1). Mullite is often associated with corundum, sillimanite, kyanite, magnetite, spinel, pseudobrookite, sanidine and cordierite (Anthony et al., 2003).

Mullite as crystalline compound was used in synthetic materials long before it was found as a mineral in nature, although the material was not known and recognized as mullite.

Around 620 AD, the Chinese were among the first who used fired clay for the production of porcelanic wares (Schneider 2005a and refs. therein). In their advanced furnaces they were able to reach temperatures as high as  $1250^{\circ}$ C to  $1400^{\circ}$ C (Kerr et al. 2004), at which mullite is formed from the raw materials.

In the late Middle Age in Europe, ceramic crucibles known as "Hessian porcelain wares" (Figure 2.2) were prominent for their superior quality. Their outstanding properties and quality were ascribed to the secret production process known as the "mystery of the Hessian wares" (Marcos et al. 2006). Hessian wares were made from a mixture of kaolinitic clay and quartz sand, which was subsequently heated at – for that time – unusually high temperatures. The processing of the raw materials at high temperatures led to the formation of mullite which acted as reinforcer of the ceramic materials (Figure 2.9).



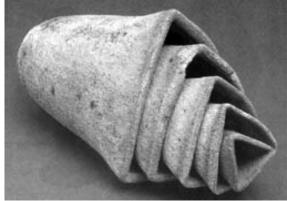


Figure 2.1: SEM photograph of mullite fibers Figure 2.2: Ceramic crucibles known as "Hessian found in volcanic rocks (Eifel Mountain, Germany, wares" (from Marcos et al. 2006). from Schneider et al. 2008).

From 1850 to the early  $20^{\text{th}}$  century, scientists observed agicular, fibrous crystals occurring in synthetic compounds produced at high temperature in the system  $\text{Al}_2\text{O}_3\text{-SiO}_2$ . Based on their habit, the crystals were identified as sillimanite (Fischer and Schneider 2005, Schneider 2005a and refs. therein) and its polymorph andalusite (Deville and Caron 1865).

Shepherd et al. (1909) concluded that in the system "Alumina-Silica" ( $Al_2O_3$ -SiO\_2), crystals identified as "sillimanite" together with corundum and cristobailte were the only phases stable at high temperature, and that the  $Al_2SiO_5$  polymorphs and alusite and kyanite transformed to sillimanite when heated above 1300°C.

Bowen et al. (1924a) recognized that this synthetic compound was in fact not sillimanite  $(Al_2O_3:SiO_2)$  but a crystalline material of which they confirmed the chemical composition of  $3Al_2O_3:2SiO_2$ . While investigating rock samples originating from the Scottish Hebridean Island of Mull, Bowen et al. (1924a) found the same material as naturally occurring mineral: "Found in fused argillaceous inclusions in Tertiary eruptive rocks of the Western Isles of Mull, Scotland. Occurs as abundant prisms associated with rare plates of corundum. It is also abundantly developed in many artificial melts and in porcelains, being the material usually called sillimanite. Its similarity to sillimanite is very great." Based on the on the locality at which it was first found, Bowen and Greig (1924b) and Bowen et al. (1924c) proposed the name "mullite" for this alumosilicate mineral.

When stating: "The discovery of mullite is a striking example of the application of physicalchemical methods to mineralogy. The mineral was first noted in artificial preparations and later sought for in argillaceous rocks that had been subjected to a high heat," Bowen et al. (1924a) could not be aware that within the same century, the material first known from synthetic products and only later found in nature would again become a thoroughly investigated and very prominent synthetic crystalline compound.

Nowadays, due to its rare occurrence in nature and the large volume of mullite used in industry, synthetic mullite plays an important role. Mullite can be produced by heat treatment from kaolinite (Al<sub>2</sub>Si<sub>2</sub>O<sub>5</sub>(OH)<sub>4</sub>) and related materials like aluminous sheet silicates (e.g. muscovite KAl<sub>3</sub>Si<sub>3</sub>O<sub>10</sub>(OH,F)<sub>2</sub>), from the polymorphs sillimanite, kyanite and andalusite (Al<sub>2</sub>SiO<sub>5</sub>) and from staurolite ((Fe<sup>2+</sup>,Mg)<sub>2</sub>Al<sub>9</sub>(Si,Al)<sub>4</sub>O<sub>20</sub>(O,OH)<sub>4</sub>) and topaz (Al<sub>2</sub>SiO<sub>4</sub>(F,OH)<sub>2</sub>) at temperatures between 1000°C and ca. 1700°C. In all cases, the formed mullite is accompanied by excess SiO<sub>2</sub>, unless an aluminum source is added (Komarneni et al. 2005). Using the mentioned raw materials is an easy large-scale production route for mullite ceramics, however, most mullites produced from those raw materials are not chemically pure. Synthesis of chemically pure mullite is possible by reaction sintering (solid-state reaction mullites or sinter mullites) of materials like (purified) quartz, cristobalite and fused silica with corundum ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub>),  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and aluminum hydroxides like gibbstie (Al(OH)<sub>3</sub>), diaspore and boehmite (both AlO(OH)). "Sinter mullites" often have compositions close to 3Al<sub>2</sub>O<sub>3</sub>:2SiO<sub>2</sub> (Komarneni et al. 2005). By controlled cooling of fused  $Al_2O_3$  and  $SiO_2$  (e.g. quartz sand) in electric arc furnaces, "fused mullite" is produced (Schneider 2004). Large mullite single crystals have been successfully produced with float zone and the Czochralski method. Both fused mullite and mullite single-crystals are known with compositions close to  $2Al_2O_3:SiO_2$  (Komarneni et al. 2005). So called "chemical mullite" (Schneider et al. 2008) is produced from precursors derived from a wide range of different sol-gel routes, which allow the production of chemically very pure mullites with  $Al_2O_3$  contents up to 90 wt.-% (Schneider 2004). Depending on the production process, controlled incorporation of other cations than Al and Si into mullite and factors like growth rate, speed and crystal size and shape can be controlled for various specialized applications (Schneider 2004, Komarneni et al. 2005, Schneider et al. 2008).

## 2.1.1 Mullite Crystal Structure

Sensu stricto, "mullite" is the name of the natural mineral mullite with a composition of  $3Al_2O_3$ :2SiO<sub>2</sub> (Bowen et al. 1924a, Bowen and Greig 1924b, Bowen et al. 1924c). It has been shown by Fischer et al. (1994; 1996) and Fischer and Schneider (2005) that mullite crystallizes as solid solution with variable  $Al_2O_3$ :SiO<sub>2</sub> ratio, which can be explained with the formula of  $Al_2(Al_{2+2x}Si_{2-2x})O_{10-x}$  within a compositional range between  $0.18 \le x \le 0.88$ . For x = 0, the composition is equal to sillimanite, for x = 1 it equals  $Al_2O_3$ . A linear relationship between the crystallographic *a*-axis and the molar  $Al_2O_3$  content is known (Fischer et al. 1994; 1996).

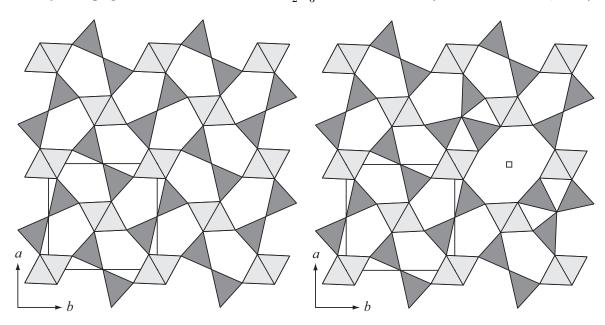


Figure 2.3: Idealized "mullite" structure visualized as sillimanite structure projected along the *c*-axis without any vacancies. (Al,Si) ordering in sillimanite is not shown. The chains of  $AlO_6$  octahedra (light gray) expand along [001]. The bowtie of (Al,Si)-tetrahedra is shown in dark gray.

Figure 2.4: Idealized "mullite" structure visualized as oxygen deficient sillimanite structure projected along [001]. A local vacancy is shown ( $\Box$ ) where the (Al,Si)<sub>2</sub>O<sub>7</sub> bowtie is broken and tetrahedra flip to nearby oxygen sites.

The average structure of mullite in space group *Pbam* with lattice parameters of about  $a \approx 7.57$  Å,  $b \approx 7.68$  Å,  $c \approx 2.88$  Å,  $V \approx 167.44$  Å<sup>3</sup> (Angel and Prewitt 1986) can be explained by the ordered structure of sillimanite: Chains of edge-connected AlO<sub>6</sub> octahedra expand along the *c*-axis (Figure 2.3). Compared to sillimanite, the *c*-axis is halved. At the height of  $z = \frac{1}{2}$  (which corresponds to  $z = \frac{1}{4}$  and  $\frac{3}{4}$  in sillimanite due to the doubled *c*-axis) adjacent octahedral chains are interlinked by (Al,Si)-tetrahedra. Two corner linked (Al,Si)-tetrahedra form bowties, which expand as a double chain along [001] (Figure 2.3). In contrast to sillimanite, Al and Si distribution is not strictly ordered, however, local ordering has been observed (Angel and Prewitt 1986, Angel et al. 1991, Brunauer et al. 2001a, Fischer and Schneider 2005, Rahman and Freimann 2005). Mullite has a higher Al content than sillimanite, hence, some Si<sup>4+</sup> is replaced by Al<sup>3+</sup>. As this exchange is not charge neutral, it can be formalized as: Si<sup>4+</sup>  $\rightarrow$  Al<sup>3+</sup>  $+ \frac{1}{2}$   $\Box_{oxygen}^{2-}$ , or  $2Si^{4+} \rightarrow 2Al^{3+} + \Box_{oxygen}^{2-}$ .

Angel and Prewitt (1986) and refs. therein showed that the oxygen vacancy is localized at the oxygen site interconnecting the (Al,Si)<sub>2</sub>O<sub>7</sub> bowtie. As a consequence of the missing ligand, the central positions of the two tetrahedra building the former bowtie (which are now only three-fold coordinated) flip to nearby oxygen sites, leaving a local void in the structure (Figure 2.4). As shown before, the number of oxygen vacancies depend on the mullite composition (Figure 2.5) and at a certain amount of  $Al_2O_3$ , charge balance in mullite consisting only of  $Al_2O_3$  and SiO<sub>2</sub> of must be compensated by oxygen vacancies at other sites.

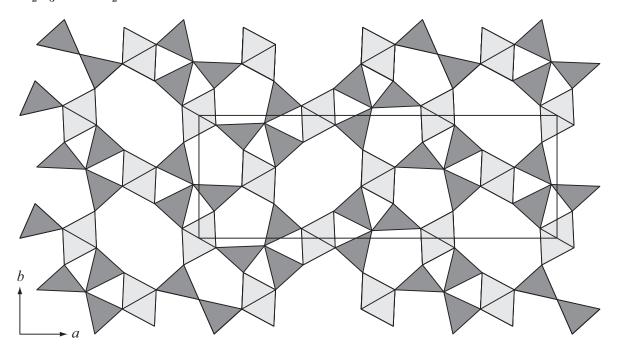
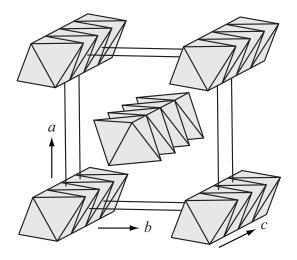


Figure 2.5: Hypothetically ordered structure of a  $4Al_2O_3$ :1SiO<sub>2</sub>-mullite. The supercell with a three-fold *a*-axis is indicated with the black frame (modified after Fischer and Schneider 2005).

## 2.1.2 Definition of Mullite-Type Structures

As more and more synthetic compounds with mullite-like properties were found, a classification was needed. To provide a consistent classification of the relations of different mullite-like structures, the term "mullite-type" was introduced by Fischer and Schneider (2005). The classification is based on mullite characteristic features (see below) and the relation between different mullite-types is based on their relation to the aristotype space group P4/mbm. Hence, "mullite-type" does not classify compounds directly on their exact structure and chemistry but rather comprises an expandable family of structures with common (i.e. mullite-like) structural features. Most noticeable are the parallel chains of edge-connected octahedra, which are interlinked with ordered or disordered arrangements polyhedra of different coordination. Fischer and Schneider (2005) defined the following four characteristics which structures must comply in order to belong to the mullite-type family:

- A mullite-type structure crystallizes in a space group, which can be traced back to the tetragonal space group P4/mbm by group-subgroup relations. The group-subgroup relations are used to classify mullite-type compounds in different groups and are extensively described in Fischer and Schneider (2005);
- 2. Chains of edge-connected  $MO_6$  octahedra are "linear representing single Einer-chains in their highest topological symmetry in space group P4/mbm". This means that  $MO_6$ chains expand linearly (i.e. as stringed together octahedra) in one direction (Figure 2.6);
- 3. In an edge-connected chain of  $MO_6$  octahedra, two oxygen apexes do not contribute to the chain linkages. An axis through these apex atoms must point towards an edge of an octahedron in an adjacent chain. This edge is parallel to the expansion direction of the chain (Figure 2.7). In a hypothetical structure consisting of the  $MO_6$ -chains only in space group P4/mbm, this angle  $\omega$  is 90°. In mullite-type crystal structures, it can vary from 30° to 90°. The lower limit of 30° is based on experimental data from investigations on mullite-type crystal structures;
- 4. Viewed along the chains of  $MO_6$  octahedra, the structure should resemble the orthogonal metric of space group P4/mbm as close as possible. The angle  $\gamma'$  (Figure 2.7) is the angle between basis vectors of the unit cell projected onto the plane perpendicular to the expansion direction of the  $AlO_6$  chains. The tolerance of  $\gamma'$  is  $90^\circ \pm 5^\circ$ . A mullite-type structure can also crystallize in an e.g. monoclinic crystal system as long as the angle of the basis vectors in a plane perpendicular to the direction of the chains is  $90^\circ \pm 5^\circ$ . The deviation of 5° from 90° is based on experimental data from known structures and may be increased when new structures with mullite-like physical properties are discovered.



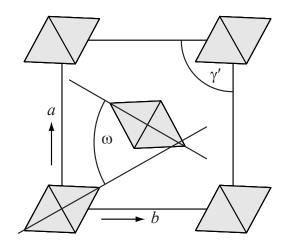


Figure 2.6: Chains of edge-connected  $AlO_6$  octahedra expanding unfolded and linear in one direction. In the simplified structure shown here (data from Angel and Prewitt 1986), the translational distance between octahedra is equal to the length of the *c*-axis.

Figure 2.7: Simplified crystal structure of mullite (data from Angel and Prewitt 1986), projected along the *c*-axis. The angle  $\omega$  between adjacent chains of octahedra and the angle  $\gamma$ ' between the basis vectors projected onto the plane perpendicular to the chain direction are displayed (modified after Fischer and Schneider 2008a).

## 2.1.3 Applications of Mullite

Mullite is widely used in ceramics industry due to outstanding properties like its very high melting point of about 1890°C, its low thermal expansion coefficient, which is the reason for good thermal shock and spalling resistance, its good insulating properties due to the low thermal conductivity, the high creep resistance, its high shear modulus, a high electrical resistance and a low gas permeability (Komarneni et al. 2005). Schneider et al. (2008) divided mullite application in three groups:

**Monolithic mullite ceramics.** This group includes traditional ceramics (i.e. tableware and porcelain) and industrial ceramic products like refractory materials and engineering ceramics (Schneider et al. 2008). Engineering materials comprise mechanical guides, brake linings, filters, supporting substrates for catalysts, corrosion resistant crucibles, casings for thermocouples and heat exchangers, substrates for electronic devices with high-temperature stability and high electrical resistance but also mechanical parts like conveyor belts (Figure 2.8) in high-temperature furnaces and reinforced parts in engines and turbines. Refractory materials include chamotte stones and mullite based bricks in various types of industrial furnaces. Transparent mullite glass-ceramics have been successfully applied as high-temperature stable windows (Okada and Schneider 2005, Schneider 2004).

**Mullite coatings.** They are used to protect surfaces of materials against corrosion and oxidation due to environmental factors like gases, acids and bases, heat etc. The thin (few µm) coatings can be applied onto substrates by chemical and physical vapor-deposition and plasma- and flame spray methods (Basu and Sarin 2005).

**Mullite matrix composites.** Composites of a mullite matrix, which is reinforced with fibers or whiskers of mullite or of different composition (e.g. SiC, Al<sub>2</sub>O<sub>3</sub>) in order to improve durability and mechanical strength, especially at high temperature (Figure 2.9) are called "mullite matrix composites". The mechanism may be compared to the use of asbestos fibers in cement materials. Mullite matrix composites are a field of ongoing research, however, potential applications may include propeller blades of gas turbines and heat shields in aerospace industry (Schneider 2005b) as well as military applications (MIL-HDBK-17-5 2002). Mullite fibers can also be used as reinforcer of glasses and glass ceramics (Travitzky 1998).

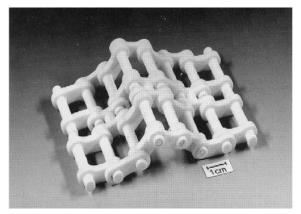
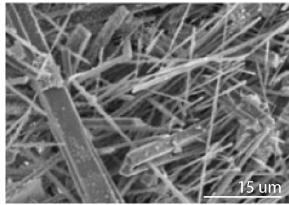


Figure 2.8: Conveyor belt made from sin- Figure 2.9: Mullite fibers in a cross-section from tered mullite. Such parts are used to feed hightemperature furnaces (from Schneider et al. 2008).



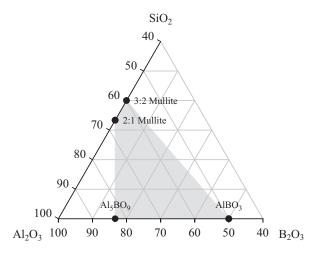
a Hessian crucible. The performance of the Hessian crucibles is ascribed to the fibers, which act as reinforcer in the ceramic matrix (from Marcos et al. 2006).

## 2.2 Boron-Mullites

From analyses of crystalline by-products obtained while investigating glasses in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$ , Dietzel and Scholze (1955) suggested a solid solution between 3:2 mullite and a silicon-free aluminoborate compound identified as  $Al_{18}B_4O_{33}$  ( $9Al_2O_3:2B_2O_3$ ) (Baumann and Moore 1942). This assumption was based on a hypothetical, although not charge neutral exchange of  $B^{3+} \rightarrow Si^{4+}$  and due to the similar physical properties of mullite and Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> (Scholze 1956, Benner and Baumann 1938, Baumann and Moore 1942).

A phase which showed the same X-ray diffraction pattern as Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> was observed by Werding and Schreyer (1984) while investigating compounds in the system MgO-Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub>-H<sub>2</sub>O. They conclude: "According to Scholze this phase exhibits solid solution with mullite, consequently, it has been referred to as "B-mullite" here." Based on their own experiments in the system Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub>-H<sub>2</sub>O, from results of Mazza et al. (1992) who produced crystalline compounds with mullite-like X-ray diffraction lines and variable boron content, and from small amounts of boron found in sillimanite (Grew and Hinthorne 1983), Werding and Schreyer (1996) introduced the term "boron-mullite".

The compositional field of the boron-mullites is constraint by 3:2 mullite  $(3Al_2O_3:2SiO_2)$ and 2:1 mullite  $(2Al_2O_3:1SiO_2)$  and the silicon free aluminoborates  $Al_5BO_9$   $(5Al_2O_3:1B_2O_3)$ with mullite-type structure and  $AlBO_3$   $(Al_2O_3:B_2O_3)$  with calcite-type structure (Capponi et al. 1972, Bither 1973, Vegas et al. 1977), as shown in Figure 2.10.



**Figure 2.10:** Plot showing the boron-mullite stability filed in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$  as defined by Werding and Schreyer (1996). Units are mol-%.

Not all compounds within the boron-mullite field necessarily crystallize with mullite-type structures (e.g.  $AlBO_3$  with calcite-type structure). On the other hand, e.g. grandidierite  $(Mg,Fe)Al_3BSiO_9$  (Stephenson and Moore 1968), is not included in the boron-mullite field despite its mullite-type crystal structure. As also a natural aluminum borosilicate named "boromullite" exists (Buick et al. 2008), a precise interpretation of the term "boron-mullite" is difficult.

Consequently, Fischer and Schneider (2008a) introduced the term "mullite-type boron compound", which includes any boron-containing compound with mullite-type structure. "Boronmullite" and "B-mullite" are reserved for chemically pure aluminoborates and aluminum borosilicates with mullite-type structures. In this context, a mullite-type boron compound must fulfill the same structural requirements as those set up for mullite-type structures (Section 2.1.2) and the mullite classification scheme of Fischer and Schneider (2005) is also applied to mullite-type boron compounds.

#### 2.2.1 Phases in the Boron Mullite Compositional Field

Phases with a chemical composition within or close to the boron-mullite compositional field (Werding and Schreyer 1996) are shown in Figure 2.11. Due to the few known aluminoborate

structures, the high-pressure phase  $Al_3BO_6$  (Capponi et al. 1972) with norbergite structure (White et al. 1965), hydrothermally synthesized  $AlBO_3$  with calcite structure (Capponi et al. 1972, Bither 1973, Vegas et al. 1977), the microporous framework structures PKU-5 ( $Al_4B_6O_{15}$ , Jing et al. 2004) and PKU-6 ( $HAl_2B_3O_8$ , Yang et al. 2007) and hexagonal jeremejevite (hydrothermally formed in granitic pegmatites, Foord et al. 1981) are also included.  $Al_3BO_6$  and  $AlBO_3$  have been found as co-products of jeremejevite syntheses (Stachowiak and Schreyer 1998, Reynaud 1977). A detailed description of the structure of jeremejevite and its thermal evolution is given in Section 5.5.4 in Chapter 5.

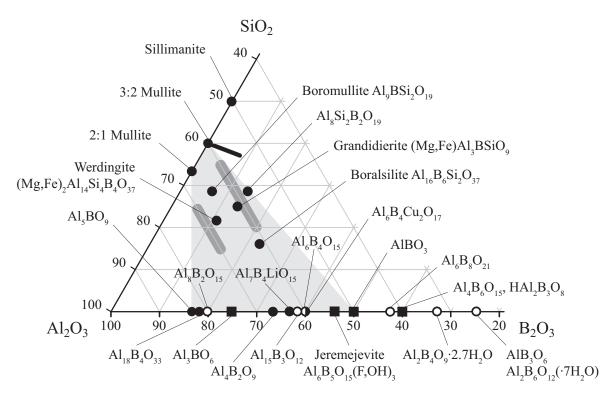


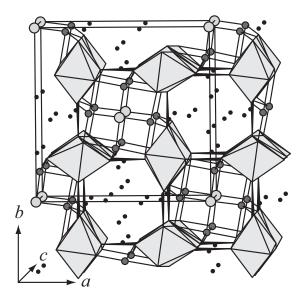
Figure 2.11: Plot showing phases with chemical compositions in or close to the boron-mullite stability field (light-gray) defined by Werding and Schreyer (1996). Except on the  $Al_2O_3$ - $B_2O_3$  line, only phases with mullite-type crystal structures are shown. For werdingite, grandidierite, jeremejevite,  $Al_7B_4LiO_{15}$ ,  $Al_6B_4Cu_2O_{17}$  and  $HAl_2B_3O_8$ , the system has to be extended according to the chemistry of the phases and the compositional projection onto the  $Al_2O_3$ - $B_2O_3$ -SiO<sub>2</sub> system is shown. Phases indicated with filled black circles crystallize with mullite-type structure (Fischer and Schneider 2008a), whereas black squares indicate phases with different crystal structures. The crystal structure of phases in synthesis experiments of jeremejevite or of phases with mullite-type structures. Incorporation of boron into mullites as proposed by Griesser et al. (2008) is indicated with the two dark-gray fields, whereas boron-doped mullites according to Zhang et al. (2010) are indicated with the black line. Units are mol-%.

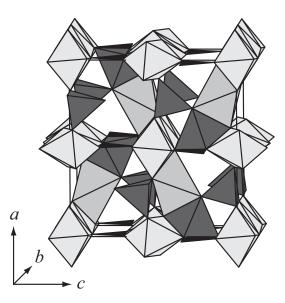
Several aluminoborate structures are still not solved and their composition is only guessed from questionable chemical analyses. During hydrothermal synthesis experiments starting from Al(OH)<sub>3</sub> and H<sub>3</sub>BO<sub>3</sub>, Lehmann and Teske (1973) found Al<sub>6</sub>B<sub>4</sub>O<sub>15</sub>, Al<sub>2</sub>B<sub>4</sub>O<sub>9</sub>·2.7H<sub>2</sub>O and  $Al_2B_6O_{12}\cdot 7H_2O$ . Reynaud (1977) reports  $Al_8B_2O_{15}$ , synthesized from  $AlF_3$ , borax and  $H_3BO_3$  at 1400°C, and  $Al_6B_8O_{21}$ , found during a hydrothermal synthesis experiment as a cophase of jeremejevite. From some of the phases, powder X-ray patterns exist, however, none of the phases has been indexed, nor are any structure solutions available so far.  $Al_6B_8O_{21}$  was subsequently confirmed by Stachowiak and Schreyer (1998) in jeremejevite synthesis experiments. The existence of  $AlB_3O_6$  is rather unclear, however, it is mentioned as  $AlB_3O_6$  in an US patent on glasses (Rogers 1950), by Stachowiak and Schreyer (1998) and as  $Al_2B_6O_{12}$  and  $Al_2B_6O_{12}\cdot 7H_2O$  by Lehmann and Teske (1973). During a series of experiments on the stability of the mineral werdingite, Werding and Schreyer (1992) obtained  $Al_8Si_2B_2O_{19}$  (Niven et al. 1991). This phase could be indexed with orthorhombic lattice parameters similar to those of  $Al_5BO_9$  (Sokolova et al. 1978). No further evidence than a reference in Stachowiak and Schreyer (1998) provides any proof for the existence of  $Al_5B_3O_{12}$ .

In this section, phases crystallizing with mullite-type crystal structures are listed according to the classification concept of Fischer and Schneider (2005). The concept has been extended by Fischer and Schneider (2008a) and unit cell settings of the structures of boronbearing mullite-type compounds are standardized to the one of mullite for better comparison. Nonetheless, structures in this chapter are reported in their standard space group setting.

 $Al_6B_4Cu_2O_{17}$  group. This group includes the borates  $Al_6B_4Cu_2O_{17}$  (Kaduk et al. 1999) and  $Al_7B_4LiO_{17}$  (Åhman et al. 1997), which both crystallize in tetragonal space group I4/m ( $a = b \approx 10.5$  Å,  $c \approx 5.6$  Å). The structure of  $Al_6B_4Cu_2O_{17}$  consists of chains  $AlO_6$  octahedra expanding parallel to the *c*-axis. Neighboring chains are interconnected by  $BO_3$  triangles and clusters of (Al,Cu)O\_5 bipyramids, which are connected at an oxygen atom that is statistically distributed around the four-fold axis. Åhman et al. (1997) already noticed analogies of  $Al_7B_4LiO_{17}$  to mullite, sillimanite and andalusite due to the  $AlO_6$ -chains expanding along the *c*-axis. The crystal structure of  $Al_7B_4LiO_{17}$  is similar to the one of  $Al_6B_4Cu_2O_{17}$ , however, empty channels in  $Al_6B_4Cu_2O_{17}$  are filled with statistically distributed Li (Figure 2.12). Both structures have recently been reinvestigated and confirmed (Plachinda and Belokoneva 2008).

**Grandidierite group.** Grandidierite ((Mg,Fe)Al<sub>3</sub>BSiO<sub>9</sub>) was first discovered in Madagascar and is named after the French naturalist and explorer Alfred Grandidier. The type locality is Andrahomana on the south coast of Madagascar where it has been found as anhedral elongated crystals up to ca. 8 cm length. It is a bluish-greenish to blackish prismatic mineral that is known from 40 localities worldwide (a comprehensive list of localities is given in Grew 1996). It occurs as rare accessory mineral in aluminous boron-rich rocks of continental or oceanic crust where it is formed during metamorphism at pressures lower than 10 kbar (e.g. pegmatites, aplites, gneisses, xenoliths, pelitic hornfels and calc-silicatic rocks). Its temperature stability at atmospheric pressure is between 550°C and ca. 1200°C and it can thus also exist in sanidinite facies rocks, e.g. volcanic xenoliths (Schreyer and Werding 1997). Typical mineral assemblages are quartz, K-feldspar, plagioclase, biotite, garnet, spinel, corundum,





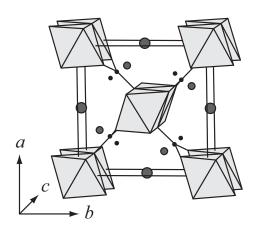
group I4/m. Al atoms are gray, AlO<sub>6</sub> octahe-(light gray sphere). Li atoms (black), are statisti- polyhedra are light gray. cally distributed at the available positions. In the  $Al_6B_4Cu_2O_{17}$  structure, the gray sphere is occupied by (Al,Cu) and the channel with Li atoms is empty.

Figure 2.12: Al<sub>7</sub>B<sub>4</sub>LiO<sub>17</sub> structure in space Figure 2.13: The crystal structure of grandidierite in space group Pnma. Al-octahedra are dra light gray and  $BO_3$  units black. Only one light gray,  $BO_3$  units are black.  $SiO_4$  tetrahedra oxygen atom around the four fold axis is shown and  $AlO_5$  polyhedra are dark gray, (Fe,Mg)O<sub>5</sub>

sillimanite, and alusite, cordierite, tourmaline, kornerupine, sapphirine, serendibite and sinhalite (Anthony et al. 2003). In pelitic rocks, its occurrence seems to be restricted to rocks, in which muscovite and other phases typical for aluminuous amphibolite-facies rocks are missing (Grew 1996).

Grandidierite (Mg,Fe)Al<sub>3</sub>BSiO<sub>9</sub>, is the Mg-rich member of the grandidierite-ominelite series (Dzikowski et al. 2007, Hiroi et al. 2001) and is of particular interest because it contains both Al and (Mg,Fe) in five-fold coordination (MacKenzie and Meinhold 1997, Dzikowski et al. 2007, Farges 2001). The grandidierite crystal structure (space group Pnma,  $a \approx 10.96$  Å,  $b \approx 5.76$  Å,  $c \approx 10.33$  Å) contains two different chains of edge-connected AlO<sub>6</sub> octahedra parallel to the *b*-axis. One chain is rather straight whereas the other is folded. Adjacent Al-octahedra are directly connected by corner-linked chains of BO<sub>3</sub>, AlO<sub>5</sub>, SiO<sub>4</sub> and MgO<sub>5</sub> polyhedra (Figure 2.13). Thermal evolution of the structure of grandidierite is discussed in Section 5.5.3 in Chapter 5.

Mullite group. Mazza et al. (1992) investigated boron incorporation into mullites and proposed a solid solution between compounds produced from starting materials with Al<sub>6-x</sub>B<sub>x</sub>O<sub>9</sub>  $(1 \le x \le 3)$  composition. Average crystal structures were solved in space group *Pbam* with pseudotetragonal lattice parameters  $a \approx b \approx 7.62$  Å and  $c \approx 2.8$  Å. With infrared investigations, BO<sub>3</sub> incorporation into Al<sub>6-x</sub>B<sub>x</sub>O<sub>9</sub> with x = 1 was proposed. The structure of Al<sub>6-x</sub>B<sub>x</sub>O<sub>9</sub> with x = 2 is similar to the one with x = 1, however, one Al site is replaced by a BO<sub>4</sub> site. The occurrence of a BO<sub>4</sub> site in this structure was confirmed by <sup>11</sup>B MAS NMR spectroscopy by MacKenzie et al. (2007). When heated to temperatures between 900°C to 1000°C, the compounds yield Al<sub>5</sub>BO<sub>9</sub> (Chapter 4) and Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> (Fischer et al. 2008b). Average structures of both phases according to Mazza et al. (1992) are shown in Figures 2.14 and 2.15.



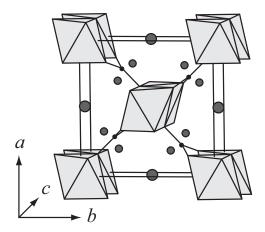


Figure 2.14: Representation of the average structure of  $Al_4B_2O_9$  according to Mazza et al. (1992) in space group *Pbam*. Al-octahedra are light gray, boron atoms are represented as black spheres. The small gray sphere represents Al, the larger gray one represents O.

Figure 2.15: Representation of the average structure of  $Al_5BO_9$  according to Mazza et al. (1992) in space group *Pbam*. Al-octahedra are light gray, boron atoms are represented as black spheres. The small gray sphere represents Al, the larger gray one represents O. In contrast to  $Al_4B_2O_9$ , the additional B site is replaced by Al.

A "mullite-like" phase with  $Al_{8+x}P_{1-x}B_{1+x}O_{16+x/2}$  with  $0 \le x \le 0.6 \pm 0.1$  is described in Mazza et al. (2001). Orthorhombic lattice parameters similar to those of mullite are presented, however, no space group is indicated. As the compounds are described as "mullite-like" containing four-fold coordinated boron, it can be assumed that structure solutions were done in space group *Pbam*, as in Mazza et al. (1992).

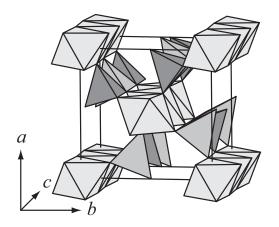
Boron incorporation into the mullite structure was also investigated by Griesser et al. (2008). They proposed a limited solid solution due to decreasing lattice parameters with increasing boron content (Figure 2.11). The final products of Griesser et al. (2008) were investigated by powder X-ray diffraction and FTIR, hence, exact mechanism of boron incorporation into the structure is not completely resolved.

In a recent study on boron-doped mullites (Zhang et al. 2010) with  $B_2O_3$  contents of up to 6.14 wt.-% (Figure 2.11), results of Griesser et al. (2008) were confirmed. Moreover, increasing boron content led to a reduction of crystallization-temperature.

Sillimanite group. Grew and Hinthorne (1983) found B- and Mg-containing sillimanite in boron-rich volcanic and sedimentary rocks subjected to granulite-facies metamorphic condi-

tions. However, high temperatures alone do not guarantee high B contents in sillimanite even when large boron quantities are available. A B  $\rightarrow$  Si substitution is suspected and seems to be coupled with Mg  $\rightarrow$  Al substitution (Grew 1996). For comparison, the structure of B-free sillimanite according to Winter and Ghose (1979) in space group *Pbnm* with  $a \approx 7.49$  Å,  $b \approx 7.68$  Å,  $c \approx 5.78$  Å is shown in Figure 2.16.

Fischer and Schneider (2008a) also included PbMBO<sub>4</sub> with  $M = Al^{3+}$ ,  $Ga^{3+}$ ,  $Fe^{3+}$ ,  $Cr^{3+}$ ,  $Mn^{3+}$  (Park et al. 2003) compounds in this group because of the same symmetry as sillimanite (in space group *Pbam*). In contrast to sillimanite,  $AlO_6$  chains are alternatingly linked with  $BO_3$  triangles or Pb atoms bonded to the chains. At least PbAlBO<sub>4</sub> is not high-temperature stable and the mullite-type structure is lost at 775°C. In Figure 2.17 the structure of PbAlBO<sub>4</sub> is shown in space group *Pnma* with  $a \approx 6.92$  Å,  $b \approx 5.71$  Å,  $c \approx 8.02$  Å.



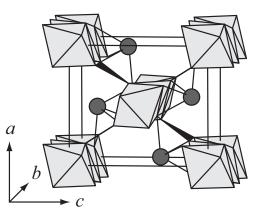
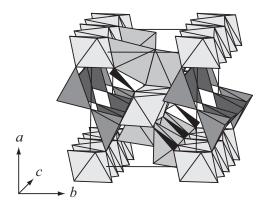


Figure 2.16: Structural representation of sillimanite (in *Pbnm*, Winter and Ghose 1979). Aloctahedra are light gray, chains of alternating  $AlO_4$  and  $SiO_4$  tetrahedra are gray and dark gray, respectively.

Figure 2.17: Crystal structure drawing of PbAlBO<sub>4</sub> according to Park et al. (2003) in space group *Pnma*. Al-octahedra are represented in light gray, whereas BO<sub>3</sub> triangles are drawn in black. Asymmetrically bonded lead atoms (four bonds to octahedra) are plotted as gray spheres.

Werdingite group. Werdingite  $((Mg,Fe)_2Al_{14}Si_4B_4O_{37})$  is a (Mg,Fe)-Al-borosilicate found in granulite-facies metamorphosed metasediments and gneisses (Anthony et al. 2003) in Namaqualand (type locality), South Africa. It is stable at pressures up to ca. 10 kbar and temperatures between ca. 700°C and 1200°C. It occurs together kornerupine, grandidierite, zircon, and rutile (Anthony et al. 2003) in hercynite-sillimanite-rich layers in a supercrustal gneissic sequence (Moore et al. 1990, Grew 1996). The Mg-endmember was successfully synthesized by Werding and Schreyer (1992). Werdingite has an idealized chemical composition of  $(Mg,Fe)_2Al_{14}Si_4B_4O_{37}$  (Moore et al. 1990) and crystallizes in triclinic space group  $P\bar{1}$  with  $a \approx 7.97$  Å,  $b \approx 8.16$  Å,  $c \approx 11.36$  Å,  $\alpha \approx 110.36^\circ$ ,  $\beta \approx 110.88^\circ$  and  $\gamma \approx 84.67^\circ$  (Figure 2.18). Two non equivalent AlO<sub>6</sub> chains are interlinked with either AlO<sub>4</sub> and BO<sub>3</sub> groups (only BO<sub>3</sub> shown in the structural drawing), with (Al,Fe)O<sub>4</sub> and SiO<sub>4</sub> tetrahedra and with AlO<sub>5</sub> and  $(Al,Mg)O_5$  polyhedra (Niven et al. 1991).



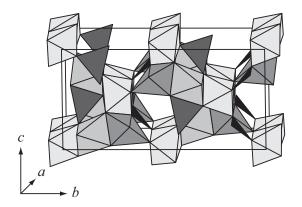


Figure 2.18: The triclinic crystal structure of Figure 2.19: Structural drawing of  $Al_5BO_q$  in werdingite in space group  $P\overline{1}$ . Al-octahedra are plotted light gray,  $BO_3$  units in black. (Al,Fe)O<sub>4</sub> and  $SiO_4$  tetrahedra are dark gray,  $AlO_5$  and  $(Mg,Fe)O_5$  polyhedra are gray.

space group  $Cmc2_1$  (Sokolova et al. 1978). Aloctahadra (light gray) are interconnected with  $\mathrm{BO}_3$  triangles (black),  $\mathrm{AlO}_4$  tetrahedra (dark gray) and two different  $AlO_5$  polyhedra. For better visualization of the mullite-type structure, the unit cell is shifted by 0, 0.1164, 0.5.

 $A_9B_2$  (Al<sub>18</sub> $B_4O_{33}$ ) group. Al<sub>18</sub> $B_4O_{33}$  was first mentioned by Baumann and Moore (1942) who determined that the compound previously cited as  $3 \text{Al}_2 \text{O}_3$ : B<sub>2</sub>O<sub>3</sub> (Mallard 1887, Benner and Baumann 1938) has a composition of  $9 \operatorname{Al}_2 \operatorname{O}_3: 2 \operatorname{B}_2 \operatorname{O}_3$ . Subsequent structural investigations were in agreement with  ${\rm Al}_5{\rm BO}_9$  composition (Sokolova et al. 1978). The boron discrepancy was explained with substitution of 9%  $B^{(IV)} \rightarrow Al^{(IV)}$  (Ihara et al. 1980) and it has been shown that up to 10% of octahedral Al can be replaced by  $Cr^{3+}$  (Garsche et al. 1991). The compound is still frequently cited as  $Al_{18}B_4O_{33}$ . Recently, Fisch et al. (2011) showed that the  $Al_{18}B_4O_{33}$  composition is rather unlikely and that the compound should be cited as  $Al_{5-x}B_{1+x}O_9$  with  $x \approx 0.03$  (Chapter 4). The structure is shown in Figure 2.19 (space group  $Cmc2_1$ ,  $a \approx 5.67$  Å,  $b \approx 15.01$  Å,  $c \approx 7.68$  Å). Isolated edge-sharing AlO<sub>6</sub> octahedra run parallel to the a-axis. Proximate octahedral chains are connected by pairs of edge-sharing  $AlO_5$  polyhedra and  $AlO_4$  linked with  $BO_3$  triangles. One side of an  $AlO_5$  polyhedra pair is connected to octahedra and linked to BO<sub>3</sub> triangles, whereas the other side links to octahedra and  $\mathrm{AlO}_4$  tetrahedra. The structure is described in more detail in Chapter 4 and thermal evolution of the compound is discussed in Section 5.5.1 in Chapter 5.

Boromullite Al<sub>9</sub>BSi<sub>2</sub>O<sub>19</sub> was reported in granulite-facies metapelitic rocks from Mount Stafford, Australia (type locality), where it is formed during anatexis (Grew et al. 2008). At its type locality, it occurs together with tourmaline, werdingite, grandidierite, sillimanite, cordierite, K-feldspar and rarely, biotite, hercynite, ilmenite and quartz (Buick et al. 2008). It crystallizes in the same space group as  $Al_{18}B_4O_{33}$  (*Cmc*2<sub>1</sub>) and has similar lattice parameters. Structurally, boromulite corresponds to a 1:1 polysome composed of Al<sub>5</sub>BO<sub>9</sub> and Al<sub>2</sub>SiO<sub>5</sub>

modules (Figure 2.20). The  $Al_2SiO_5$  module contains all Si and has the topology of sillimanite with a doubled *b*-axis (resulting in a composition of  $Al_4Si_2O_{10}$ ) whereas the  $Al_5BO_9$  module agrees with the structural model of Sokolova et al. (1978).

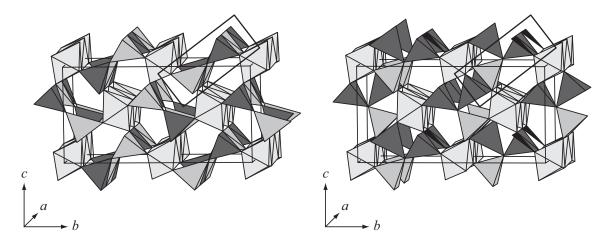
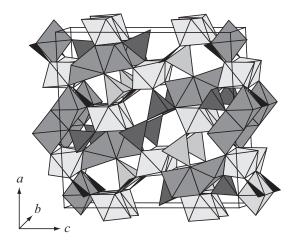


Figure 2.20: Structural representation of the sillimanite (left) and  $Al_5BO_9$  (right) modules of boromullite. Al-octahedra are represented in light gray. The important structural difference between the two modules with respect to the 1:1 polysome boromullite is marked with the black square. For better comparison to the  $Al_5BO_9$  structure, the unit cell is shifted by 0.25443, 0.12121, 0.65357.

**Boralsilite** (Al<sub>16</sub>B<sub>6</sub>Si<sub>2</sub>O<sub>37</sub>) and Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> group. The mineral boralsilite (Al<sub>16</sub>B<sub>6</sub>Si<sub>2</sub>O<sub>37</sub>) was found as the first anhydrous Al-B-silicate relatively widespread in granulite-facies granitic pegmatites in metapelitic rocks at Larsemann Hills, East Antarctica (type locality). It crystallizes as bundles of prisms in assemblages with quartz, K-feldspar, tournaline, plagioclase, werdingite, dumortierite, grandidierite and sillimanite (Grew et al. 1998). The crystal structure has been solved by Peacor et al. (1999) in space group C2/m with  $a \approx 14.77$  Å,  $b \approx$ 5.57 Å,  $c \approx 15.08$  Å,  $\beta \approx 91.96^{\circ}$ . Three different AlO<sub>6</sub> chains prolong parallel to the *b*-axis. They are interconnected by three different polyhedral arrangements. One consists of an Si<sub>2</sub>O<sub>7</sub> bowtie, the second of AlO<sub>5</sub> polyhedra, AlO<sub>4</sub> tetrahedra and BO<sub>3</sub> triangles, and the third is built up from AlO<sub>5</sub> polyhedra and BO<sub>4</sub> tetrahedra (Figure 2.21).

Several ambiguous structure solutions for mullite-type  $Al_4B_2O_9$  exist but Scholze (1956) was the first who indexed it with orthorhombic lattice parameters a = 14.8(2) Å, b = 15.1(2)Å and c = 5.6(1) Å. After recognizing the similarity of the  $Al_4B_2O_9$  structure to the one of boralsilite, Fischer et al. (2008b) provided a promising, although not completely solved structure in space group C2/m ( $a \approx 14.81$  Å,  $b \approx 5.54$  Å,  $c \approx 15.05$  Å,  $\beta \approx 90.91^{\circ}$ ) based on powder XRD, <sup>11</sup>B and <sup>27</sup>Al MAS-NMR data: Three types of edge-connected chains of  $AlO_6$ octahedra extend parallel to the *b*-axis, interlinked by three different polyhedral configurations (Figure 2.22). One of them is a complex arrangement of distorted  $AlO_5$ ,  $AlO_4$  polyhedra and  $BO_3$  triangles. The second one comprises  $AlO_{4+1}$  and  $BO_4$  tetrahedra which are disordered with  $AlO_4$  tetrahedra and  $BO_3$  triangles due to a partially occupied oxygen site. The third type contains  $BO_3$  and/or  $BO_4$  units and replaces the  $Si_2O_7$  bowtie in boralsilite. Its exact configuration has not been resolved yet. Thermal evolution of  $Al_4B_2O_9$  is discussed in Section 5.5.2 in Chapter 5.



**Figure 2.21:** Crystal structure of boralsilite (in C2/m). AlO<sub>6</sub> chains are drawn in light gray. The chains are interconnected with AlO<sub>5</sub> and AlO<sub>4</sub> polyhedra, BO<sub>3</sub> triangles and Si<sub>2</sub>O<sub>7</sub> bowties (dark gray).

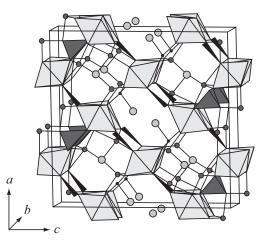


Figure 2.22:  $Al_4B_2O_9$  structure according to Fischer and Schneider (2008a) in space group C2/m.  $AlO_6$  octahedra (light gray) are interconnected with Al-polyhedra (Al is represented with the small sphere) and  $BO_4$  tetrahedra (dark gray). The  $Si_2O_7$  bowtie of boralsilite is replaced by  $BO_3$ units (black) and further two O sites and two B sites of which the exact configuration has not been solved yet (modified after Fischer and Schneider 2008a).

## 2.2.2 Applications of Boron Mullites

With regard to industrial application, the compounds  $Al_{18}B_4O_{33}$  and  $Al_4B_2O_9$  are most cited. Due to mullite-like properties and the low-cost production (ca.  $\frac{1}{20}$  to  $\frac{1}{10}$  of the cost of SiC, Kim et al. 2004 and Song et al. 2007), easy fabrication in large quantities (Wada et al. 1991, Ray 1992, Sokolov and Gasparyan 2004, Zhang et al. 2006, Peng et al. 2006, Elssfah et al. 2007, Xinyong et al. 2007, Wang et al. 2008, Zhou et al. 2009, Chen et al. 2009), high strength (Ray 1992, Xinyong et al. 2007, Gatta et al. 2010) and low thermal expansion and conductivity (Wada et al. 1993; 1994),  $Al_{18}B_4O_{33}$  and  $Al_4B_2O_9$  are used as reinforcer in ceramic and ceramic/metal matrix composites and in glass ceramics (Peng et al. 2006, Wang et al. 2008, Fan et al. 2009, Feng et al. 2009, Lee et al. 2010, Geng et al. 2010).  $Al_{18}B_4O_{33}$  has been proposed as reinforcer in fire insulations in ships due to its low weight, as construction components in nuclear reactors because of its neutron absorbing capabilities and in refractory linings due to the high resistance against borosilicate melts (Garsche et al. 1991 and refs. therein). Recently, nanotubes have been reported (Ma et al. 2002, Li and Chang 2006) and  $Al_{18}B_4O_{33}$  fibers have been coated with boron nitride,  $Bi_2O_3$  and  $ZnAl_2O_4$  to enhance strength between matrix and aluminoborate nanowires or whiskers (Zhang et al. 2005, Song et al. 2007, Fei et al. 2010, Liu et al. 2010a;b, Yue et al. 2010).

## 2.3 Element Boron

The metalloid boron is the fifth element of the periodic table and belongs to group 13 (boron group) with an electronic configuration of a full K shell (1s<sup>2</sup>) and three electrons in the L shell (2s<sup>2</sup> 2p<sup>1</sup>) resulting in a trivalent boron ion B<sup>3+</sup>. It has two stable isotopes, <sup>11</sup>B and <sup>10</sup>B, with a ratio of about 80:20 (Palmer and Swihart 1996). Boron does not occur in nature as pure element. Except from a few rare fluorine-minerals (e.g. avogadrite ((K,Cs)BF<sub>4</sub>), barberiite (NH<sub>4</sub>BF<sub>4</sub>), ferrucite (NaBF<sub>4</sub>), Anovitz and Grew 1996) and one nitride (cubic BN, although its genesis is unclear, Dobrzhinetskaya et al. 2008), boron is always associated with oxygen in nature (Kistler and Helvaci 2006).

The name "boron" originates from "baurach", which is the Arabian word for borax (Smith 1985) and the Persian word for white (Oganov 2010). Boron-containing minerals are known since Babylonian times when they were used as flux for gold-welding. In the Middle Ages, borax (tincal) was brought from Asia and used as welding and enameling agents. At the beginning of the  $18^{\text{th}}$  century, boric acid (H<sub>3</sub>BO<sub>3</sub>) was produced from borax, and in 1778, sassolite (natural boric acid) was discovered in fumaroles at Sasso, Italy (Anovitz and Grew 1996).

Elemental boron was recognized in 1808 from the reaction of boric acid with potassium independently by Joseph Louis Gay-Lussac and Louis Jacques Thénard, naming it bore, and by Sir Humphry Davy, who called it boracium in analogy to other metals he found. Subsequently, he discovered that the new element was not a metal, hence, Davy changed the name from boratium to boron, combining the names of borax and carbon (Anovitz and Grew 1996).

In 1864, large borax deposits were discovered in California, Nevada and Oregon, starting the highly profitable "borax era" in the USA. Until the 1980's, California was the main global supplier of boron-bearing compounds (Smith and Medrano 1996). Today, more than 80% of the world's boron demand is supplied by the USA (California), Turkey and South America (Chile, Argentina, Bolivia and Peru) (Garrett 1998, Smith 2002).

Major materials sold and processed are boric acid  $(H_3BO_3)$  and borax penta- and decahydrate  $(Na_2B_4O_5(OH)_4 \cdot 5H_2O, Na_2B_4O_5(OH)_4 \cdot 8H_2O)$  (Carpenter and Kistler 2006). Of further importance are colemanite  $(Ca_2B_6O_{11} \cdot 5H_2O)$ , ulexite  $(NaCaB_5O_6(OH)_6 \cdot 5H_2O)$  and kernite  $(Na_2B_4O_6(OH)_2 \cdot 3H_2O)$ .

#### 2.3.1 Geological Occurrence and Enrichment of Boron

Boron occurs widespread in upper crustal rocks, found in sedimentary, volcanic, plutonic and metamorphic environments. It is an enriched upper crustal element (ca. 15 ppm) and abundant in the lower crust (ca. 2 ppm), relative to the primitive mantle, which contains ca. 0.6 ppm B (Smith and Medrano 1996). The B concentration of the primitive mantle is similar to the one in C1 chondrites (ca. 0.7 ppm, Leeman and Sisson 1996). Due to hydrothermal alteration, fractional crystallization, interaction with oceanic sediments (which contain up to 100 ppm B, Vils et al. 2008) and serpentinization (interaction with seawater, Bonatti et al. 1984), estimation of the boron content of unaltered oceanic crust is difficult.

Wherever boron is observed, its distribution is very uneven. Contents in granites can reach 40 ppm, 180 ppm in oceanic sediment and up to 270 ppm in altered oceanic crust (Grew et al. 2011 and refs. therein).

Boron is soluble and mobile in aqueous solutions which transport and disperse it. Such aqueous solutions comprise e.g. salt and borax lakes (up to ca. 9000 ppm), brines (up to ca. 2200 ppm), hot springs and fumaroles (up to ca. 2000 ppm), seawater (ca. 4.5 ppm), sedimentary waters (up to ca. 2 ppm) and fresh water (up to ca. 0.01 ppm) at the Earth's surface. At depth, boron is distributed by hydrothermal fluids (Henry and Dutrow 1996 and refs. therein).

Despite its relative enrichment in the crust, boron has to get further enriched by geochemical processes (magmatic or transport in aqueous fluids) in order to allow formation of boron-bearing minerals. If the amount of available boron during crystallization is too low for formation of boron minerals (e.g. tourmaline), traces of boron will replace major elements in minerals (e.g. trace boron incorporation into feldspars, Deegan et al. 2010). With increasing boron content, partial or full substitution of  $B \rightarrow Al$  have been observed in e.g. boron-bearing feldspars (Stamatakis 1989, Higgins and Shaw 1984), boron-feldspars (reedmergenite, NaBSi<sub>3</sub>O<sub>8</sub>, Appleman and Clark 1965) and borosilicates.

Local ares, where B is concentrated in minerals and thus may form various kinds of boron deposits, can be categorized in three different processes (Watanabe 1975): (1) Endogenic processes include crystallization of boron minerals in melts and solutions, rock-alteration (metasomatism) due to magmatic and hydrothermal activity and metamorphic mobilization and concentration of boron; (2) Accumulation of boron due to volcanic activities (fumaroles, volcanoes); (3) Exogenic processes such as mechanical accumulation and chemical precipitation.

**Endogenic processes.** Melting of upper earth crust alone does not provide enough B to a melt in order to form a high content of boron minerals. Most B is brought to the melt by remelting of low metamorphosed metasediments, in which B is already present in e.g. tourmaline and sheet silicates. As long as boron is contained in the melt, it acts as flux for silicate melting, lowers the viscosity of the melt and allows it to persist to lower temperatures (Dingwell et al. 1996). Boron behaves as incompatible element in igneous systems and is concentrated in the melt by crystal fractionation. It tends to be enriched in felsic and silicic differentiates of magmatism (granites including alkaline to aluminous types and its textural variants pegmatites, porphyries and rhyolites). After formation of rock forming silicates, boron is accumulated in the residual solution until a crystal-melt equilibrium is reached or it is dissolved and expelled from the melt in an aqueous phase (London et al. 1996). As boron is a highly soluble element, boron content of igneous melts is directly related to the fluid content of the melt, and boron will escape the melt when the first fluids are released. Concentrated boron in the residual melt and/or boron concentrated in an escaping fluid reacts with wallrock forming metasomatic minerals (mostly tourmaline, axinite (Ca<sub>2</sub>(Fe<sup>2+</sup>,Mg,Mn<sup>2+</sup>)Al<sub>2</sub>BSi<sub>4</sub>O<sub>15</sub>(OH)) and other boron-bearing silicates). Amount and type of metasomatic minerals depends on the type of wallrock. In contact with carbonate formations (skarns), boron mineralization is common and a few commercial deposits have formed (Garrett 1998). Danburite (CaB<sub>2</sub>(SiO<sub>4</sub>)<sub>2</sub>) and datolite (CaBSiO<sub>4</sub>(OH)) are typical skarn borate minerals (Smith 2002).

As soon as boron is transported in fluids, the distinction between igneous and metamorphic conditions is rather difficult. In connection with metamorphic processes, boron (e.g. from boron bearing phyllosilicates in sediments) is mobilized and subsequently fixed by crystallization (Watanabe 1975).

During low-grade hydrothermal overprint (T  $< 350^{\circ}$ C), boron and other fluid-mobile elements can be leached before any textural change of the host-rock is observed.

During contact metamorphism, intrusion of fluids or magma can act as effective transfer media for boron, resulting in depletion of boron in the intrusive rock. Depending on the type of surrounding rock and the interaction with the B-rich fluid, tourmaline or skarn deposits containing other borosilicates are formed.

Subduction of either continental or oceanic crust leads to devolatilization, recrystallization and redistribution of elements within mineral assemblages. Boron mobility depends on fluid availability, P-T-stability of boron-bearing phases and the relative partitioning between these phases (Leeman and Sisson 1996).

Tournaline-group minerals are the main metamorphic borosilicates because they are commonly present in felsic metamorphic rocks (Grew 1996). Boron content is, however, not dominated by B in tournaline, but by ab- and adsorbed boron in clays and sheet silicates (especially muscovite,  $KAl_3Si_3O_{10}(OH,F)_2$ ) due to their high occurrence (Leeman and Sisson 1996). Other borosilicates are only present if the host rock is highly enriched in boron or due to special bulk composition.

**Volcanic activities.** Boron is often encountered in fumarolic vapor, in thermal springs and in hydrothermal fluids. The occurrence of these products may be related to magmatic activity. It is important to note that boron is very unevenly distributed in volcanic products. High-temperature volcanic gases can contain up to ca. 35 ppm of B, whereas the boron content of

e.g. hot springs seems to depend vastly on the country rock type (Christ and Harder 1974, Garrett 1998). Nonetheless, surface water of geothermal springs, containing boron leached from country rock, is an important continental boron source. In contrast, volcanic activities seem to play a minor role as boron source in marine deposits (Grew et al. 2011).

Sassolite  $(H_3BO_3)$  is a prominent boron mineral occurring at active volcanic sites. At Sasso (Tuscany), fumaroles mainly consisting of steam,  $CO_2$ , and only ca. 0.01% of  $H_3BO_3$ are responsible for the formation of sassolite deposits (type locality) (Garrett 1998, Anthony et al. 2003). Other fumaroles are known in active areas such as Vulcano (Lipari islands), where emitted  $H_3BO_3$  reacted with surrounding rocks to form H-, Na- and Ca-borates. After the 1851 eruption, sassolite was found at Mount Vesuvius. In Japan, sassolite encrustations are found in fumaroles of the Shows-Shinzan and Asama volcanoes (Watanabe 1975).

**Exogenic processes.** Exogenic processes leading to boron accumulation can be a result of enrichment due to mechanical accumulation (e.g. weathering resistant tourmaline is concentrated in sediments) or chemical precipitation processes. Chemical precipitation precesses can be further divided into boron concentration due to evaporation on land and evaporation of seawater (Watanabe 1975).

The principal source of terrestrial evaporite zones are long-lasting boron-rich springs. They are often related to faulted basins at active collision zones (Smith and Medrano 1996), where fluids from volcanic and hydrothermal processes have been enriched with boron due to previous metamorphism of boron-rich sediments. However, this cannot be applied as a general rule because e.g. large United States boron deposits are far away from large-scale fault zones and the reason of their occurrence is not completely known. Nevertheless, all deposits lie on local rift zones perpendicular to large-scale fault zones and appear to have been formed by geothermal springs in areas with active volcanoes (Garrett 1998). Another possibility to form boron rich waters (surface waters) is due to weathering of silicates. Whereas tourmaline is relatively weathering resistant, boron enriched materials, such as B-bearing feldspars and micas are decomposed and B is going into solution and is subsequently transported. Transport distance is related to the country rock type and weathering conditions. Boron in solution is further rapidly ab- and adsorbed if e.g. Fe- and Al-hydroxides and phyllosilicates are present (Christ and Harder 1974).

Once available, boron-rich waters are concentrated in closed basins in arid regions with high evaporation rate leading to precipitation of boron minerals (mainly borax and ulexite, possibly also colemanite, Grew et al. 2011). Due to the high mobility of boron in fluids, continental evaporite boron deposits do not survive longer geological times unless they become protected, e.g. by a layer of clay (Smith 2002). If a boron deposit is not shielded from meteoric water and groundwater, soluble boron is leached from the surface and transported into deeper sediments where new boron containing minerals form.

After initial deposition, most sedimentary deposits were altered at higher temperature

and/or pressure, often surrounded by solutions with compositions differing from the sediment composition (Smith 2002). After deposition and subsequent burial, minerals are subjected to thermal diagenesis, often combined with transformation to a lower hydration state (Smith and Medrano 1996). This process has been observed e.g. in the deposit at Boron, Kramer district, CA, where initially deposited borax transformed to kernite  $+ 6 H_2O$ . Reaction diagenesis occurs, when primary or already thermally altered borate sediments interact with solutions of different chemical composition. After thermal diagenesis, the deposit at Boron was altered by reaction diagenesis due to a contact of Ca-bearing ground water. Borax transformed to ulexite, which subsequently transformed to colemanite. As a result, cross-cutting veins of ulexite and colemanite have been observed at the deposit at Boron (Smith and Medrano 1996).

Continuous evaporation of seawater can also lead to formation of boron deposits. The rather low amount of B in seawater is concentrated and boron is mainly bound (ab- and adsorbed) by clay minerals (Watanabe 1975). In connection with salt domes, boron may be concentrated on the surface of the dome. The only large marine deposit is the Inder formation in Russia (Garrett 1998). It was most likely formed by intruding high-borate geothermal springs and its borate minerals are localized in overlying gypsum layers.

#### 2.3.2 Commercially Relevant Boron Deposits

The only two large boron deposits prospected so far, located in California (Mojave Desert) and in Turkey (provinces of Eskişehir, Kütahya and Balıkesir), are of sedimentary origin. More than 80% of the world's boron demand are covered by these two deposits, however, many small deposits exist in other countries, e.g. Russia, China, Tibet, South America (Garrett 1998, Smith 2002). Principal sedimentary borate ore minerals are borax, kernite and ulexite in the United States and borax, ulexite and colemanite in Turkey. The same minerals as well as hydroboracite (CaMgB<sub>6</sub>O<sub>8</sub>(OH)<sub>6</sub>·3H<sub>2</sub>O) are found in South American deposits. Calcium borosilicates danburite and datolite, found in skarn-type metamorphic rocks in Russia and the metamorphic magnesium borate szaibelyite (MgBO<sub>2</sub>(OH)) found in China are less important.

Processing of the boron minerals is usually done on-site by relatively simple procedures. After mining, the boron ore is sorted, ground, and subsequently leached. A hot (ca. 90°C) inorganic acid (e.g. sulfuric acid) reacts with borax, kernite, colemanite, ulexite and hydroboracite to form boric acid, which is then added to new raw material. The process is repeated until the desired purity-grade is reached (Garrett 1998). Boric oxide ( $B_2O_3$ ) can be produced from boric acid by heat treatment (Smith 1985).

#### 2.3.3 Applications of Boron

**Boron.** In its pure elemental form, boron is an important dopant for germanium and silicon in semiconductor industry and often used in pyrotechnic applications (flares, ignitors,

explosives and delay compositions) (Baudis and Fichte 1985). In nuclear reactors, boron (often as enriched <sup>10</sup>B) acts as neutron moderator and is used in emergency shutdown systems (Schwetz 1985).

**Boron oxide and boric acid.** Both compounds are intensively used in ceramic and glass industry for the production of borosilicate glasses, insulation and textile fiberglasses and as fluxing agents in porcelain enamels and ceramic glazes (Smith 1985). Minor uses of  $H_3BO_3$  are as flux in welding and soldering, as disinfectant, fungicide and insecticide (wood protection, cockroach control), water softener (Anovitz and Grew 1996), as additive in nuclear reactor cooling water (Smith 1985).

**Borates.**  $B_2O_3$  and  $H_3BO_3$  are used as fluxing agents in welding and smelting (K-, Na-, Caborates), as water softener in washing powders (borax, Hammond 1985), as neutron absorbers in nuclear reactors (Na-, Ca-borates), lubricants (K-borates). Barium borates are used as fire retardants and lithium borates are used in glass production (Smith 1985).

**Boron carbide and boron nitride.** Synthetic boron carbide and nitride compounds are very hard materials, comparable to SiC, WC and diamond. Boron carbide (composition close to  $B_4C$ , Balakrishnarajan et al. 2007) is used as abrasive material (for polishing and grinding hard materials as WC and ceramics) and in resistant engineering components (e.g. nozzles for sand-blasting equipment). Due to its lower production cost, <sup>10</sup>B-enriched  $B_4C$  has almost completely replaced elemental boron from nuclear reactors (Schwetz 1985).

Boron nitride is the second hardest material after diamond and is applied in similar applications as those of diamond (grinding, drilling, etc.). Sintered BN parts are used as cutting tools for hardened steels and in rock-drills. Because of its high chemical resistance, BN is a preferential material for strong lab ware (Schwetz 1985).

**Iron-boron alloys.** Small amounts of boron (max. 30 ppm) increase hardness of C-, Mg-, Cr-, Mo-hardened steels. It further improves the cutting properties of high-speed steels. Boron-bearing alloys are called ferroboron, if the boron content is higher than ca. 10–20 wt.-%. Boron is also used in the production of metallic glasses (Baudis and Fichte 1985).

**Other applications.** Metal borides are used for electrolytic production of aluminum  $(TiB_2)$  and as Cu deoxidizer  $(CaB_6)$ . Other applications are lab ware (crucibles, boats) and high electronic emission applications (e.g.  $LaB_6$  as electrodes in electron microscopes) (Schwetz 1985).

Organic boron compounds comprise boron halides, B-F compounds and boranes (Brotherton and Guibert 1985). Boron halides are used in organic synthesis (BF<sub>3</sub>), for the production of boron fibers (BCl<sub>3</sub>), and for polymerization of alkenes (BBr<sub>3</sub>). Fluoroboric acid HBF<sub>4</sub> is used for the production of fluoroborate salts and as etching agent for silicon and glass. Molten alkali and ammonium fluoroborates are used as fluxing agents in soldering and  $NH_4BF_4$  combined with nitro compounds is explosive. Boranes (H-B compounds) are highly reactive, often used as reducers in organic chemistry and in organic synthesis. Trimethylborate ((CH<sub>3</sub>O)<sub>3</sub>B) is the most important boric acid ester because of its use in the production of diborane. It is further used as antioxidant in fuels (Garrett 1998).

## 2.3.4 Crystal Structure of Boron

Elemental boron forms a structurally unique class of materials. The common structural unit is a  $B_{12}$  icosahedron, in which each boron atom has five neighbors (Nelmes et al. 1993). Elemental boron crystallizes in four different crystal structures.

 $\alpha$ -B<sub>12</sub> boron (space group  $R\bar{3}m$ ,  $a \approx 4.9$  Å,  $c \approx 12.58$  Å) is stable at ambient conditions and has a B<sub>12</sub> icosahedral configuration at each corner of the unit cell. The icosahedra are rather regular with B-B distances in the range of 1.75 Å to 1.81 Å (Will and Kiefer 2001).

 $\beta$ -B<sub>106</sub> boron crystallizes at ambient conditions in the same space group as  $\alpha$ -boron and is stable up to ca. 2500K (Oganov et al. 2009), however, all B icosahedra are rotated relative to the one in the center. Combined with vacancies, this results in a supercell with almost doubled lattice parameters than those of  $\alpha$ -boron (Widom and Mihalkovič 2008).

 $\gamma$ -B<sub>28</sub> has been produced by Oganov et al. (2009) at temperatures between 1800K and 2000K and at pressures higher than 19 GPa. It crystallizes in space group *Pnnm* ( $a \approx 5.05$  Å,  $b \approx 5.61$  Å,  $c \approx 6.98$  Å) and consists of B<sub>12</sub> icosahedra and B<sub>2</sub> pairs. The structure can be compared to a NaCl structure in which icosahedra and B<sub>2</sub> pairs play the role of "anion" and "cation", respectively (Oganov et al. 2009).

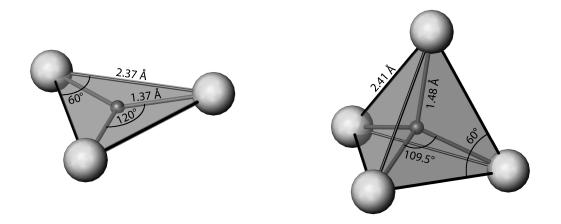
Tetragonal boron T-192 has been reported by Vlasse et al. (1979) at high temperatures and intermediate pressures. Possible space groups are  $P4_1$  or  $P4_3$  and the unit cell content consists of  $B_{21} \cdot 2B_{12} \cdot B_{2.5}$ . Due to partially filled interstitial sites, the unit-cell content varies from 190 to 192 atoms (Oganov et al. 2009).

## 2.4 Crystal Chemistry of Boron-Bearing Minerals

Boron polyhedra in crystal structures are isolated or share vertices with each other or with other tetrahedra (Anovitz and Grew 1996). The possible polyhedral arrangements of boron polyhedra have been systematized by Burns et al. (1995), Grice et al. (1995) and Hawthorne et al. (1996).

Generally, the structures of borate minerals are based on  $B(O,OH)_3$  and  $B(O,OH)_4$  polyhedra (Figures 2.23 and 2.24).  $B(O,OH)_3$  groups have an average bond valence of  $z/\rho$  (where z is the cation charge and  $\rho$  is the coordination number) of 1 valence unit, whereas a  $B(O,OH)_4$  polyhedron has 0.75 valence units. It is therefore possible that borate groups occurring as discrete oxyanions (e.g.  $BO_3^{3-}$ ) polymerize forming finite clusters, chains, sheets or frameworks. At ambient conditions, only corner-linked B-polyhedra have been observed so far (Filatov and

Bubnova 2000), however, edge-connected  $BO_x$  groups are known from structures synthesized at high pressure and high temperature (Filatov and Bubnova 2008). The clusters are referred to as "fundamental building blocks" (Burns et al. 1995). Due to different bond valence sums, it is energetically not favorable to build large clusters consisting solely of B(O,OH)<sub>4</sub> tetrahedra.



BO<sub>3</sub><sup>3-</sup> equilateral triangle. Bond lengths and an- BO<sub>4</sub><sup>5-</sup> tetrahedron. Bond lengths and angles are gles are indicated.

Figure 2.23: Structural drawing of an ideal Figure 2.24: Structural drawing of an ideal indicated.

It is interesting to point out that an isotope fractionation of <sup>11</sup>B and <sup>10</sup>B occurs between trigonal and tetrahedral coordinated boron species. The heavier <sup>11</sup>B is preferentially three-fold coordinated, whereas the <sup>10</sup>B isotope is concentrated in four-fold coordinated sites. The isotope fractionation factor ( $[^{11}B/^{10}B_{trigonal}]/[^{11}B/^{10}B_{tetragonal}]$ ) is ca. 1.03 at room temperature and decreases with temperature (Palmer and Swihart 1996, Kakihana et al. 1977).

From refined crystal structures containing B(O,OH)<sub>3</sub> and/or B(O,OH)<sub>4</sub> groups, Hawthorne et al. (1996) investigated central-ligand bond lengths and angles. The average center-ligand distance in  $B(O,OH)_3$  triangles is 1.370 Å, with a rms deviation of 0.017Å. A Gaussian distributed minimum of 1.322 and maximum of 1.428 Å has been observed with a variation range of 0.106 Å. In B(O,OH)<sub>4</sub>, the mean distance is 1.476 Å with a rms deviation of 0.025 Å. The Gaussian distributed minimum and maximum observed distances in  $B(O,OH)_4$  are 1.397 Å and 1.512 Å, respectively with a variation of 0.115 Å. For an ideally coordinated  $B(O,OH)_3$ group, coordination would represent an equilateral triangle with (O,OH)-B-(O,OH) angles of 120° and (O,OH)-(O,OH)-(O,OH) angles of 60°. B-(O,OH)-B angles in an ideally coordinated  $B(O,OH)_4$  polyhedron are 109.47°. Actual angles depend on the linkage of  $BO_x$  groups with neighboring polyhedra.

Isolated  $BO_3$  and  $BO_4$  polyhedra have practically no thermal expansion when heated and behave similar to  $AlO_4$  and  $SiO_4$  tetrahedra (Filatov and Bubnova 2000; 2008, Hazen and Prewitt 1977, Tucker et al. 2000). Crystal structures containing  $BO_x$  groups can be regarded as special because not only isolated BO<sub>x</sub> groups but also polymerized groups of BO<sub>x</sub> polyhedra remain rigid at high temperature. However, with rising temperature, single  $BO_x$  groups or rigid polymerized  $BO_x$  group clusters act as structural hinges in order to allow structural changes. Thermal invariability of the  $BO_x$  groups or clusters is the main reason for highly anisotropic and sometimes even negative thermal expansion of  $BO_x$  containing crystals. At high-pressure conditions, borates become denser due to decreasing bond lengths, increased coordination number and  $BO_4$  polymerization via O-O edges (Filatov and Bubnova 2008).

## 2.5 Analytical Techniques for Boron Analysis

Boron was recognized as mineral constituent in the late 18<sup>th</sup> century. More or less reliable wetchemical analyses were possible only about 100 years later, however, they remained a challenge until today (e.g. due to boron loss at higher temperatures). Easy and fast quantitative analysis of boron is still difficult because B is not detected with most X-ray fluorescence methods and also electron microprobes require special analyzer crystals (Anovitz and Grew 1996).

Electron microprobes can be equipped with thin-window detectors for energy-dispersive spectrometry (EDS) and layered synthetic microstructure (LSM, optimized for low Z elements) crystals for wavelength dispersive spectrometry (WDS). With EDS methods, boron can be identified, however, quantification is almost impossible due to the low yield of X-rays of boron. Even with WDS, the quantification of light elements is a challenge, and it largely depends on the instrument setup and the matrix of the sample and the standard (McGee 1996). Electron microprobe measurements have been performed on aluminoborate samples early in this study, however, results were not reproducible and this method was not further pursuited.

Quantitative boron isotope analyses can be performed with high precision with magnetic mass spectrometers. The sample is usually thermally ionized and positive or negative boroncomplex ions are measured. Standardization, sample preparation and instrumental calibration and setup are difficult and require much time (Swihart 1996). Especially during sample preparation (boron must be isolated from the sample) contamination and boron-loss are an issue. SIMS (secondary ion mass spectrometry) is a favorable method for boron quantification because it combines the relatively easy sample preparation of the electron microprobe with high-precision results of mass spectrometry. Standardization is problematic and a calibration curve from known standards is needed. Once set up, due to the high quantification precision, measurements of diffusion coefficients are possible, as well as element mapping and isotope measurements (Hervig 1996). Quadrupole mass spectrometers are less sensitive but easier to operate and may thus be used to good advantage for certain experiments (Swihart 1996). It has been shown in this study, that laser-ablation coupled with an ICP (inductively coupled plasma) quadrupole mass spectrometer allows good estimation of boron content, however not boron isotope measurements. The advantage of this method is that time-consuming and complex sample preparation could be avoided.

X-ray diffraction measurements with conventional single-crystal laboratory diffractometers

offer the possibility to locate boron in crystal structures. Correlations with (anisotropic) displacement parameters and the overall scale factor make precise quantification (site occupancy factors) difficult. Due to the low amount of information obtained from powder X-ray diffraction measurements, it is doubtful if boron sites can be identified and occupancy fractions of boron sites are not reliable.

Fourier transformed infrared, Raman and nuclear magnetic resonance (NMR) spectroscopy do not allow straightforward and precise quantification of boron content, however, both provide evidence on boron occurrence in a sample and unambiguous information about the coordination of boron can be obtained. In Raman and FTIR spectroscopy, the difference in the mass of the two B isotopes (10 rel.-%) has to be accounted for (Siebert 1966, Farmer 1974, Harris and Bertolucci 1989). In NMR experiments, <sup>10</sup>B and <sup>11</sup>B have different resonance frequencies and can therefore be measured independently. Due to different frequency and quantum spins, <sup>10</sup>B measurements are less sensitive (D. Rentsch, personal communication). In this study, all three methods were used to obtain additional structural information of the investigated compounds.

Further techniques for boron analysis include nuclear methods and parallel electron energyloss spectroscopy (EELS). With EELS, a sample is probed with monoenergetic electrons. After interaction with the sample, electron energy distribution is analyzed. The resulting spectra provide information on elemental composition, coordination and site symmetry, valence, spin state and bonding of atoms at very high spatial resolution. The high potential of EELS can be increased in combination with transmission electron microscopy (Lawrence and Buseck 1996). Nuclear methods comprise particle-induced gamma-ray (PIGE) and particle emission (NRA), prompt-gamma neutron activation analysis (PGNAA) and fast neutron activation analysis (FNAA). All of these methods have the great advantage that the bulk sample (or depths up to 100  $\mu$ m in thin sections) is analyzed. Furthermore, matrix effects are subordinate. Nuclear methods are based on nuclear reactions due to interaction of nuclei with charged particles or neutrons. They offer precise information on isotope type and concentration in a bulk sample. However, particle accelerators or nuclear reactors are required and minimum sample mass is rather large (up to ca. 1000  $\mu$ g) when neutrons are used (Robertson and Dyar 1996).

## **3 Boron Mullite Synthesis Experiments**

Synthesis experiments in this study were performed in a Nabertherm LHT04/17 furnace equipped with Superkanthal (MoSi<sub>2</sub>) heating elements in air atmosphere. Charges were placed in lid covered platinum crucibles or in sealed platinum capsules cut from a platinum tube with 4 mm inner diameter. Solid starting materials were thoroughly ground in agate mortars and pressed as tablets before heating. Prior to further investigations, excess  $B_2O_3$  or  $H_3BO_3$  was washed away with deionized  $H_2O$  in samples with aluminoborate compositions. Chemical compounds used for precursor preparation are listed in Table 3.1.

Preliminary sample identification was done by powder X-ray diffraction analyses in combination with the ICDD PDF-2 2004 database. As the phases  $Al_5BO_9$  and  $Al_{18}B_4O_{33}$  cannot be distinguished by X-ray diffraction methods, the term "ABO" is used in this chapter for  $Al_5BO_9$  and  $Al_{18}B_4O_{33}$  in space group  $Cmc2_1$ .  $Al_4B_2O_9$  in space group C2/m was identified using structural information from Fischer et al. (2008b). Identification of mullite-like phases was done with LeBail refinements starting with lattice parameters a = 7.6 Å, b = 7.5 Å and c = 2.8 Å in space group *Pbam*.

Chemical compound, formula	Catalogue no.	Vendor
Aluminium nitrate nonahydrate 98+%, $\rm Al(NO_3)_3\cdot 9H_2O$	237973-100G	Sigma-Aldrich
Aluminium nitrate nonahydrate p.a., $\rm Al(\rm NO_3)_3\cdot 9 H_2O$	1.01063.0500	Merck
Aluminum oxide -100 mesh 99%, $\mathrm{Al}_2\mathrm{O}_3$	$23,\!474-5$	Sigma-Aldrich
Aluminium sulfate hexaceda hydrate p.p.a, $\rm Al_2(SO_4)_3\cdot 16H_2O$	06421	Fluka
Borax p.a., $\mathrm{Na_2B_4O_5(OH)_4} \cdot 10\mathrm{H_2O}$	6308	Merck
Boric acid p.a., $H_3BO_3$	165	Merck
Boron oxide 97.5% min., $\mathrm{B_2O_3}$	012290 500G	Alfa Aesar
Boron oxide 99.98%, $\mathrm{B_2O_3}$	339075-100G	Sigma-Aldrich
Glycerol 85%, $C_3H_5(OH)_3$	07-3800-01	Haenseler AG
Molybdenum(VI) oxide p.p.a., ${\rm MoO}_3$	$69850 \ 100 { m g}$	Sigma-Aldrich
Potassium carbonate p.a., $\mathrm{K}_{2}\mathrm{CO}_{3}$	1.04928.0500	Merck
Potassium sulfate p.a., $K_2(SO_4)_2$	5153	Merck
Silica gel, grade 62, 60–200 mesh, 150Å, ${\rm SiO}_2$	112926-00-8	Aldrich
Tetra ethyl orthosilicate (TEOS), $\mathrm{C_8H_{20}O_4Si}$	$91689\text{-}250\mathrm{ML}$	Sigma-Aldrich
Yttrium oxide 99.99%, $\mathrm{Y}_{2}\mathrm{O}_{3}$	$205168\text{-}50\mathrm{G}$	Sigma-Aldrich

Table 3.1: Chemical compounds used in this study.

## 3.1 Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub> Compounds with Mullite-Type Structure

## 3.1.1 Overview

Until the late  $20^{\text{th}}$  century, routes for the synthesis of aluminoborates with mullite-type structures (Al<sub>5</sub>BO<sub>9</sub> (5Al<sub>2</sub>O<sub>3</sub>:B<sub>2</sub>O<sub>3</sub>), Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> (9Al<sub>2</sub>O<sub>3</sub>:2B<sub>2</sub>O<sub>3</sub>) and Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> (4Al<sub>2</sub>O<sub>3</sub>:2B<sub>2</sub>O<sub>3</sub>)) aimed at the production of single-crystal or powder samples for subsequent structural or crystal-chemical analyses and for the investigation of material properties. In this context, "single-crystal" is used for crystals with dimensions suitable for measurements with conventional laboratory single-crystal X-ray diffractometers (a few 100 µm or larger), whereas the term "powder" is used for loose crystals or sintered aggregates of crystals with a few µm length, which are too small to be measured with single-crystal X-ray diffractometers.

Since the  $21^{st}$  century, methods for large-scale synthesis of ABO and  $Al_4B_2O_9$  as advanced materials were investigated. This includes production of µm-large whiskers, micro- and nanotubes, nanorods and nanowires with and without coatings for specialized applications (Cheng et al. 2004, Ma et al. 2002, Kim et al. 2004, Zhang et al. 2005, Li and Chang 2006, Peng et al. 2006, Zhang et al. 2006, Elssfah et al. 2007, Song et al. 2007, Wang et al. 2008, Lee et al. 2010 and more). Preparation methods for such specialized materials are not considered here.

Four different methods for single-crystal or powder synthesis of aluminoborates ABO and  $Al_4B_2O_9$  are repeatedly reported in literature: (1) Synthesis from  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub> with additional excess B<sub>2</sub>O<sub>3</sub> as flux; (2) Synthesis of an aluminum and a boron source ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, Al(OH)<sub>3</sub>, Al<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub>,... and B<sub>2</sub>O<sub>3</sub> or H<sub>3</sub>BO<sub>3</sub>) in a flux of different composition; (3) Solid-state reactions of  $\alpha$ - or  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub> or H<sub>3</sub>BO<sub>3</sub>; (4) Synthesis from amorphous precursors prepared via sol-gel routes.

#### Flux of different chemical composition (K<sub>2</sub>CO<sub>3</sub>, MoO<sub>3</sub>, Y<sub>2</sub>O<sub>3</sub>, K<sub>2</sub>SO<sub>4</sub>, CaO)

ABO single-crystals have been synthesized by Sokolova et al. (1978) from  $Al_2O_3$  and  $B_2O_3$ in a  $K_2CO_3$ -MoO<sub>3</sub>-Y<sub>2</sub>O<sub>3</sub> flux with unknown composition by cooling the melt from 1150°C to 950°C.

Azizov et al. (1982) and Nekrasova and Leonyuk (2008) showed that ABO may crystallize during synthesis of (Yb,Y)Al<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> and YAl<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> from K<sub>2</sub>MoO<sub>3</sub>O<sub>10</sub> flux (= K<sub>2</sub>CO<sub>3</sub> +  $3MoO_3$ ) as a side phase. Starting from a mixture CaO, Al<sub>2</sub>O<sub>3</sub> and H<sub>3</sub>BO<sub>3</sub> with johachidolite composition (CaAlB<sub>3</sub>O<sub>7</sub>, Kadiyski et al. 2008), Leonyuk and Shvanskiy (1998) produced ABO. The same observations were made when K<sub>2</sub>MoO<sub>4</sub> flux was added.

Ihara et al. (1980) reported prismatic crystals of up to 1 mm length from a mixture of 25.0 mol-%  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, 43.8 mol-% B<sub>2</sub>O<sub>3</sub> and 31.2 mol-% CaO (ratio Al<sub>2</sub>O<sub>3</sub>:B<sub>2</sub>O<sub>3</sub> = 4.5:8) by cooling the melt from 1440°C to 1250°C.

ABO whiskers of a few  $\mu$ m length were produced from Al(OH)<sub>3</sub> and H<sub>3</sub>BO<sub>3</sub> in KCl flux by cooling the melt from 900°C and from Al<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub> and H<sub>3</sub>BO<sub>3</sub> in K<sub>2</sub>SO<sub>4</sub> flux by cooling from 1100°C (Wada et al. 1991). A similar procedure was used by Gönenli and Tas (2000), who

dissolved  $Al_2(SO_4)_3 \cdot 18 H_2O$  and  $H_3BO_3$  with  $K_2SO_4$  in water. The material obtained from the dried solution was heated at 1073°C and 1150°C. In a TEM study, Carazeanu et al. (2004) determined formation temperature of ABO nanocrystals at around 900°C. The crystals were produced from  $Al_2(SO_4)_3$ ,  $H_3BO_3$  and  $K_2SO_4$  with  $K_2CO_3$  or KCl as fluxing agent. The optimum heating temperature found was 1000°C for 50%  $K_2SO_4$  flux.

## Isochemical flux (B<sub>2</sub>O<sub>3</sub>)

Benner and Baumann (1938) synthesized ABO crystals of up to 2 cm length (reported by Baumann and Moore 1942) by melting  $B_2O_3$  in molten  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> using an electric arc furnace.

By heating 30 wt.-%  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> in 70 wt.-% molten B<sub>2</sub>O<sub>3</sub> at 1100°C, Scholze (1956) was able to crystallize ABO. With the same precursor, he produced Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> at 1000°C. As Scholze (1956) measured optical properties of the materials, it is likely that he obtained large crystallites, however, crystal size is not mentioned.

#### Solid-state reaction

Dietzel and Scholze (1955) synthesized ABO crystals by heating a stoichiometric mixture of  $9\alpha$ -Al<sub>2</sub>O<sub>3</sub>:2B<sub>2</sub>O<sub>3</sub> in a molybdenum boat under reducing atmosphere and Ray (1992) produced ABO powder by calcining the same mixture at 1300°C. His observations showed that at temperatures lower than 700°C, no aluminoborate phase formed. Between 800°C and 1000°C, ca. 10% of the mixture reacted to Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> and between 1000°C and 1100°C, formation of ABO started at the expense of Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>. After 1100°C, only ABO was observed. The temperatures are in agreement with those Rymon-Lipinski et al. (1984) determined with thermogravimetric analyses and those of the phase diagram by Gielisse and Forster (1962), which was derived by preparing samples from  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub> in sealed platinum capsules between 600°C and 1650°C.

Garsche et al. (1991) obtained ABO powder as well as single-crystal samples from  $9\alpha$ -Al<sub>2</sub>O<sub>3</sub> and 2B<sub>2</sub>O<sub>3</sub>. Powder samples were produced by annealing the mixture at 1350°C for 2h. Single-crystals were obtained from the same starting material, however, heat treatment was performed at 1500°C for 7 days in sealed platinum capsules.

#### Sol-gel derived precursors

From a precursor made from a denitrified gel of stoichiometric amounts of  $Al(NO_3)_3 \cdot 9 H_2O$ and  $H_3BO_3$ , Mazza et al. (1992) produced metastable mullite-like " $Al_5BO_9$ " and " $Al_4B_2O_9$ " in space group *Pbam* at 880°C and 917°C, respectively. Further heating of the products above 900°C – 1000°C yielded  $Al_4B_2O_9$  and ABO. The same compounds were also produced and investigated with powder X-ray diffraction, <sup>11</sup>B and <sup>27</sup>Al MAS NMR by MacKenzie et al. (2007). In a study on Al-B-Si compounds, both compounds were also found by Griesser et al. (2008) from a similar precursor and Fischer et al. (2008b) used the preparation method of Mazza et al. (1992) to crystallize  $Al_4B_2O_9$  from a non-stoichiometric ( $Al_2O_3:2B_2O_3$ ) starting mixture at 950°C.

#### Precipitation process derived precursors

To avoid boron evaporation at higher temperatures, Duoy (2005) dissolved  $Al(NO_3)_3 \cdot 9H_2O$ and  $H_3BO_3$  in distilled  $H_2O$ . The aqueous solution was subsequently precipitated into a mixure of ammonium bicarbonate and ammonium carbamate (yielding solid ammonium borate and ammonium nitrate) before the liquid was evaporated. The resulting precursors were heated at 1000°C, when ABO and Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> were found and at 1100°C, ABO formed together with  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>. At temperatures below 950°C, formation of transition alumina phases ( $\gamma$ -Al<sub>2</sub>O<sub>3</sub>, Chakravorty 1993, Chakraborty and Das 2003 and Gutierrez et al. 2002) is reported.

# 3.1.2 Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub> Synthesis Experiments

## Flux of different chemical composition (K<sub>2</sub>CO<sub>3</sub>, MoO<sub>3</sub>, Y<sub>2</sub>O<sub>3</sub>, K<sub>2</sub>SO<sub>4</sub>, Na<sub>2</sub>B<sub>4</sub>O<sub>7</sub>·10H<sub>2</sub>O)

Assuming that Sokolova et al. (1978) were performing synthesis experiments for preparation of an Al-B-Y-O phase from a  $K_2CO_3 + 3MoO_3$  flux, the first experiment for synthesizing ABO single crystals (# 1) was performed by starting from  $9\alpha$ -Al<sub>2</sub>O<sub>3</sub>, 2B<sub>2</sub>O<sub>3</sub> and K<sub>2</sub>CO<sub>3</sub> with  $3MoO_3$  yielding  $K_2Mo_3O_{10}$  as flux. Weighed ratios and heating rates are shown in Table 3.2. Elongated prismatic ABO single-crystals of a maximum length of about 1 mm and a rather smooth surface, intergrown in radial aggregates (Figures 3.1 - 3.3) could easily be separated from the flux in boiling H<sub>2</sub>O. The crystals were mostly twinned, as estimated from polarized light microscopy, and Al<sub>2</sub>O<sub>3</sub> inclusions could be identified with electron microprobe analyses (Figure 3.4). The inclusions are the reason for excess corundum as identified with an X-ray powder diffraction pattern of ground crystals (Figure 3.5). SEM images of the crystals of this synthesis experiment are shown in Figures 3.6 and 3.7.

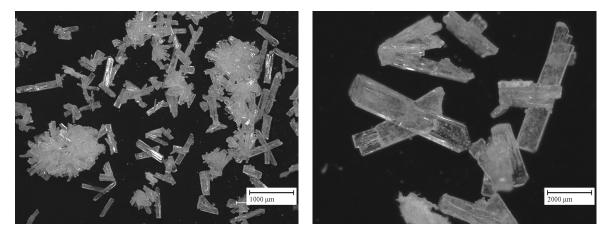


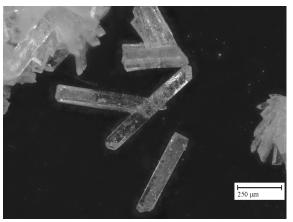
Figure 3.1: Microscope photograph of the ABO Figure 3.2: Microscope photograph of the ABO crystals and radially grown crystal aggregates pro- crystals from run #01 at higher magnification. duced in run #01.

	Table 3.2	List of	: ABO s	ynthesis expe	riments	perform	ed with	$K_2CO_3 + 3$	$3MoO_3$ f	lux.
#	$\alpha\text{-}\mathrm{Al}_2\mathrm{O}_3$	$\mathrm{B}_2\mathrm{O}_3$	A:B	(AB):Flux	Wt.	$T_{\rm max}$	$T_{\rm end}$	$\nabla$	Type	Phases
	(g)	(g)	(mol)	(wt.)	(g)	$(^{\circ}C)$	$(^{\circ}C)$	$(^{\circ}C/h)$		(XRD)
01	0.197	0.03	9:2	0.111	2.29	1100	600	10	Х	1, 2
02	0.659	0.20	9:4	0.368	2.92	1100	600	10	Р	1, 2
04	0.395	0.06	9:2	0.152*	3.45	1100	700	5	Р	1, 2
05	0.395	0.07	9:2.2	0.225	2.61	1100	600	5	Р	1, 2
06	1.977	0.30	9:2	0.083	29.66	1150	600	3	Р	1, 2
07	2.372	0.36	9:2	0.111	27.47	1100	600	10	Р	1, 2, 3
08	2.768	0.42	9:2	0.129	27.93	1100	600	10	Р	1,2,3,4
09	1.713	0.26	9:2	0.098	21.97	1100	600	10	Р	1, 2, 4
10	1.858	0.282	9:2	0.086*	26.88	1100	400	10	Р	2, 4, U
11	1.858	0.282	9:2	0.086	26.88	1100	400	10	Р	1, 2, 4
12	0.395	0.06	9:2	0.111	4.58	1100	600	10	Р	1,2,3,4
13	3.295	0.50	9:2	0.460	12.04	1100	450	10	Р	1, 2, U
16	0.659	0.10	9:2	0.153	5.71	1100	450	10	Р	1, 2
18	1.318	0.20	9:2	0.153	11.41	1100	450	10	Р	1, 2, 4, 5
19	0.659	0.14	9:2.8	0.153	6.18	1100	450	10	Р	1,  4,  5
20	0.659	0.14	9:2.8	0.153	6.18	1250	1000	10	Р	1,2,3,4
21	0.659	0.14	9:2.8	0.153	6.18	1250	1000	10	Р	1, 4
22	0.659	0.14	9:2.8	0.115	7.76	1100	600	10	Р	1, 4, U
31	0.329	$0.08^{+}$	5:2	0.099	4.53	1100	600	10	Р	1, 4, 5
45	0.197	0.20	2:3	0.111	2.46	1100	800	5	Р	1, 2, U
46	0.197	0.20	2:3	0.111	2.46	1100	750	5	Р	1, 3
$\mathbf{X} = 0$	crystals, $P =$	powder			1 = ABO			$4 = \text{KAl}(\text{MoO}_4)_2$		
* sing	gle-phase K <sub>2</sub>	$Mo_3O_{10}$ a	as flux		$2 = \alpha - A$	$Al_2O_3$		$5 = \mathrm{K}_{2}\mathrm{Mo}_{3}\mathrm{O}_{10}\cdot 3\mathrm{H}_{2}\mathrm{O}$		
$\dagger H_3 H$	$\mathrm{BO}_3$ instead	of $B_2O_3$			$3 = K_2$	$\mathrm{Mo}_4\mathrm{O}_{13}$		U = unic	lentified	phase(s)

Table 3.2: List of ABO synthesis experiments performed with  $K_2CO_3+3MoO_3$  flux

In another 20 runs with  $K_2CO_3 + 3MoO_3$  flux, the total precursor amount and the ratio  $(Al_2O_3+B_2O_3)$ :flux were varied from ca. 0.46 to 0.08, up to 100 wt.-% of excess boron were added and the  $K_2CO_3 + 3MoO_3$  flux was replaced by previously prepared  $K_2Mo_3O_{10}$ .  $T_{max}$  was varied from 1100°C to 1250°C,  $T_{end}$  from 400°C to 1000°C and the cooling gradient from 3°C/h to 10°C/h (Table 3.2) but the synthesis yielding large single-crystals could not be reproduced. Samples consisted either of only small crystals of ABO powder or of sintered tablets (Figure 3.8), accompanied by excess corundum (Figure 3.10), or of large, yellowish KAl(MoO\_4)\_2 crystals (Figure 3.9) together with small ABO crystals and excess corundum (Figure 3.11). With the  $K_2CO_3 + 3MoO_3$  flux method, no synthesis run yielded single-phase ABO.

Similar to Gönenli and Tas (2000) and Wada et al. (1991), a synthesis with  $H_3BO_3$ ,  $Al_2(SO_4)_3 \cdot 16 H_2O$  and  $K_2SO_4$  was performed. The precursor material consisted of 7.32g of  $Al_2(SO_4)_3 \cdot 16 H_2O$  and 0.35g  $H_3BO_3$  (molar ratio of  $8Al_2O_3:2B_2O_3$ ) and 5g  $K_2SO_4$ . The compounds were diluted in deionized  $H_2O$  which was evaporated after the solution was homogeneous. The obtained powder was subsequently heated at 1100°C and slowly cooled to 450°C at 10°C/h. The synthesized product consisted of ABO,  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and at least one unidentifiable phase.



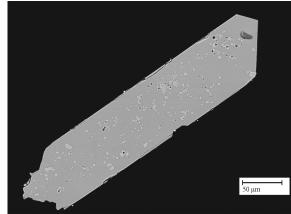


Figure 3.3: Microscope photograph of single ABO crystals from run #01. The bright spots in the crystals were identified as  $\rm Al_2O_3$  inclusions.

Figure 3.4: Electron microprobe photograph of an ABO crystal from run #01. The bright spots were identified as  $\mathrm{Al}_2\mathrm{O}_3$  inclusions.

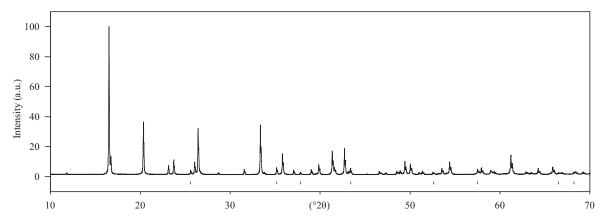


Figure 3.5: Powder XRD pattern of ground ABO crystals from run #01. hkl-ticks of corundum are indicated.

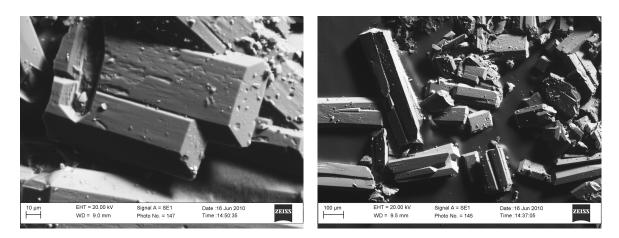
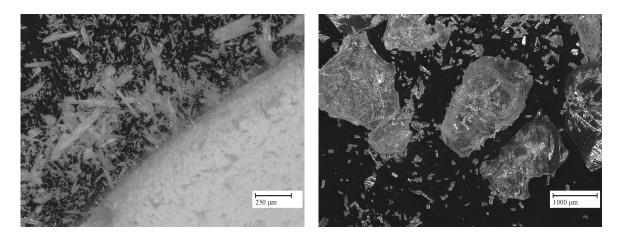


Figure 3.6: Secondary electrons SEM photo- Figure 3.7: Secondary electrons SEM photograph of the ABO crystals from run #01 showing the rather smooth crystal surface.

graph of the ABO crystals from run #01.



tered tablet of ABO crystals and corundum from  $KAl(MoO_4)_2$  crystals from run #09. run #05.

Figure 3.8: Small crystals and a piece of a sin- Figure 3.9: Small ABO crystals and large

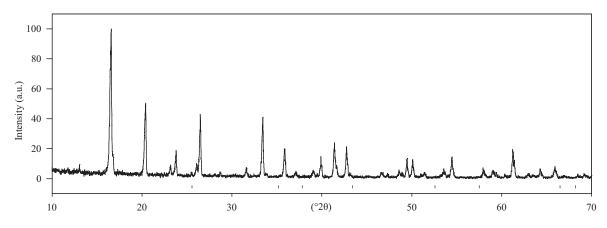


Figure 3.10: Powder XRD pattern of the ground sintered tablet of ABO crystals shown in Figure 3.8 from run #05.  $hkl\mbox{-ticks}$  of corundum are indicated.

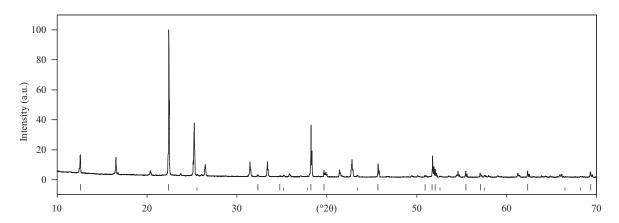


Figure 3.11: Powder XRD pattern of the large  $KAl(MoO_4)_2$  crystals, ABO and corundum, shown in Figure 3.9. Short hkl-ticks indicate corundum peaks, long hkl-ticks indicate peaks assigned to  $\mathrm{KAl}(\mathrm{MoO}_4)_2.$ 

Sokolova et al. (1978) used a mixture of  $Y_2O_3$ - $K_2CO_3$ - $MoO_3$  as flux without indicating molar ratios. Nekrasova and Leonyuk (2008) and Azizov et al. (1982) showed that ABO may crystallize during synthesis of (Yb,Y)Al<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> from  $K_2MoO_3O_{10}$  flux. Y and Er have similar ionic radii, thus, the molar amount of  $Er_2O_3$  was replaced by  $Y_2O_3$  in a recipe for the synthesis of  $ErAl_3(BO_3)_4$  crystals in a  $K_2Mo_3O_{10}$  flux (Teshima et al. 2006 and Beregi et al. 1999). The syntheses yielded ABO with  $KY(MoO_4)_2$  and ABO with  $KY(MoO_4)_2$  together with  $K_2Mo_3O_{10}$ . The presence of  $K_2Mo_3O_{10}$  was apparently due to incomplete washing in  $H_2O$  and  $KY(MoO_4)_2$  was also found in similar experiments by Azizov et al. (1982). After washing in hot HCl, ABO crystals with  $\mu$ m size remained as the only solid.

Syntheses of ABO with borax  $(Na_2B_4O_7 \cdot 10H_2O)$  as flux were prepared as listed in Table 3.3. In all runs, a glass remained in the platinum crucible which was soluble in warm  $HNO_3$ . No crystals were found after dissolving the glass. However, few very thin crystals grew at the rim of the platinum crucible but they could not be removed. Scratched-off pieces of the crystalline products were identified as ABO.

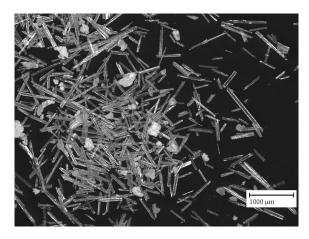
#	$\stackrel{\rm Al_2O_3}{\rm (g)}$	$\substack{\mathbf{B_2O_3}\\(\mathbf{g})}$	A:B (mol)	(AB):Flux (wt%)	Wt. (g)	$\begin{array}{c} T_{\max} \\ (^{\circ}\mathrm{C}) \end{array}$	${T_{ m end}}$ (°C)	$ abla (^{\circ}C/h) $	Type	Phases (XRD)
40	0.764	0.104	5:1	80:20	3.53	1100	900	3	X, G	ABO
41	1.274	0.174	5:1	60:40	3.35	1100	900	3	X, G	ABO
42	1.784	0.243	5:1	40:60	3.17	1100	900	3	X, G	ABO
43	1.019	0.139	5:1	70:30	3.45	1100	900	3	X, G	ABO
44	0.892	0.121	5:1	75:25	3.49	1050	850	2	X, G	ABO

Table 3.3: List of ABO syntheses performed with borax  $(Na_2B_4O_7 \cdot 10H_2O)$  as flux.

X = crystals, G = glass

## Isochemical flux (B<sub>2</sub>O<sub>3</sub>)

Four different synthesis routes were tested starting from  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and H<sub>3</sub>BO<sub>3</sub>: (1) A pressed tablet of 0.2g Al<sub>2</sub>O<sub>3</sub> and 0.8g H<sub>3</sub>BO<sub>3</sub> was covered with ca. 3g H<sub>3</sub>BO<sub>3</sub> in a platinum crucible. The charge was heated at 1500°C, then heated to 1550°C and held at this temperature for one hour and subsequently quenched. Yield were small crystals at the rim of the crucible. (2) A pressed tablet of 0.2g Al<sub>2</sub>O<sub>3</sub> and 0.8g H<sub>3</sub>BO<sub>3</sub> was fired for 2h at 1200°C. The ground ABO powder was subsequently mixed with more H<sub>3</sub>BO<sub>3</sub> and pressed as tablet, which was covered with ca. 3g H<sub>3</sub>BO<sub>3</sub> in a platinum crucible and heated for 1h at 1550°C and then slowly cooled to 1250°C at 15°C/h. Yield were ABO crystals (Figure 3.12 and 3.13) which could be removed from the crucible by dissolving remaining B<sub>2</sub>O<sub>3</sub> in hot water. The crystals were coexisting with loose aggregates of tiny ABO crystals (Figures 3.14 and 3.15). SEM images of the crystals show that the crystals have a rather flaky surface (Figures 3.16 and 3.17). From powder X-ray diffraction patterns, the product was identified as single-phase ABO (Figure 3.18). (3) A presintered piece of 0.1g ABO was put into a platinum tube. The tube was filled with H<sub>3</sub>BO<sub>3</sub> which was heated to obtain B<sub>2</sub>O<sub>3</sub>, sealed and subsequently heat treated at 1500°C for 72h and then cooled to 1200°C in 72h. After the heat treatment, the presintered piece of ABO remained almost unaltered (Figure 3.19). (4) A platinum tube was filled with roughly 50:50 wt.-%  $Al_2O_3$  and  $B_2O_3$ . The sealed tube was heated for 72h at 1500°C and then cooled to 1250°C in 72h. After the heat treatment, the tube contained pieces of  $\rm B_2O_3$ in which small crystals were trapped (Figure 3.20). After dissolving  $B_2O_3$  in hot water, few small ABO crystals remained together with corundum.



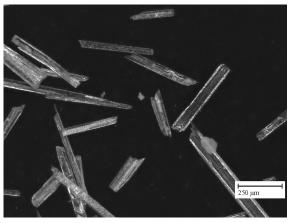


Figure 3.12: ABO crystals grown from the B<sub>2</sub>O<sub>3</sub> Figure 3.13: ABO crystals grown from the B<sub>2</sub>O<sub>3</sub> flux (#62-2).

flux (#62-2).

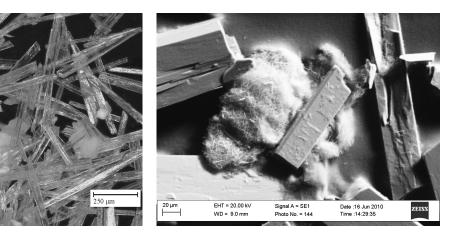


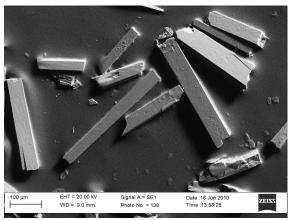
Figure 3.14: ABO crystals grown from the B<sub>2</sub>O<sub>3</sub> flux coexisting with aggregates of tiny needle-like ABO crystals (#62-2).

Figure 3.15: Secondary electrons SEM picture of ABO crystals from run #62-2 coexisting with aggregates of tiny ABO crystals.

## Solid-state reaction

ABO powder samples were prepared by solid-state reactions of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub> or H<sub>3</sub>BO<sub>3</sub>.

Sample #24 (Table 3.4) was prepared according to Garsche et al. (1991) but one sample with a stoichiometric ratio of  $Al_2O_3:B_2O_3$  of 5:1 (#24a) and another with a ratio of 9:2 (#24b). After 15 days of heating at 1500°C of 0.36g of the precursor, a sintered piece of ABO



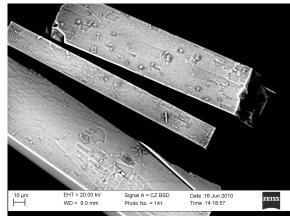


Figure 3.16: Secondary electrons SEM picture of ABO crystals from run #62-2.

Figure 3.17: Backscattered electrons SEM picture showing the rather flaky surface of the crystals from run #62-2.

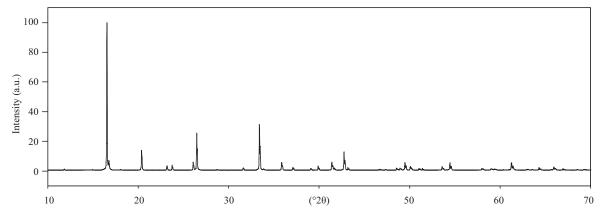


Figure 3.18: Powder XRD pattern of ground ABO grystals grown from the  $B_2O_3$  flux (#62-2). The different relative intensities when compared to powder XRD patterns of other samples are ascribed to preferred-orientation effects due to a different sample preparation method.

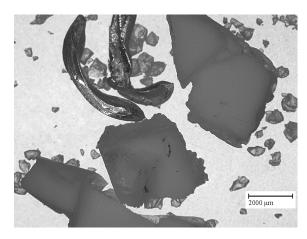
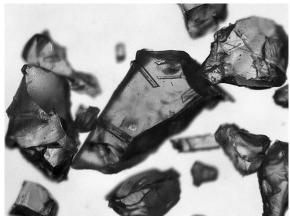


Figure 3.19: Pieces of the Pt capsule and sin- Figure 3.20: Small ABO crystals trapped in tered ABO after heat treatment with  $B_2O_3$  in a pieces of  $B_2O_3$  (run #62-4). Length of the crystals sealed platinum tube (run #62-3).



is ca.  $0.1~\mathrm{mm}.$ 

was found in the sealed platinum tube with the starting mixture with a 5:1 ratio, whereas the other platinum tube exploded in the furnace. Weighing the sealed capsule before and after the heat treatment indicated that it was not airtight.

Sample #38 was produced in an Anton-Paar furnace and in-situ measured with X-ray powder diffraction to monitor the solid-state reaction. Heating rate was  $10^{\circ}$ C/min and measurements were made at 25°C and then each 100°C up to 1100°C. An additional measurement was performed at 450°C (melting point of B<sub>2</sub>O<sub>3</sub>). Each measurement took 20 min. After heating the sample for 2h at 1100°C, it was cooled to 25°C and measured again. A selection of the temperature dependent diffraction patterns is shown in Figure 3.21. Crystallization temperatures are in agreement with previous findings.

$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$								*	-	
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	#	2 3	2,0	0, 0					•	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	23	3.295	0.5		9:2	3.80	1200	800	100	ABO, 7% $\alpha$ -Al <sub>2</sub> O <sub>3</sub>
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	24a	7.322	1.0		5:1	8.32	1500	1500	-, 7d	ABO, 30% α-Al <sub>2</sub> O <sub>3</sub>
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	24b	6.590	1.0		9:2	7.59	1400	1400	-, 7d	-
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	28	0.2		1.1	2:9	1.3	1200	400	20	ABO
50 3.66 0.50 5:1 4.16 1200 800 100 ABO, 8% $\alpha$ -Al <sub>2</sub> O <sub>3</sub>	29	0.2		1.0	2:8	1.2	1200	500	7	ABO
. 20	38	0.26	0.044		8:2	0.304	1100	25	10	ABO, $30\% \alpha$ -Al <sub>2</sub> O <sub>3</sub>
54 1.0 1.0 1.2:1 2.0 1200 1200 -, 44h ABO	50	3.66	0.50		5:1	4.16	1200	800	100	ABO, 8% $\alpha$ -Al <sub>2</sub> O <sub>3</sub>
	54	1.0		1.0	1.2:1	2.0	1200	1200	-, 44h	ABO

Table 3.4: List of solid-state reaction ABO synthesis experiments.

#### Sol-gel derived precursors

Precursors for synthesis experiments of  $Al_4B_2O_9$  were prepared according to Mazza et al. (1992) and Fischer et al. (2008b).  $Al(NO_3)_3 \cdot 9 H_2O$  and  $H_3BO_3$  with 10 wt.-% of glycerol were heated to 80°C while continuously stirring with a magnetic stirrer. After homogenization, the solution was further heated to 110°C, when denitrification started. As release of NO<sub>x</sub> gasses stopped, the foamy material was dried at 200°C, then crushed and heated at 300°C for 2h to remove remaining organics. Precursors of samples #51 and #53 (Table 3.5) were heated at 950°C and 1100°C, yielding  $Al_4B_2O_9$  and ABO powder samples, respectively. The precursor of sample #52 was heated in an Anton-Paar furnace with 10°C/min and in-situ measured with powder X-ray diffraction (Figure 3.22). Each measurement took 40 min. Results are in agreement with those of previous findings, however, when 1100°C were reached,  $Al_4B_2O_9$  and ABO were coexisting for a short time (< 2h), which is also reported by Duoy (2005). After heating for 2h at 1100°C, the remaining  $Al_4B_2O_9$  phase transformed completely to ABO.

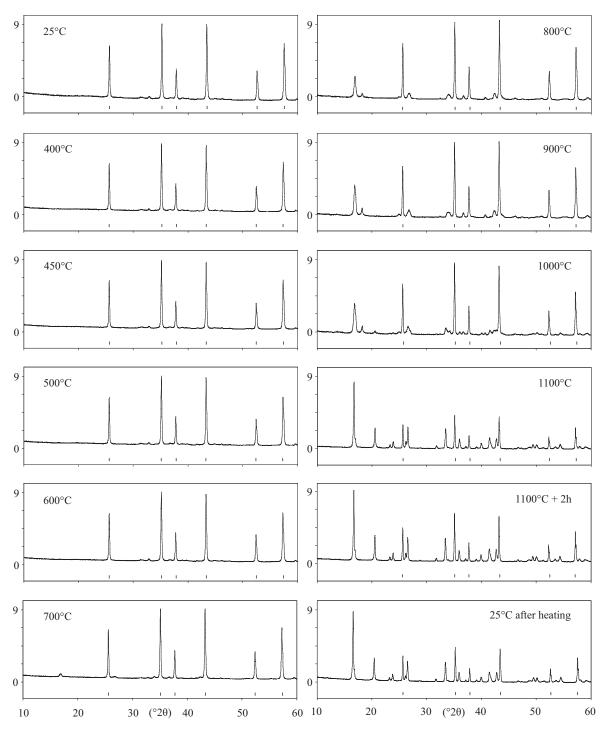
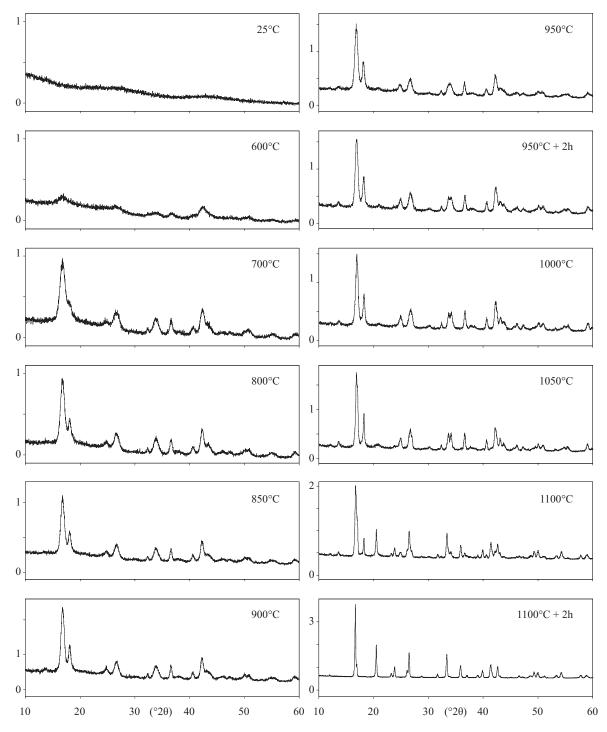


Figure 3.21: In-situ powder XRD patterns of the solid-state reaction experiment #38 starting from  $8\alpha$ -Al<sub>2</sub>O<sub>3</sub> and 2B<sub>2</sub>O<sub>3</sub>. At 25°C, only corundum can be identified (*hkl*-ticks). At 700°C, Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> starts to crystallize (confirmed by the small peak at ca. 16° 20) and the phase is stable up to 1000°C. At 1100°C, Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> is completely transformed to ABO. After 2h at 1100°, ABO peaks have slightly smaller peak width (FWHM of the (021) reflection at ca. 16.5° 20 0.139(2)° 20 instead of 0.146(2)° 20), indicating increasing size of coherently scattering domains ("crystallite size"). However, this may be an effect of sample sintering. After cooling to 25°C, ABO and corundum remain as stable phases. Intensity scale is counts  $\cdot 10^3$ .



**Figure 3.22:** In-situ powder XRD patterns of the sol-gel precursor #52 consisting of Al(NO<sub>3</sub>)<sub>3</sub> · 9 H<sub>2</sub>O and 2H<sub>3</sub>BO<sub>3</sub>. Up to ca. 600°C, the precursor material is amorphous. At 700°C, Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> starts to crystallize and crystal growth continues to about 1050°C (FWHM of the (202) reflection at ca. 16.6° 20 decreases from 0.74° 20 at 850°C to 0.28° 20 at 1050°C). At 1050, ABO formation starts. At 1100°C, both phases are coexisting, but after 2h at 1100°C, the Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> phase has completely transformed to ABO, which is the only phase present after cooling to room temperature. Intensity scale is counts  $\cdot 10^3$ .

#	$\begin{array}{c} \operatorname{Al}(\operatorname{NO}_3)_3 \cdot 9\operatorname{H}_2\operatorname{O} \\ (g) \end{array}$	H <sub>3</sub> BO <sub>3</sub> (g)	A:B (mol)	Wt. (g)	$\begin{array}{c} T_{\max} \\ (^{\circ}\mathrm{C}) \end{array}$	$\begin{array}{c} T_{\mathrm{end}} \\ (^{\circ}\mathrm{C}) \end{array}$	∇ (°C/h)	Phases (XRD)
51a	9.238	0.762	2:1	10.0	950	950	0; 15h	$Al_4B_2O_9$
51b	9.238	0.762	2:1	10.0	950	950	0; 10h	ABO
52	7.521	2.479	1:2	10.0	25	1100	10	$Al_4B_2O_9$ , ABO, A
53a	15.042	4.958	1:2	20.0	950	950	0; 5h	$Al_4B_2O_9$
53b	15.042	4.958	1:2	20.0	1100	1100	0; 5h	ABO

Table 3.5: List of sol-gel derived precursor ABO synthesis experiments.

 $\mathrm{A}=\alpha\text{-}\mathrm{Al}_2\mathrm{O}_3$ 

 $Al_4B_2O_9$  according to the structure of Fischer et al. (2008b)

# 3.2 Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> Compounds with Mullite-Type Structure

#### 3.2.1 Overview

Three different methods for the production of compounds with mullite-type crystal structure in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$  are reported in literature: (1) Dietzel and Scholze (1955) synthesized compounds identified as "mullites" with powder X-ray diffraction from a solidstate reaction route. They used crystalline  $B_2O_3$ ,  $\alpha$ - $Al_2O_3$  and quartz as starting materials. The charge was heated in molybdenum boats under reducing atmosphere at temperatures between 1385°C and 1775°C. (2) From a sol-gel precursor consisting of  $Al(NO_3)_3 \cdot 9 H_2O$ ,  $H_3BO_3$  and tetra ethyl orthosilicate (TEOS,  $C_8H_{20}O_4Si$ ), Griesser et al. (2008) were able to synthesize mullite-type structures. (3) Werding and Schreyer (1992) obtained a compound with an assigned stoichiometry of  $Al_8Si_2B_2O_{19}$  at temperatures between 800°C and 830°C and pressures from 1 kbar to 4 kbar. They noticed the similarity of powder X-ray diffraction lines to those of ABO (Sokolova et al. 1978). The exact structure of the compound has not been solved yet.

## 3.2.2 Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> Synthesis Experiments

#### Solid-state reaction

Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> samples were prepared from  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, previously heated silica gel and B<sub>2</sub>O<sub>3</sub>. In a first attempt, the stoichiometry of boralsilite Al<sub>16</sub>Si<sub>2</sub>B<sub>6</sub>O<sub>37</sub> (sample #25) (Grew et al. 2008) and a composition of Al<sub>2</sub>O<sub>3</sub>:B<sub>2</sub>O<sub>3</sub>:SiO<sub>2</sub> of 8:3:4 (sample #26) were chosen. Subsequently, a hypothetical, although not charge neutral exchange of Si<sup>4+</sup>  $\rightarrow$  Al<sup>3+</sup> was tested with starting mixtures representing stoichiometries of Al<sub>4.5</sub>Si<sub>0.5</sub>BO<sub>9.25</sub> (#57), Al<sub>4</sub>SiBO<sub>9.5</sub> (#58) and Al<sub>3.5</sub>Si<sub>1.5</sub>BO<sub>9.75</sub> (#59). Due to overlapping peaks of mullite, sillimanite, cristobalite and ABO, phase identification with powder X-ray diffraction was ambiguous and results may not be accurate. Starting material compositions and results from powder X-ray diffraction patterns are listed in Table 3.6.

#	$\begin{array}{c} \mathrm{Al_2O_3} \\ \mathrm{(g)} \end{array}$	$\substack{\mathbf{B}_2\mathbf{O}_3\\(\mathbf{g})}$		Theoretical stoichiometry	A:B:Si (mol)	Wt. (g)	Temp. (°C)	Time	Phases (XRD)
25	1.425	0.364	0.210	$Al_{16}Si_2B_6O_{37}$	8:3:2	2.0	1300	2h	ABO*
26	0.712	0.182	0.210	$Al_{16}Si_4B_6O_{41}$	8:3:4	1.1	1300	2h	ABO, A
57	0.714	0.192	0.090	$\mathrm{Al}_{4.5}\mathrm{Si}_{0.5}\mathrm{BO}_{9.25}$	4.5:1:0.5	1.0	1200	12d	ABO, A
58	0.626	0.189	0.184	$Al_4SiBO_{9.5}$	4:1:1	1.0	1200	9d	ABO, A, C
59	0.810	0.280	0.409	$Al_{3.5}Si_{1.5}BO_{9.75}$	$3.5{:}1.5{:}1$	1.5	1200	4.5d	ABO, A, C

Table 3.6: List of Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> solid-state reaction synthesis experiments.

 $\mathrm{A}=\alpha\text{-}\mathrm{Al}_2\mathrm{O}_3,\,\mathrm{C}=\mathrm{cristobalite}\;(\mathrm{SiO}_2)$ 

\*<br/>one peak at 27.908° $2\theta$  fits sillimanite

#### Sol-gel derived precursors

The method used by Griesser et al. (2008) was slightly modified to prepare twelve sol-gel precursors. Unlike as in Griesser et al. (2008), 5g of glycerol were added to each precursor resulting in a total weight of 23g per precursor. Glycerol forms ester complexes with boric acid at room temperature and thus prevents volatilization of  $H_3BO_3$  at higher temperatures (Mazza et al. 1992). Furthermore, handling of the precursors is easier with glycerol. Al(NO<sub>3</sub>)<sub>3</sub> · 9 H<sub>2</sub>O,  $H_3BO_3$ , TEOS ( $C_8H_{20}O_4Si$ ) and glycerol were diluted with ethanol p.a. and stirred for ca. 2h at 50°C. With further heating, the solutions turned brown and at ca. 75°C, brownish NO<sub>x</sub> gases were released. The resulting foamy gel was dried at 150°C in a furnace for 5h. After grinding the material and subsequent heat treatment at 400°C, a brownish powder was obtained. This precursor material was heat treated at 950°C and 1250°C for 5h.

#	Theoretical stoichiometry	$\begin{array}{c} \mathrm{Al}_2\mathrm{O}_3\\ (\mathrm{mol}\text{-}\%) \end{array}$	$\substack{\mathrm{B}_2\mathrm{O}_3\\(\mathrm{mol}\text{-}\%)}$	$\begin{array}{c} \mathrm{SiO}_2 \\ \mathrm{(mol-\%)} \end{array}$	Phases (XRD)
А	$Al_5BO_9$	83.33	16.67	0.00	Mullite-type*†
В		76.55	13.35	10.00	Mullite-type*†
С		70.00	10.00	20.00	Mullite-type*†
D	Boromullite $Al_9Si_2BO_{19}$	64.29	7.14	28.57	Mullite-type*†
Е		56.57	3.35	40.00	Mullite-type*†
$\mathbf{F}$	Sillimanite $Al_2SiO_5$	50.00	0.00	50.00	$\gamma$ -Al <sub>2</sub> O <sub>3</sub> , Quartz
G		53.37	5.59	40.00	Mullite-type*†
Η	close to $Al_8Si_2B_2O_{19}$	56.61	13.38	30.00	$Mullite-type*\dagger$
Ι		60.00	20.00	20.00	Mullite-type <sup>*</sup> , Quartz
J	Boralsilite $Al_{16}B_6Si_2O_{37}$	61.54	23.08	15.38	Mullite-type, Quartz?
Κ		63.31	26.69	10.00	Mullite-type*, Quartz?
L	$Al_4B_2O_9$	66.67	33.33	0.00	$Al_4B_2O_9$ (Fischer et al. 2008b)

Table 3.7: List of Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> sol-gel derived precursor synthesis experiments at 950°C.

\* from LeBail refinements starting with a = 7.6 Å, b = 7.5 Å, c = 2.8 Å in space group *Pbam* † very broad bump at 20° 20 may fit  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>

Tables 3.7 and 3.8 show a list of the precursor compositions and possible phase identification according to matching reflection positions of powder XRD data.

Samples prepared in the  $Al_2O_3$ - $B_2O_3$ -SiO<sub>2</sub> system were not easy to identify by means of

#	Theoretical stoichiometry	$\begin{array}{c} \mathrm{Al}_2\mathrm{O}_3\\ (\mathrm{mol}\text{-}\%) \end{array}$	$\substack{\mathrm{B}_2\mathrm{O}_3\\(\mathrm{mol}\text{-}\%)}$		Phases (XRD)
А	$Al_5BO_9$	83.33	16.67	0.00	ABO-type $(Cmc2_1)^*$ , A
В	· ·	76.55	13.35	10.00	ABO-type $(Cmc2_1)^*$ , A
С		70.00	10.00	20.00	ABO-type $(Cmc2_1)^*$ , A
D	Boromullite $Al_9Si_2BO_{19}$	64.29	7.14	28.57	ABO-type $(Cmc2_1)^*$ , A
Е		56.57	3.35	40.00	ABO-type $(Cmc2_1)^*$ , A
$\mathbf{F}$	Sillimanite $Al_2SiO_5$	50.00	0.00	50.00	Mullite-type $(Pbam)$ <sup>†</sup> , C, A
G		53.37	5.59	40.00	ABO-type $(Cmc2_1)^*$ , A
Η	close to $Al_8Si_2B_2O_{19}$	56.61	13.38	30.00	ABO-type $(Cmc2_1)^*$
Ι		60.00	20.00	20.00	ABO-type $(Cmc2_1)^*$
J	Boralsilite $Al_{16}B_6Si_2O_{37}$	61.54	23.08	15.38	ABO-type $(Cmc2_1)^*$
Κ		63.31	26.69	10.00	ABO-type $(Cmc2_1)^*$
L	$\rm Al_4B_2O_9$	66.67	33.33	0.00	ABO-type $(Cmc2_1)^*$

**Table 3.8:** List of  $Al_2O_3$ - $B_2O_3$ - $SiO_2$  sol-gel derived precursor synthesis experiments at 1250°C.

 $A = \alpha - Al_2O_3$ ,  $C = cristobalite (SiO_2)$ 

\* from LeBail refinements starting with a = 5.67 Å, b = 15.01 Å, c = 7.68 Å in space group  $Cmc2_1$ 

 $\dagger$  from LeBail refinements starting with a = 7.6 Å, b = 7.5 Å, c = 2.8 Å in space group Pbam

powder X-ray diffraction and Al, B and Si sites in mullite-type structures cannot be distinguished. Nevertheless, reasonable LeBail fits with a mullite-like unit cell ( $a \approx 7.6$  Å,  $b \approx 7.5$ Å,  $c \approx 2.8$  Å in space group *Pbam*) were achieved in powder X-ray diffraction patterns of all 950°C samples except F and L. Sample F was almost amorphous, however, peaks may fit nanocrystalline  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (Chakravorty 1993, Chakraborty and Das 2003, Gutierrez et al. 2002, Duoy 2005, Cividanes et al. 2010). According to Chakravorty (1993) and Chakraborty and Das (2003), this phase can also be interpreted as an Al-Si spinel phase. The powder pattern of sample L was in agreement with the structure of Fischer et al. (2008b).

All 1250°C samples except boron-free sample F could be indexed using a unit cell with similar lattice parameters as of ABO and boromullite (Buick et al. 2008) in space group  $Cmc2_1$ . Surprisingly, no excess SiO<sub>2</sub> was found, except in sample F, which was identified as a mixture of mullite, corundum and tetragonal cristobalite. The synthesis temperature of 1250°C is below the stability range of cristobalite. However, impure tetragonal has been observed at lower temperatures (Smallman and Bishop 1999).

# 3.3 Discussion

Different types of synthesis routes yield ABO and  $Al_4B_2O_9$  and the production of those materials is easy. It is, however, rather difficult to obtain single-crystals larger than ca. 100 µm. With respect to this, only two of the synthesis experiments were successful: Unreproducible run #01, yielding ABO single-crystals with a fair amount of  $Al_2O_3$ -inclusions grown from  $K_2CO_3 + 3MoO_3$  flux and run #62-2, yielding single-phase ABO crystals from presintered ABO in excess  $B_2O_3$  as flux. It seems that  $T_{end}$  and the cooling gradient is crucial in synthesis experiments with  $K_2MoO_3O_{10}$  as flux because the compound decays to several intermediate phases (e.g.  $K_2Mo_4O_{13}$ , Leonyuk et al. 1980). The fact that ABO was presintered and the higher formation temperature in run #62 may be the reason that no excess corundum was found in this experiment. In any system different to  $Al_2O_3$ - $B_2O_3$ - $SiO_2$ , products contained secondary phases, which could not be easily separated. Crystal size was very small.

All solid-state reaction experiments were successful and similar formation temperatures of  $Al_4B_2O_9$  and ABO as known from literature were observed with in-situ high-temperature powder X-ray diffraction measurements. In stoichiometrically weighed precursor materials without excess boron, ABO crystallized together with excess  $Al_2O_3$ . Obviously, lack of boron due to  $B_2O_3$  evaporation at high temperature is the limiting factor for single-phase ABO formation from such precursors. Interestingly, experiments in sealed platinum capsules to avoid boron evaporation did not yield single-crystals suitable for XRD experiments as in Garsche et al. (1991). Nevertheless, from FWHM peak-widths of powder XRD patterns, a positive correlation of the size of coherently scattering domains ("crystallite size") and excess  $B_2O_3$  in the starting mixture could be derived. This can have two reasons: Either, excess  $B_2O_3$  lowers crystallization temperature of the phases, as shown by Duoy (2005) and Zhang et al. (2010), and thus larger crystallites grow, or excess  $B_2O_3$  aids sintering of the ABO crystals. The latter argument may be supported with the observation that sintered tablets of samples produced with excess  $B_2O_3$  were noticeably harder to grind, which is in agreement with findings of Sokolov and Gasparyan (2004).

All synthesis experiments performed with a morphous precursors obtained by the sol-gel route yielded ABO and/or  $\rm Al_4B_2O_9$ . However, X-ray powder diffraction peak widths of the  $\rm Al_4B_2O_9$  phase at 1000°C were noticeably broader than those of  $\rm Al_4B_2O_9$  produced from  $\alpha$ -Al\_2O\_3 and  $\rm B_2O_3$  (FWHM of the (202) peak ca. 0.38° 20 instead of ca. 0.28° 20 at 1000°C), indicating that a morphous sol-gel precursors may not be the optimal starting materials for the production of crystalline  $\rm Al_4B_2O_9$ . FWHM of ABO samples prepared by this route were comparable to those produced from solid-state reactions. This is not surprising as ABO is formed from the metastable intermediate  $\rm Al_4B_2O_9$  phase in both solid-state reaction and from sol-gel derived precursors according to the following equation:

$$\begin{array}{l} 5\operatorname{Al}_4\mathrm{B}_2\mathrm{O}_9 \xrightarrow{1050^\circ\mathrm{C}-1100^\circ\mathrm{C}} 4\operatorname{Al}_5\mathrm{BO}_9 + 3\operatorname{B}_2\mathrm{O}_3, \text{ or} \\ 9\operatorname{Al}_4\mathrm{B}_2\mathrm{O}_9 \xrightarrow{1050^\circ\mathrm{C}-1100^\circ\mathrm{C}} 2\operatorname{Al}_{18}\mathrm{B}_4\mathrm{O}_{33} + 5\operatorname{B}_2\mathrm{O}_3 \end{array}$$

Compared to the in-situ measured solid-state reaction synthesis experiment with  $Al_2O_3$ and  $B_2O_3$  as precursor, the crystallization temperature of  $Al_4B_2O_9$  was slightly lower and the phase coexisted with ABO at 1100°C for a short time (< 2h). This confirms the observation of Duoy (2005) that the temperature dependent stability field of  $Al_4B_2O_9$  seems to be slightly extended when amorphous precursors are used.

Most samples in the system  $Al_2O_3-B_2O_3-SiO_2$  yielded phases with only slightly different unit cell parameters than those of already known mullite-type phases in the systems  $Al_2O_3$ -

 $\rm B_2O_3$  and  $\rm Al_2O_3\text{-}SiO_2.$  From preliminary sample investigation with powder X-ray diffraction, no conclusions about Al, Si and/or B incorporation into the structures and hypothetical solid solutions can be made.

# 4 Crystal-Chemistry of Mullite-Type Aluminoborates Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> and Al<sub>5</sub>BO<sub>9</sub>: A stoichiometry puzzle

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# 4.1 Abstract

Orthorhombic  $Al_2O_3$ -rich aluminoborate is an important ceramic material for which two slightly different compositions have been assumed:  $Al_5BO_9$  ( $5Al_2O_3:B_2O_3$ ) and  $Al_{18}B_4O_{33}$ ( $9Al_2O_3:2B_2O_3$ ). The formula  $Al_{18}B_4O_{33}$  (=  $Al_{4.91}B_{1.09}O_9$ ) was derived from results of chemical analyses when crystal structure data were not yet available. Subsequent structural investigations indicated  $Al_5BO_9$  composition. Nevertheless,  $Al_{18}B_4O_{33}$  was still accepted as the correct stoichiometry assuming that additional B replaces 9% Al.

Powder samples of both compositions and ones with excess boron were prepared by solidstate reactions between  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub>/H<sub>3</sub>BO<sub>3</sub> at temperatures above 1100°C and singlecrystals were grown from flux at 1100°C and 1550°C. Products were investigated by singlecrystal and powder XRD, <sup>11</sup>B and <sup>27</sup>Al solid-state MAS-NMR, Raman and FTIR spectroscopy as well as Laser-ablation ICP-MS. No indication of the predicted 9% B  $\rightarrow$  Al substitution was found. LA ICP-MS indicated 12.36(27) wt.-% B<sub>2</sub>O<sub>3</sub> corresponding to Al<sub>4.97</sub>B<sub>1.03</sub>O<sub>9</sub>. Hence, the suggested Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> stoichiometry can be excluded for all synthesized samples. A very low amount of Al vacancies at a five-fold coordinated site are likely, charge balanced by an additional nearby three-fold coordinated B site. All evidence indicates that the title compound should be reported as Al<sub>5-x</sub>B<sub>1+x</sub>O<sub>9</sub> with x < 0.038(6), which is close to Al<sub>5</sub>BO<sub>9</sub>.

# 4.2 Introduction

Mullite used as ceramic raw material is an important and versatile compound in the system  $SiO_2-Al_2O_3$ . Its properties include high thermal stability, very low thermal expansion, low heat conductivity, high creep and corrosion resistance and high stiffness (Schneider and Komarneni 2005). The composition of mullite is rather variable:  $Al_2(Al_{2+2x}Si_{2-2x})O_{10-x}, 0.18 \le x \le 0.88$ . However, the most common compositions are 3:2 mullite,  $3Al_2O_3:2SiO_2$ , x = 0.25 and 2:1 mullite,  $2Al_2O_3:SiO_2$ , x = 0.4.

In the system  $SiO_2-Al_2O_3-B_2O_3$ , Werding and Schreyer (1984) and Werding and Schreyer (1996) introduced the term "boron-mullite" for compounds within a compositional range between 3:2 and 2:1 mullite and two silicon-free aluminoborate members (Figure 4.1),  $Al_5BO_9$  ( $5Al_2O_3:B_2O_3$ ) with a mullite-type structure (Sokolova et al. 1978) and  $AlBO_3$  ( $Al_2O_3:B_2O_3$ ) of calcite structure-type, stable under hydrothermal high-pressure conditions (Capponi et al. 1972, Vegas et al. 1977). More recently, the term mullite-type boron compound has been defined (Fischer and Schneider 2008a) for corresponding structures in the binary system  $Al_2O_3-B_2O_3$  whereas the name boromullite (Buick et al. 2008) is reserved for a mineral with sillimanite- and  $Al_5BO_9$ -like modules. In addition,  $Al_3BO_6$  (Capponi et al. 1972) with norbergite structure-type (White et al. 1965) was synthesized above 25 kbar and 800°C and rhombohedral  $Al_4B_6O_{15}$  with a microporous framework consisting of  $AlO_6$  octahedra and  $BO_3$  units has been produced from  $AlCl_3$  and  $H_3BO_3$  at 350°C (Jing et al. 2004).

In the  $Al_2O_3$ -rich part of the phase diagram,  $Al_{18}B_4O_{33}$  (9  $Al_2O_3:2B_2O_3$ ) is composition-

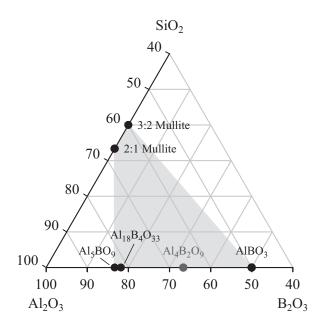


Figure 4.1:  $Al_2O_3$ -rich part of the ternary system  $SiO_2$ - $Al_2O_3$ - $B_2O_3$ . The boron-mullite stability field according to Werding and Schreyer (1984) and Werding and Schreyer (1996) is represented by the gray area. All units in mol-%.

ally very close to  $Al_5BO_9$  ( $5Al_2O_3:B_2O_3$ ). This material gained industrial interest because of mullite-like properties. Due to its low-cost production ( $\frac{1}{10}$  to  $\frac{1}{20}$  of the cost of SiC, Kim et al. 2004 and Song et al. 2007), its easy fabrication in large quantities (Zhang et al. 2006, Wang et al. 2008, Wada et al. 1991, Sokolov and Gasparyan 2004, Peng et al. 2006, Elssfah et al. 2007), its high strength (Ray 1992, Xinyong et al. 2007, Gatta et al. 2010), and its low thermal expansion and conductivity (Wada et al. 1993; 1994) the compound is used as reinforcer in metal matrix composites (Wang et al. 2008, Peng et al. 2006). Further applications include reinforcement in fire insulations for ships, construction components in nuclear reactors because of neutron absorbing capabilities, and in refractory linings due to high resistance against boron-rich glass melts (Garsche et al. 1991 and refs. therein). Recently, nanotubes have been synthesized (Ma et al. 2002, Li and Chang 2006) and aluminoborate fibers have been successfully coated by boron nitride to lower interface reactions and enhance strength between matrix and aluminoborate nanowires or whiskers (Song et al. 2007, Zhang et al. 2005, Fei et al. 2010).

Although the aluminoborate discussed above is of high importance, results of crystalchemical investigations have remained inconsistent. This mullite-type boron compound is reported with two slightly different stoichiometries,  $Al_{18}B_4O_{33} = Al_{4.91}B_{1.09}O_9 =$  $9Al_2O_3:2B_2O_3$  and  $Al_5BO_9 = 5Al_2O_3:B_2O_3$  (Table 4.1), both crystallizing in the same mullite-type structure. There are no studies to date clarifying whether both similar compounds exist, probably because applications of this material were considered to be more important than its exact chemical characterization. Persistent assumption of  $Al_{18}B_4O_{33}$  stoichiometry may thus have complied with tradition as this compound was originally defined on early analytical data without considering the crystal structure.

	Table 4.1: 1	neoretical compo	SITION OF $AI_5 DO_9$	and $AI_{18}D_4O_{33}$ .	
	$\mathrm{Al}_2\mathrm{O}_3{:}\mathrm{B}_2\mathrm{O}_3$	$\rm Al_2O_3~(wt\%)$	$\mathrm{B}_2\mathrm{O}_3$ (wt%)	$Al_2O_3 \ (mol-\%)$	$\mathrm{B}_2\mathrm{O}_3~(\mathrm{mol}\text{-}\%)$
$Al_5BO_9$	5:1	87.99	12.01	83.33	16.66
$\mathrm{Al}_{18}\mathrm{B}_{4}\mathrm{O}_{33}$	9:2	86.83	13.17	81.81	18.18

Table 4.1: Theoretical composition of  $Al_5BO_9$  and  $Al_{18}B_4O_{33}$ .

 $Al_{18}B_4O_{33}$  and  $Al_5BO_9$  contain only trivalent cations, though of different radius.  $Al^{3+}$  and  $B^{3+}$  may both occur in four-fold coordination. Therefore, both Al and B may exist as a solid solution at a tetrahedrally coordinated site of a mullite-type structure. On the other hand, Al-O distances in an AlO<sub>4</sub> tetrahedron are ca. 1.75 Å (Kunz and Armbruster 1990) whereas B-O distances are ca. 1.476 Å (Hawthorne et al. 1996). This bond length difference of ca. 15% seems to contradict an extensive solid solution. Furthermore, it could be expected that the degree of this substitution is temperature dependent. The aim of this study is an investigation of aluminoborate samples, produced above 1100°C by different synthesis routes with a compositional range allowing formation of  $Al_{18}B_4O_{33}$  and/or  $Al_5BO_9$ . Using single-crystal and powder X-ray diffraction, solid-state <sup>11</sup>B and <sup>27</sup>Al MAS-NMR, FTIR, Raman spectroscopy and Laser-ablation ICP-MS compositional measurements, the composition of this important compound will be clarified.

# 4.3 Historical Background

In the late 19<sup>th</sup> century, Mallard (1887) studied needle-like, facetted crystals with orthorhombic symmetry of supposed  $3 \text{Al}_2\text{O}_3$ : B<sub>2</sub>O<sub>3</sub> composition, synthesized by Ebelmen (1851). Until then, the only known aluminoborate was the mineral jeremejevite,  $Al_6(BO_3)_5(F,OH)_3$  (Mallard 1887, Websky 1883). Benner and Baumann (1938) patented synthesis of an aluminoborate starting from molten Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub>. These acicular, orthorhombic crystals were assumed to have 3 Al<sub>2</sub>O<sub>3</sub>:B<sub>2</sub>O<sub>3</sub> composition. Mullite-like physical properties qualified this aluminoborate as a new ceramic raw material. Cell dimensions (Table 4.2) derived from powder X-ray diffraction were presented by Baumann and Moore (1942). From the density of 2.93 g/cm<sup>3</sup>, they concluded that the correct formula of the crystals is  $9 \operatorname{Al}_2 \operatorname{O}_3 : 2 \operatorname{B}_2 \operatorname{O}_3 (\operatorname{Al}_{18} \operatorname{B}_4 \operatorname{O}_{33})$ , with a unit cell content of Z = 1.1. Dietzel and Scholze (1955) studied glasses in the system SiO<sub>2</sub>-Al<sub>2</sub>O<sub>3</sub>- $B_2O_3$ . From analyses of crystalline by-products they proposed a solid solution between 3:2 mullite and Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub>. Subsequent experiments in the system Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub> (Scholze 1956) indicated a new mullite-type phase of  $2 \operatorname{Al}_2 O_3 : B_2 O_3$  (Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>) composition (Mazza et al. 1992, Fischer et al. 2008b) obtained by heating Al<sub>2</sub>O<sub>3</sub> in a B<sub>2</sub>O<sub>3</sub> flux at 1000°C. According to the  $Al_2O_3$ - $B_2O_3$  phase diagram,  $Al_4B_2O_9$  transforms to  $Al_{18}B_4O_{33}$  at 1035°C (Scholze 1956, Gielisse and Forster 1962). Sokolova et al. (1978) grew aluminoborate single-crystals by cooling an Al<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-Y<sub>2</sub>O<sub>3</sub>-K<sub>2</sub>O-MoO<sub>3</sub> melt from 1150°C to 950°C. Cell dimension and

orthorhombic symmetry were in accordance with previous findings for  $Al_{18}B_4O_{33}$ . However, single-crystal X-ray structure refinement yielded  $Al_5BO_9$  composition with Z = 4 instead of  $Al_{18}B_4O_{33}$  (Table 4.2). The structure was described with one  $AlO_6$  octahedron, three  $AlO_4$ tetrahedra and one planar  $BO_3$  group. Results of unspecified chemical analyses yielded 14 wt.-%  $B_2O_3$  (no esd's given), which is closer to  $Al_{18}B_4O_{33}$  than to  $Al_5BO_9$  (Table 4.1). Singlecrystals grown in the system CaO- $Al_2O_3$ - $B_2O_3$  were also studied by Ihara et al. (1980). Results of gravimetric analyses agreed with  $Al_{18}B_4O_{33}$ , Z = 1.09 (Table 4.2), but the structural data were in agreement with those of Sokolova et al. (1978). They concluded that two of the Al-tetrahedra specified by Sokolova et al. (1978) are more precisely characterized as five-fold coordinated  $Al^{(V)}$  polyhedra. The discrepancy between the crystallographically derived composition of  $Al_5BO_9$  and the chemical composition pointing to  $Al_{18}B_4O_{33}$  was explained with a disordered structure in which  $\frac{1}{11} \approx 9\%$  of  $Al^{(IV)}$  is substituted by tetrahedrally coordinated  $B^{(IV)}$ .

Garsche et al. (1991) produced single-crystals by fusion of  $Al_2O_3$  and  $B_2O_3$  in sealed platinum capsules at 1500°C. X-ray powder and single-crystal diffraction data (Table 4.2) are in agreement with findings of Sokolova et al. (1978) and Ihara et al. (1980). Without further investigations, the compound was reported as  $Al_{18}B_4O_{33}$ , based on the tetrahedral B substitution proposed by Ihara et al. (1980).

The Al-coordination in "Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub>" has also been investigated by <sup>27</sup>Al MAS-NMR (Massiot et al. 1995, Gan et al. 2002, Hung et al. 2006) on samples produced according to Garsche et al. (1991). One Al<sup>(IV)</sup>, two distinct Al<sup>(V)</sup> and one Al<sup>(VI)</sup> site were assigned with multiplicities of 1:1:1:2 in accordance with previous crystal structure data (Sokolova et al. 1978, Garsche et al. 1991, Ihara et al. 1980).

Mazza et al. (1992) synthesized crystalline products from an amorphous precursor of  $H_3BO_3$ and  $Al(NO_3)_3 \cdot 9 H_2O$  and proposed a solid solution  $Al_{6-x}B_xO_9$  with  $1 \le x \le 3$ . The structures were solved using a pseudotetragonal mullite model in space group *Pbam* (Table 4.2). FTIR spectra of a member with x = 1 showed no evidence of  $B^{(IV)}$ , which was later also confirmed by <sup>11</sup>B MAS-NMR (MacKenzie et al. 2007). Mazza et al. (1992) consider  $Al_5BO_9$  as being stable in the temperature regime between 900°C and 1000°C, leading to  $Al_{18}B_4O_{33}$  with space group  $Cmc2_1$  upon heating at higher temperature.

# 4.4 Experimental

## 4.4.1 Sample Preparation

Using stoichiometric mixtures and mixtures containing excess boron to exclude  $B_2O_3$  as the limiting factor for  $Al_{18}B_4O_{33}$  formation (due to boron evaporation at high temperature), starting materials were prepared in order to allow formation of  $Al_{18}B_4O_{33}$  and/or  $Al_5BO_9$ . All syntheses were performed in lid-covered platinum crucibles in air. Three different synthesis routes were followed: (1) Powder samples *b*, *c* and *e* were prepared by solid-state reactions

	I	able 4.2: Theoret	<b>Table 4.2:</b> Theoretical composition of $Al_5BO_9$ and $Al_{18}B_4O_{33}$ .	$O_9$ and $Al_{18}B_4O_{33}$	÷		
	Baumann and	Dietzel and	Scholze $(1956)$	Sokolova et	Ihara et	Garsche et	Mazza et
	al. $(1942)$	Scholze $(1955)$		al. (1978)	al. (1980)	al. (1991)	al. $(1992)$
a-axis	15.0		5.68(5)	5.6673(7)	5.682(13)	7.6942(1)	7.621
b-axis	7.5		14.98(10)	15.011(2)	14.973(34)	15.0110(2)	7.621
c-axis	5.67		7.68(5)	7.693(1)	7.692(17)	5.6689(1)	2.833
Volume				654.4(2)		654.74	
23	1.1		1.09	4	1.09	1.09	1.0
Space group	Orthorhombic		$Cmc2_1, C2cm, Cmcm$	$Cmc2_1$	$Cmc2_1$	$A2_1 am$	Pbam
Density $(g/cm^3)$	2.93		2.94(1)	2.96	2.93(1)		
Assigned formula	$\mathrm{Al}_{18}\mathrm{B}_4\mathrm{O}_{33}$	$\mathrm{Al}_{18}\mathrm{B}_4\mathrm{O}_{33}$	$\mathrm{Al_{18}B_4O_{33}}$	$\mathrm{Al}_5\mathrm{BO}_9$	$\mathrm{Al}_{18}\mathrm{B}_4\mathrm{O}_{33}$	$\mathrm{Al}_{18}\mathrm{B}_4\mathrm{O}_{33}$	$\mathrm{Al}_5\mathrm{BO}_9$
$B_2O_3 (wt\%)$		13.2	13.3	14	13.1		
$Al_2O_3 (wt\%)$		86.6	86.7	86	86.9		
Stoichiometry		$\mathrm{Al}_{18}\mathrm{B}_4\mathrm{O}_{33}$	$Al_{17.96}B_{4.04}O_{33}$	$Al_{17.77}B_{4.23}O_{33}$	$Al_{18.02}B_{3.98}O_{33}\\$		
from wt%		$Al_{4.91}B_{1.09}O_9{}^a$	$\rm Al_{4.90}B_{1.10}O_9{}^a$	$Al_{4.85}B_{1.15}O_9{}^{\rm a}$	$Al_{4.91}B_{1.09}O_9{}^a$		
<sup>a</sup> normalized to 9 oxygen							

of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> with B<sub>2</sub>O<sub>3</sub> or H<sub>3</sub>BO<sub>3</sub>. The powders were thoroughly mixed, pressed to pellets and subsequently heated at 1100°C or 1200°C. (2) Powder sample *d* was produced from an amorphous precursor prepared with Al(NO<sub>3</sub>)<sub>3</sub> · 9 H<sub>2</sub>O and H<sub>3</sub>BO<sub>3</sub>, according to Fischer et al. (2008b). The denitrified raw material was heated at 1100°C. (3) Single-crystals were grown from a mixture of 9 Al<sub>2</sub>O<sub>3</sub>:2 B<sub>2</sub>O<sub>3</sub> in a K<sub>2</sub>CO<sub>3</sub> + 3 MoO<sub>3</sub> flux (borate/flux ratio ca. 1:9) by cooling the melt at 10°C/h from 1100°C to 600°C (sample *a*) and from a pressed pellet (previously heated at 1200°C) consisting of Al<sub>2</sub>O<sub>3</sub>:3.3H<sub>3</sub>BO<sub>3</sub> in B<sub>2</sub>O<sub>3</sub> flux by slow cooling (15°C/h) from 1550°C to 1250°C (sample *f*). Fluxes were dissolved in hot deionized H<sub>2</sub>O. Prior to further investigations, samples were washed in warm deionized H<sub>2</sub>O to eliminate any remaining H<sub>3</sub>BO<sub>3</sub> or B<sub>2</sub>O<sub>3</sub>. Details on samples and synthesis conditions are given in Table 4.3.

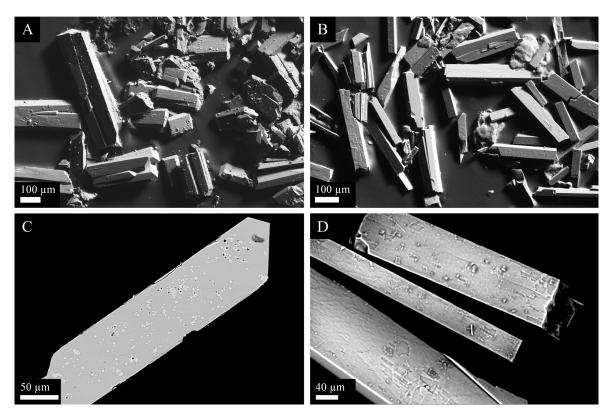
Both  $K_2CO_3 + 3MoO_3$  and  $H_3BO_3$  flux methods yielded elongated single-crystals with different habits. Radially grown crystals from sample *a* were of elongate prismatic shape with a rather smooth surface but with  $Al_2O_3$  inclusions. No inclusions were found in crystals from sample *f*, but they had a rather rough and flaky surface (Figure 4.2). A sample of  $Al_4B_2O_9$ (Mazza et al. 1992, Fischer et al. 2008b) was prepared according to Fischer et al. (2008b) and used as reference compound for tetrahedrally coordinated  $B^{(IV)}$  in FTIR, Raman and <sup>11</sup>B MAS-NMR spectroscopy.

Sample	$\mathrm{Al}_2\mathrm{O}_3{:}\mathrm{B}_2\mathrm{O}_3$	Starting materials	Temperature
a	9:2	$Al_2O_3$ , $B_2O_3$ in $K_2CO_3 + 3 MoO_3$ flux	$1100^{\circ}\mathrm{C} \rightarrow 600^{\circ}\mathrm{C}$
b	9:2	$Al_2O_3, B_2O_3$	$1200^{\circ}C$ for $10h$
c	5:1	$Al_2O_3, B_2O_3$	$1200^{\circ}C$ for $10h$
d	1:2	$Al(NO_3)_3 \cdot 9H_2O, H_3BO_3$	$1100^{\circ}C$ for 5h
e	1.2:2	Al <sub>2</sub> O <sub>3</sub> , H <sub>3</sub> BO <sub>3</sub>	$1200^{\circ}C$ for 44h
f	1:3.3	$Al_2O_3$ , $H_3BO_3$ in $B_2O_3$ flux	$1550^{\circ}\mathrm{C} \rightarrow 1250^{\circ}\mathrm{C}$
Sample	Product	Analyzed with	
a	Single-crystals	Single-crystal and powder XRD, MAS-NM	R, LA ICP-MS
b	Powder	Powder XRD, MAS-NMR	
с	Powder	Powder XRD, MAS-NMR	
d	Powder	Powder XRD, MAS-NMR	
e	Powder	Powder XRD, MAS-NMR, Raman, FTIR	
f	Single-crystals	Single-crystal and powder XRD, MAS-NM	R, LA ICP-MS

 Table 4.3: Starting materials, synthesis conditions and analytical methods used for sample characterization.

# 4.4.2 X-ray Diffraction

A full intensity dataset was measured with an Enraf Nonius CAD4 diffractometer on a singlecrystal of sample *a* (Table 4.4). Lattice parameters were determined from 24 reflections centered at four high-angular settings with  $\pm 37.9^{\circ} < \theta < \pm 44.4^{\circ}$  in order to reduce crystal and beam alignment errors. After correcting the data for Lorentz-polarization and absorption ef-



**Figure 4.2:** SEM image of samples a and f: (A) secondary electron (SE) picture of sample a showing the prismatic habit of the crystals. (B) SE picture showing the elongated prismatic crystals from sample f. (C) Backscattered electron (BSE) image of Al<sub>2</sub>O<sub>3</sub> inclusions as bright dots in a polished crystal of sample a. (D) BSE image showing the rough and flaky surface of crystals from sample f.

fects with WinGX v. 1.80.05 software package (Farrugia 1999), the structure was solved by direct methods and refined with Bruker ShelXTL v. 6.10 (Bruker-AXS 2000) using neutral atomic scattering factors in space group  $Cmc2_1$  (no. 36) with Z = 4. The Flack parameter (Flack 1983) of 0.43(10) indicated 1:1 merohedral twinning of the selected crystal.

Due to its significantly smaller size, a full intensity dataset of a crystal from sample f was collected with a Bruker Smart Apex2 CCD diffractometer (Table 4.4). Subsequently 15 reflections centered at four angular settings with  $\pm 11.0^{\circ} < \theta < \pm 30.0^{\circ}$  were measured with the CAD4 for direct comparison of lattice parameters with those of sample a. CCD data were integrated and empirically absorption-corrected using Apex2 v. 2009-11.0 software package (Bruker-AXS 2009). The structure was refined with Bruker ShelXTL v. 6.10 (Bruker-AXS 2000) using the structural model obtained from sample *a*. A Flack parameter of 0.60(11) indicated 1:1 merohedral twinning of the crystal. Single-crystal data collection parameters are reported in Table 4.4.

XRD powder patterns were measured with a PANalytical X'Pert MPD diffractometer equipped with a Cu X-ray source (40 kV/40 mA) and an X'Celerator detector. Automatic divergence slits and 0.02 radian soller slits were used. Patterns were collected from  $10^{\circ}$  to

	Sample $a$	Sample $f$
Measurement type	Full intensity dataset	Full intensity dataset
Diffractometer	Enraf Nonius CAD4	Bruker Smart Apex2 CCD
X-ray radiation	$MoK\alpha$	MoKα
X-ray power	50 kV, 40mA	50  kV, 40  mA
Crystal size	$0.2 \ge 0.13 \ge 0.05 \text{ mm}^3$	$0.15 \ge 0.025 \ge 0.05 \text{ mm}^3$
Measurement time	Max. $120 \text{ s/step}$	60  sec/frame
Temperature (°C)	25	25
Space group	$Cmc2_1$	$Cmc2_1$
a-axis (Å)	5.6686(2)	$5.6618(7)^{\mathrm{a}}$
b-axis (Å)	15.0060(9)	$14.9981(12)^{\rm a}$
c-axis (Å)	7.6892(4)	$7.6806(7)^{a}$
Volume $(Å^3)$	654.07(6)	$652.21(12)^{a}$
Z	4	4
$ ho (g/cm^3)$	2.942	2.950
Reflections collected	4294	5915
Max. 2θ (°)	69.93	69.94
Index range $h$	-9 9	-99
Index range $k$	-24 24	-24 23
Index range $l$	-12 12	-12 11
Resolution range (Å)	$1 \dots 0.6$	$1 \dots 0.62$
Unique reflections	1166	1130
$\operatorname{Reflections} > 4 \sigma(I)$	1045	1088
R(int)	0.0510	0.0399
$R(\sigma)$	0.0375	0.0294
L. S. parameters	83	83
GoF	0.990	1.046
$R1,I>4{f \sigma}(I)$	0.0173	0.0186
R1, all data	0.0273	0.0199
w $R2$ on $F^2$	0.0295	0.0370
$\Delta \rho_{\rm max}$ , close to	$-0.21 e^{-} \cdot A^{-3}$ , O1	$-0.24 \text{ e}^- \cdot \text{Å}^{-3}$ , Al1
$\Delta \rho_{\rm max}$ , close to	$0.24 e^{-} \cdot A^{-3}, B1$	$0.24 e^{-} \cdot A^{-3}$ , Al2

 Table 4.4: Starting materials, synthesis conditions and analytical methods used for sample characterization.

<sup>a</sup>lattice parameters from measurement with CAD4 diffractometer

 $80^{\circ} 2\theta$  with a step size of  $0.002^{\circ} 2\theta$ /step at 100 s/step. Lattice parameters were derived from Pawley refinements to precisely extract peak maxima by treating the data independently from the structural model. Excess  $Al_2O_3$  (if present) was quantified by Rietveld refinements. It was not possible to estimate the amount of excess  $B_2O_3$  or  $H_3BO_3$  prior to the final washing step due to fast hydration of  $B_2O_3$ , resulting in a poorly crystalline mix of  $B_2O_3$  and  $H_3BO_3$ . All powder XRD data were handled with Topas-Academic v.4.1 (Coelho 2007a) using the fundamental parameter approach for peak-shape modeling.

# 4.4.3 Solid-State <sup>11</sup>B and <sup>27</sup>AI MAS-NMR

 $^{11}\text{B}$  and  $^{27}\text{Al}$  MAS-NMR spectra were recorded on a Bruker Avance 400 NMR spectrometer (9.4 T) using a 2.5 mm CP/MAS probe. The  $^{11}\text{B}$  and  $^{27}\text{Al}$  MAS-NMR spectra were recorded

at 128.38 MHz and 104.26 MHz, respectively using the following parameters: 0.3  $\mu$ s  $\pi/6$  pulse widths, 25 kHz MAS rate, 3 s (0.2 s for <sup>27</sup>Al) relaxation delays, appropriate number of scans for reasonable signal to noise ratios and 77 kHz SPINAL-64 proton decoupling (Fung et al. 2000).  $^{11}B$  and  $^{27}Al$  chemical shifts were referenced to external samples of 1 M aqueous  $H_3BO_3$  at 19.6 ppm (Kroeker and Stebbins 2001) and  $1.1 \text{ m Al}(\text{NO}_3)_3$  solution at 0.0 ppm (Harris et al. 2001). The observed <sup>27</sup>Al NMR line shape was exactly the same when no proton decoupling was applied, whereas the  $^{11}B$  quadrupolar powder pattern of e.g.  $H_3BO_3$  subtly depended on the efficiency of the decoupling field (Klochko et al. 2009). Owing to the presence of boron nitride devices in the probe, a  $\approx 16$  kHz broad, asymmetric background signal was present in the <sup>11</sup>B NMR spectra. Before further analysis of these spectra, the background signal recorded with an empty spinner was subtracted from the spectrum of interest. Quadrupolar parameters and the relative amounts of B<sup>(III)</sup> and/or B<sup>(IV)</sup> atoms were determined by non-linear least-square fits of the regions of interest using the software Dmfit v. 20080716 (Massiot et al. 2002). For H<sub>3</sub>BO<sub>3</sub>, the parameters  $C_q = 2.51$  MHz,  $\delta_{iso} = 19.7$  ppm and  $\eta = 0.04$  were in good accordance with literature data (Klochko et al. 2009). For resonances originating from single components only, the parameters for amplitude, position, quadrupolar coupling constant, asymmetry parameter and the zero order base line were automatically and independently optimized by the fit routine. In the case where  $B^{(III)}$  and  $B^{(IV)}$  resonances were observed simultaneously  $(Al_4B_2O_q,$ Mazza et al. 1992 and Fischer et al. 2008b), the asymmetry parameter was kept constant for the trigonal B<sup>(III)</sup> site and a Gaussian shape was chosen for tetrahedrally coordinated B<sup>(IV)</sup> in order to compare results to those of Fischer et al. (2008b).  $^{27}\mathrm{Al}$  NMR parameters ( $\delta_{\mathrm{iso}},$  $C_{\alpha}$ ,  $\eta$ ) were determined from slices of a z-filtered MQ-MAS-NMR spectrum and the 1D NMR spectra were subsequently simulated keeping the preliminary evaluated parameters C<sub>q</sub> and  $\eta$  constant, while the parameters for amplitude, chemical shift and zero order base line were optimized.

### 4.4.4 FTIR and Raman Spectroscopy

FTIR and Raman spectra were collected for sample e and  $Al_4B_2O_9$  as reference for tetrahedrally coordinated BO<sub>4</sub>. IR powder spectra were acquired from 300 cm<sup>-1</sup> to 4000 cm<sup>-1</sup> on a Bruker Tensor 27 FTIR spectrometer equipped with a globar MIR light source, a KBr beam splitter, and a DLaTGS detector. Sample and background spectra were averaged from 100 scans at 4 cm<sup>-1</sup> resolution. Two methods were used to obtain absorption spectra: (1) The finely ground sample was dispersed in KBr ( $\approx$ 1:200), pressed to transparent pellets, and measured in transmission mode in the usual sample compartment. (2) The undiluted sample powder was pressed on the diamond window of a Harrick MVP 2 diamond ATR accessory. Background spectra were obtained from a pure KBr pellet and from the empty ATR unit. Data handling was performed with OPUS v. 5.5 software (Bruker Optik GmbH 2005).

Powder Raman spectra were obtained from 70 cm<sup>-1</sup> to 1670 cm<sup>-1</sup> on a confocal edge filterbased Renishaw RM1000 micro-Raman system equipped with a 17 mW HeNe-laser (632.8 nm excitation) and a 50 mW multimode  $Ar^+$ -laser (488 nm and 514.5 nm excitation, each  $\approx 20$  mW), a 1200 lines/mm grating, using a thermo-electrically cooled CCD detector. Raman intensities were collected with a Leica DMLM microscope with a 50x/0.85 n.a. objective. Excitation at 488 nm and 10 min acquisition time yielded Raman spectra with reasonable signal to noise ratio at a resolution of 5–6 cm<sup>-1</sup>. Data was processed with Grams32 software v. 4.14 (Galactic Industries Corp. 1996).

## 4.4.5 Laser-Ablation ICP-MS

 $Al_2O_3$  and  $B_2O_3$  wt.-% were measured in-situ on epoxy-embedded single-crystals from samples *a* and *f* (Table 4.3) by laser-ablation inductively-coupled-plasma mass-spectrometry (LA ICP-MS). The system consists of a pulsed 193 nm ArF Excimer laser Geolas Pro system (Lambda Physik, Germany) coupled with a Perkin Elmer ELAN DRCe quadrupole mass spectrometer. Details on the setup and optimization strategies to minimize matrix effects by setting up robust plasma conditions can be found in Pettke (2008).

 $Al_2O_3$  inclusion-free crystal domains in sample a were usually smaller than 50 µm; Hence, the beam diameter for analysis was set to 32 µm. To minimize matrix-load induced elemental fractionation (Kroslakova and Günther 2007), ablation rate tests were performed at 6 J/cm<sup>2</sup> energy density on the sample with a 10 Hz pulse rate, revealing a much higher ablation rate for NIST SRM 610 used for calibration than for the aluminoborate crystals. Calibration shots were thus made with 16 µm beam size, resulting in near-equal aerosol masses ablated per unit time. Only 10 s signals were used for quantification, to ensure a crater aspect (depth/diameter) ratio < 1. The aerosol was transported to the ICP-MS using mixed He-Ar gas. For samples *a* and *f*, a total of 11 and 12 individual spot analyses were acquired, respectively. Data quantification was done using SILLS v. 1.2.0 (Guillong et al. 2008), employing 356 ±7 µg/g for B (Pearce et al. 1997) and 2.04 wt-% Al<sub>2</sub>O<sub>3</sub>. Internal standardization was carried out by summing  $B_2O_3$  and  $Al_2O_3$  to 100 wt-%.

# 4.5 Results

## 4.5.1 X-ray Diffraction

Single-crystal X-ray (CAD4) refined cell dimensions of sample *a* yielded V = 654.07(6) Å<sup>3</sup> whereas corresponding parameters (CAD4) for sample *f* were significantly smaller: V = 652.21(12) Å<sup>3</sup> (Table 4.4).

Three refinement strategies were followed for both crystals a and f: (1) The structures were refined with all positions fully occupied. (2) The occupancy of the Al4 tetrahedron was fixed at 9% B  $\rightarrow$  Al substitution according to the suggestion of Ihara et al. (1980), and (3) a mixed Al, B population was refined at each Al site in separate refinements. In strategy (3), intensity data were restricted to 0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å (high-angle data) in order to reduce contributions from bonding electrons. X-ray diffraction data containing bonding electron information are not properly modeled in standard site-occupancy refinements and lead to incorrect occupancy fractions (Armbruster et al. 1990).

After refinement in accordance with model (1), atomic displacement parameters ( $U_{eq.}$ ) were very similar for all Al sites except for Al2.  $U_{eq.}$  of Al2 was about 12% higher than the average of the remaining three.

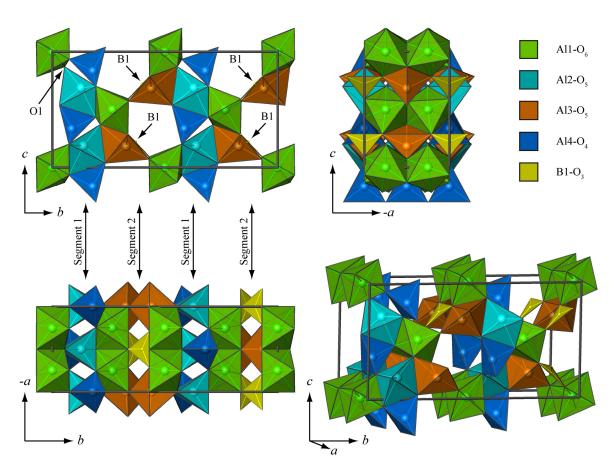
Strategy (2) led to increased least squares agreement factors compared to (1) and most important, the atomic displacement parameter of the tetrahedrally coordinated Al4 site decreased to ca. 50% of the value refined in strategy (1).

Strategy (3) resulted in fully occupied Al-polyhedra, except for Al2. This type of population refinement has to be performed stepwise for each Al site to reduce correlations with the scale factor. Due to correlations between displacement parameters and occupancy, vacancies at the Al2 site were also refined with an isotropic displacement parameter of Al2 constrained to the one of Al3. No significant difference in vacancy concentration was observed, thus, we chose to refine all displacement parameters individually. For Al2, final occupancies were Al:B = 0.973(5):0.027(5) for sample *a* and Al:B = 0.962(6):0.038(6) for sample *f*. As an alternative approach, Al vacancies were refined at the Al2 site, resulting in 2.1(4)% vacancies for sample *a* and 3.2(4)% for sample *f*. All refinements of strategy (3) reduced the originally (according to strategy (1)) enlarged atomic displacement factor  $U_{eq}$  of Al2 to values similar to the ones of other Al sites within the same structure.

Atomic coordinates, isotropic and anisotropic displacement parameters of the vacancy model (strategy 3) for samples a and f and calculated bond valences of sample a are listed in Appendix A.1.

The main features of the structure (Figure 4.3) are isolated mullite-like chains of edgesharing Al1-octahedra running parallel to the *a*-axis. Compared to mullite, the cell is doubled along [100] and [010] resulting in an additional symmetry equivalent octahedral chain at b/2. Proximate octahedral chains are connected by pairs of irregular edge-sharing AlO<sub>5</sub> polyhedra (Al2 and Al3) and, additionally, by alternating AlO<sub>4</sub> tetrahedra (Al4) and BO<sub>3</sub> triangles (B1). One side of an edge-connected AlO<sub>5</sub> polyhedra pair is edge-connected to octahedra and corner-linked to BO<sub>3</sub> triangles, whereas the other side is edge-connected to octahedra and corner-linked to tetrahedra. As a result, octahedral chains are separated along [010] by two different alternating segments within (101). One segment comprises AlO<sub>4</sub> and AlO<sub>5</sub> polyhedra (segment 1, Figure 4.3); The other consists of BO<sub>3</sub> and AlO<sub>5</sub> polyhedra (segment 2). The Al1 octahedron, the tetrahedron and the BO<sub>3</sub> triangle are rather regular:  $\Delta$ (Al1-O)<sub>max</sub> = 0.0553 Å,  $\Delta$ (Al4-O)<sub>max</sub> = 0.0250 Å and  $\Delta$ (B1-O)<sub>max</sub> = 0.0138 Å, whereas the AlO<sub>5</sub> polyhedra are fairly distorted with  $\Delta$ (Al2-O)<sub>max</sub> = 0.2966 Å, and  $\Delta$ (Al3-O)<sub>max</sub> = 0.3692 Å(Appendix A.1). The AlO<sub>5</sub> polyhedra are more precisely described as AlO<sub>4+1</sub>, as Al occupies the center of the four closest O ligands, which is also responsible for the increased distortion.

After preferred-orientation corrections in XRD Rietveld refinements, all patterns from samples a-f were matching the Al<sub>5</sub>BO<sub>9</sub> structure, differing only in the amount of excess corun-



**Figure 4.3:** Structural drawing of an  $Al_5BO_9$  unit cell projected along the *a*-axis (top left), the *b*-axis (top right), the c-axis (bottom left) and in an arbitrary view (bottom right). Segments dividing the mullite-like  $AlO_6$  chains are indicated with arrows: Segment 1 consists of  $Al2-O_5$  polyhedra and  $Al4-O_4$  tetrahedra, whereas segment 2 contains  $Al3-O_5$  polyhedra and  $B1-O_3$  triangles. Note that Al2 and Al3 atoms are not in the center of the polyhedron. For better illustration of the similarity to the mullite structure, the unit cell has been shifted by 0, 0.1164, 0.5.

dum. Pawley-refinement-derived lattice parameters of all samples are shown in Table 4.5. They are the same within a maximum deviation of 6 esd's, except for the significantly smaller values of sample f, confirming our previous single-crystal results. Excess corundum was only found in samples a-c (Table 4.3) whereas no corundum was found in samples d-f, prepared with excess  $B_2O_3/H_3BO_3$  in the starting mixtures (Table 4.5).

# 4.5.2 Solid-State <sup>11</sup>B and <sup>27</sup>AI MAS-NMR

In the <sup>11</sup>B MAS-NMR spectrum of  $Al_4B_2O_9$ , which was collected as a reference, signals of highly symmetric four-fold coordinated  $B^{(IV)}$  and second-order quadrupolar broadened planar three-fold coordinated  $B^{(III)}$  were observed at distinct positions (Figure 4.4). The ratio of  $\approx 25:75$  determined by line shape analysis for the  $B^{(IV)}/B^{(III)}$  sites fits reasonably to the value of  $\approx 20:80$  estimated by Fischer et al. (2008b). Nevertheless, accurate quantitative results can

Sample	a (Å)	b (Å)	c (Å)	V (Å <sup>3</sup> )	GoF	$\rm Al_2O_3~(wt\%)$
a	5.66848(4)	15.00512(12)	7.68971(5)	654.057(8)	1.04	8
b	5.66738(3)	15.00687(10)	7.69230(4)	654.225(7)	1.52	7
С	5.66789(4)	15.00604(14)	7.69157(6)	654.188(9)	1.47	8
d	5.66775(3)	15.00642(6)	7.68943(3)	654.006(5)	1.60	0
e	5.66869(2)	15.00741(7)	7.69005(3)	654.211(5)	1.47	0
f	5.66421(3)	15.00214(7)	7.68428(3)	652.974(5)	1.83	0
$a^{\mathrm{a}}$	5.6686(2)	15.0060(9)	7.6892(4)	654.07(6)	n/a	n/a
$f^{\mathbf{a}}$	5.6618(7)	14.9981(12)	7.6806(7)	652.21(12)	n/a	n/a

Table 4.5: Lattice parameters and excess  $Al_2O_3$  content of all samples determined by powder X-ray diffraction. For better comparison, lattice parameters obtained from single-crystal XRD are given in the two bottom rows.

<sup>a</sup>lattice parameters from single-crystal XRD for comparison

only be evaluated including the signal intensities of spinning sidebands.

The <sup>11</sup>B MAS-NMR spectra of samples a-f showed the same shapes of resonances (Figure 4.4). Chemical shifts and quadrupolar parameters obtained by line shape analysis (Table 4.6) yielded convergent results for all synthesis routes (Table 4.3). The influence of a simulated  $B^{(IV)}$  signal on the expected experimental line shape of sample a is shown in Figure 4.5. For fractions as low as 2–3% of  $B^{(IV)}$ , a distinct deviation of the quadrupolar broadened line shape of the spectra can be observed. As shown in Figure 4.4, none of the measured <sup>11</sup>B MAS-NMR spectra suggest evidence for  $B^{(IV)}$ .

<sup>27</sup>Al MAS-NMR spectra of samples a-f match each other (Figure 4.6), only the amount of excess corundum varies, depending on synthesis conditions. In the <sup>27</sup>Al 3Q-MAS-NMR spectrum of sample a (Appendix A.1) at least three distinct aluminum species were observed and from chemical shift arguments  $Al^{(IV)}$ ,  $Al^{(V)}$  and  $Al^{(VI)}$  are present. The spectrum showed the same resonances already observed by Gan et al. (2002), recorded at the same magnetic field. Since they also collected MQ-MAS data at higher frequencies, it is evident that two different  $Al^{(V)}$  sites are present. DOR NMR diffusion experiments (Hung et al. 2006) allowed assignment of the  $Al^{(V)}$  signals to  $Al^{(V)}_1$  (Al2 in this study) and  $Al^{(V)}_2$  (Al3 in this study). All 1D <sup>27</sup>Al MAS-NMR line shape simulations performed in our study (Figure 4.7) yielded four different Al species:  $Al^{(IV)}$ ,  $Al^{(V)}_1$ ,  $Al^{(V)}_2$  and  $Al^{(VI)}$  (Appendix A.1) with a ratio of ca. 1:1:1:2. The shape of the 1D <sup>27</sup>Al NMR spectra of sample d and e (lowest amounts of  $Al_2O_3$ ) are both in perfect agreement with literature data (Massiot et al. 1995).

## 4.5.3 Raman and FTIR Spectroscopy

FTIR powder spectra are shown in Figure 4.8. Note that IR band positions from the KBr and ATR techniques are almost identical, although different intensities result in somewhat different band patterns. Moreover, compared to KBr spectra, ATR band positions are systematically shifted to slightly lower values, which is a common effect in ATR spectra (Harrick 1967). In contrast, due to different selection rules, Raman spectra appear quite different (Fig-

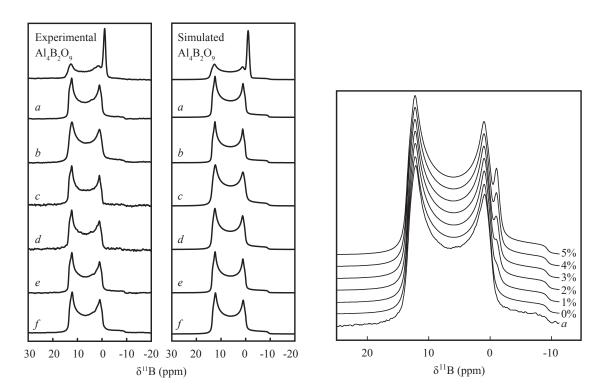


Figure a-f. In the spectra of Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>, the narrow signals of 0-5% tetrahedrally coordinated BO<sub>4</sub>. resonance around -1 ppm is assigned to BO<sub>4</sub>.

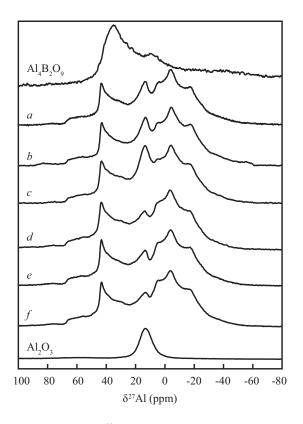
4.4: Experimental and simulated Figure 4.5: Experimental <sup>11</sup>B{<sup>1</sup>H} MAS-NMR <sup>11</sup>B $^{1}$ H MAS-NMR spectra (128.38 MHz) of spectra of powdered single-crystals from sample a reference  $Al_4B_2O_9$  powder and powdered samples with a series of simulated spectra considering the

Sample	$\delta^{11} {\rm Biso}~({\rm ppm})$	$C_q$ (MHz)	η
$H_3BO_3$	19.7	2.51	0.04
$Al_4B_2O_9^a$	17.4	2.66	0.10
a	16.8	2.62	0.09
b	16.7	2.61	0.09
С	16.8	2.61	0.09
d	16.8	2.61	0.08
e	16.8	2.62	0.08
f	16.7	2.61	0.09

Table 4.6: Chemical shifts and <sup>11</sup>B quadrupolar parameters obtained by line shape simulation.

 $^{\rm a}{\rm the}~{\rm BO}_4$  site was simulated by a Gaussian shape at -1.0 ppm (line width of  ${\approx}140$  Hz). Ratio determined:  $B^{(III)}/B^{(IV)} \approx 75:25$ 

ure 4.9). Considering the structural units of the investigated borates, the spectral regions (both IR and Raman) may be assigned to certain vibrations. The antisymmetric stretching vibrations of the BO<sub>3</sub> group are expected at  $\approx 1450 \text{ cm}^{-1}-1200 \text{ cm}^{-1}$  (predominant in IR spectra, Figure 4.8), the symmetric stretching vibration at  $\approx 900 \text{ cm}^{-1}$ -1050 cm<sup>-1</sup> (predominant in Raman spectra, Figure 4.9). The characteristic IR-active antisymmetric stretching bands of the BO<sub>4</sub> tetrahedron in the vibrational region at  $\approx 950 \text{ cm}^{-1}-1200 \text{ cm}^{-1}$  are only observed



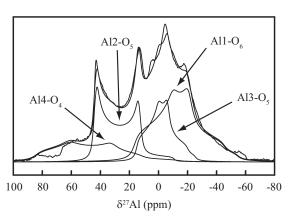
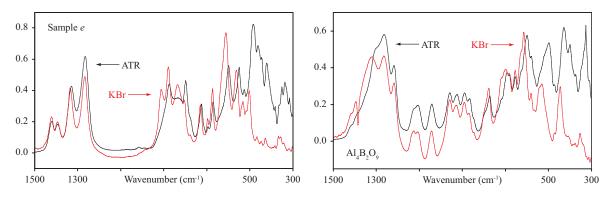


Figure 4.6: 1D <sup>27</sup>Al MAS-NMR spectra (104.26 Figure 4.7: 1D <sup>27</sup>Al MAS-NMR spectrum of corundum (bottom).

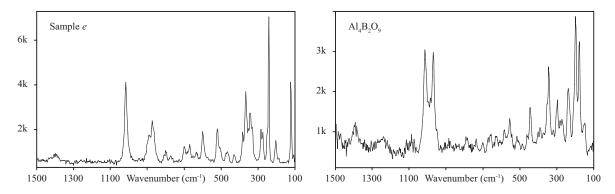
MHz) of  $Al_4B_2O_9$  (top), samples a-f and pure sample b with spectrum simulation using the quadrupolar parameters extracted from simulation of the  $^{27}\mathrm{Al}$  3Q-MAS-NMR spectrum.

for the  $Al_4B_2O_9$  reference but not in sample e (Figure 4.8). The strong  $BO_3$  characteristic symmetric stretching mode of the  $BO_3$  group occurs in Raman spectra at 1016 cm<sup>-1</sup> (Figure 4.9). The Raman spectrum of the reference material  $Al_4B_2O_9$  shows in addition a strong  $BO_4$ characteristic band at  $\approx 960 \text{ cm}^{-1}$ , which is absent in sample e (Figure 4.9). The bending motions of the  $BO_3$  group and all other vibrations of the  $AIO_x$  polyhedra and lattice vibrations are contained in the complex band region between 900 cm<sup>-1</sup> and 100 cm<sup>-1</sup>. Because of this complexity, the latter will not be discussed.

To answer the question for boron speciation  $(B^{(III)} \text{ and/or } B^{(IV)} \text{ coordination, one or more})$ structural sites) the possible vibrations of a  $BO_3$  group need to be considered. The ideal anion group (e.g. in solution or in a calcite-type structure) has symmetry  $D_{3h}$  resulting in four possible vibrations (Appendix A.1). Two of them, i.e. the E' species  $v_3$  and  $v_4$ , are doubly degenerate and active in both IR and Raman spectra. In contrast,  $v_1$  is only Raman-active and  $v_2$  is only IR-active. If the symmetry of the BO<sub>3</sub> group is lowered in a crystal structure, the degenerate E' species split into two bands, and formerly inactive vibrations may become active. In  $Al_5BO_9$  the Wyckoff site of boron is 4a with site symmetry m (Appendix A.1). The effective symmetry of the  $BO_3$  group, however, is higher, as is indicated by the very similar bond lengths of B-O3 (1.36 Å) and 2 x B-O7 (1.38 Å) (Appendix A.1), and almost identical bond angles close to 120°, i.e. it represents a flat isosceles triangle with symmetry  $C_{2v}$  (mm2). Independently, if the true site symmetry or the effective symmetry are preferred, the vibrational E' modes split up and the selection rules are released.



**Figure 4.8:** ATR and KBr pellet FTIR spectra of sample e (left) and  $Al_4B_2O_9$  (right). The well resolved peaks between 1250 cm<sup>-1</sup> and 1450 cm<sup>-1</sup> (left) are assigned to the BO<sub>3</sub> group in sample e. Peaks of stretching vibrations of BO<sub>4</sub> in reference  $Al_4B_2O_9$  (right) are between 950 cm<sup>-1</sup> and 1200 cm<sup>-1</sup> (right). In the spectra of sample e, no indication for BO<sub>4</sub> can be found. Intensity scale is (ATR) absorbance



**Figure 4.9:** Raman spectra of sample e (left) and reference  $Al_4B_2O_9$  (right). The single strong Raman band at 1016 cm<sup>-1</sup> is assigned to the symmetric stretching mode of the BO<sub>3</sub> group (left), whereas the double peak in the  $Al_4B_2O_9$  pattern corresponds to vibrations of BO<sub>3</sub> and BO<sub>4</sub> groups. Intensity scale is arbitrary units.

In addition to the symmetry considerations above, it is important to note that boron consists of two abundant natural isotopes <sup>11</sup>B and <sup>10</sup>B with a ratio of about 80:20 (Hannaford and Lowe 1977). This isotope ratio is also found in the common boron-bearing chemicals, such as  $B_2O_3$  and  $H_3BO_3$  used in the syntheses of the present study. Due to the different mass of the isotopes (10 rel.-%), the frequencies of vibrations, where motions of boron isotopes are involved (therefore not in  $v_1$ !), are different by several tens of cm<sup>-1</sup> (Siebert 1966, Farmer 1974, Harris and Bertolucci 1989).

## 4.5.4 Laser-Ablation ICP-MS

LA ICP-MS measurement data for samples a and f are reported in Table 4.7. Compositions of the two synthetic products are identical. They are marginally higher with  $\approx 12.35$  wt.-%  $B_2O_3$  than expected for pure  $Al_5BO_9$  (12.01 wt.-%  $B_2O_3$ ), but significantly lower than for  $Al_{18}B_4O_{33}$  (13.17 wt.-%  $B_2O_3$ , Table 4.1). External reproducibility of one standard deviation uncertainties on  $B_2O_3$  are ca. 0.25 wt.-% which is of the same order as the uncertainty quoted for the B concentration of the standard reference material NIST SRM 610 used for calibration (Pearce et al. 1997).

Spot no.	Sample $a$		Sample $f$	
	$Al_2O_3$	$B_2O_3$	$Al_2O_3$	$B_2O_3$
1	87.71	12.29	87.65	12.35
2	87.75	12.25	87.65	12.35
3	87.63	12.37	87.99	12.01
4	87.72	12.28	87.33	12.67
5	87.69	12.31	87.72	12.28
6	87.75	12.25	87.29	12.71
7	87.60	12.40	87.72	12.28
8	87.72	12.28	87.69	12.31
9	87.84	12.16	87.78	12.22
10	87.09	12.91	87.08	12.92
11	87.71	12.29	87.89	12.11
12	87.86	12.14	n/a	n/a
Average	87.66	12.33	87.62	12.38
$\operatorname{ESD}$	0.20	0.20	0.27	0.27
ESD SRM610 <sup>a</sup>	n/a	0.25	n/a	0.25

Table 4.7:  $Al_2O_3$  and  $B_2O_3$  concentrations (wt.-%) from LA ICP-MS.

<sup>a</sup>standard deviation according to uncertainty of NIST SRM 610 (Pearce et al. 1997)

# 4.6 Discussion

The excellent agreement of cell dimensions (Table 4.5) of samples a-e synthesized along different routes and starting compositions suggests that they are structurally and chemically very similar differing in the amount of excess corundum. This is also confirmed by <sup>11</sup>B and <sup>27</sup>Al MAS-NMR data (Figures 4.4 and 4.6, Table 4.6). LA ICP-MS compositional data of samples a and f are equal within esd's (Table 4.7). Atomic coordinates of samples a and f(single-crystal X-ray data) are also identical. Indications for differences between sample a and f are: (1) The cell dimensions (Table 4.5), which are most sensitive due to their ability to sum up small structural differences, (2) the crystal from sample f had also systematically larger atomic displacement parameters than the crystal from sample a (ca. 20% for cation sites, corresponding to ca. 10 esd's, and ca. 10% for O sites, corresponding to 5 esd's). Both data sets were of corresponding quality but were measured on different machines. Nevertheless, we have previously tested reference crystals to corroborate that both machines produce comparable results. Thus, the increased displacement parameters of sample f must be related to crystal properties, such as increased strain compared to the crystal from sample a. The only systematic experimental difference between samples a-e and sample f, distinct by their unit cell volumes (Table 4.5), is the higher synthesis temperature for sample f (at 1550°C cooled to 1250°C) whereas samples a-e were treated at 1200°C or below. The higher crystallization temperature and subsequent air quenching to ambient conditions could explain the suspected increased strain. But is  $B \rightarrow Al$  substitution the origin of the significantly smaller unit cell volume of sample f? LA ICP-MS results (Table 4.7) seem to exclude this interpretation. A possible answer will be discussed below.

The "hypothetical"  $Al_{18}B_4O_{33}$  (=  $Al_{4.91}B_{1.09}O_9$ ) composition with the structure of  $Al_5BO_9$  can only be achieved if:

- 9% B substitutes for Al at the Al4 tetrahedron according to Ihara et al. (1980). B is only known in three- and four-fold coordination by O (Hawthorne et al. 1996, Leonyuk 1997). Therefore, B is not expected to occupy one of the other five- or six-fold coordinated sites within the mullite type Al<sub>5</sub>BO<sub>9</sub> structure.
- (2) Minor Al vacancies at one of the other Al sites (six-coordinate Al1 or five-coordinate Al2 and Al3) may be charge balanced by a BO<sub>4</sub> or BO<sub>3</sub> polyhedron at a nearby interstitial position. In case of Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> = Al<sub>4.91</sub>B<sub>1.09</sub>O<sub>9</sub>, 9% Al vacancies may either be located (2a) at one single position, or (2b) 9% vacancies are statistically distributed throughout all Al positions in the structure.

Ad (1): Single-crystal XRD investigations on crystals from samples a and f do not show any indication for the 9% B  $\rightarrow$  Al4 substitution as proposed by Ihara et al. (1980). If applying a constraint of 9% B and 91% Al to the tetrahedral Al4 site, agreement factors increase (R1: from 0.0198 to 0.0203 for sample a and 0.0212 to 0.0218 for sample f) and most important, the atomic displacement factor ( $U_{eq.}$ ) at Al4 becomes halved. This  $U_{eq.}$  behavior has a straightforward explanation: If the electron density (occupancy) at a structural site is underestimated in the refinement model, the probability density cloud around the atomic site (represented by  $U_{eq.}$ ) contracts because integration over the observed electron density is already satisfied (according to the model) for a smaller cloud volume. Theoretically (Kunz and Armbruster 1990) the opposite should be observed: Occupational disorder leads to increased displacement parameters. Furthermore, the Al4 tetrahedron is rather undistorted with  $\Delta$ (Al4-O)<sub>max</sub> = 0.0250 Å and the average bond length <Al4-O>tet = 1.7464(12) Å is in agreement with a fully occupied AlO<sub>4</sub> tetrahedron (Kunz and Armbruster 1990). By assuming a mean tetrahedral bond length for BO<sub>4</sub> of 1.476 Å (Hawthorne et al. 1996), the average bond length of a tetrahedron occupied with 91% Al and 9% B is expected to be  $\approx 1.72$  Å.

<sup>11</sup>B MAS-NMR spectra show no evidence for tetrahedrally coordinated  $B^{(IV)}$  in samples a-f. Due to the low detection limit and the good agreement of the  $B^{(III)}:B^{(IV)}$  ratio for

 $Al_4B_2O_9$  (Fischer et al. 2008b), partially occupied tetrahedral  $BO_4$  sites with occupancies > 2% should be clearly detectable (Figure 4.5). From FTIR spectra, 9% of  $B^{(IV)}$  in tetrahedral coordination can also be excluded, because the characteristic vibrational region of the IR-active antisymmetric stretching vibrations of the  $BO_4$  tetrahedron at  $\approx 950 \text{ cm}^{-1}$ -1200 cm<sup>-1</sup> (Farmer 1974) is empty (Figure 4.8). Moreover, there is only one Raman band at 1016 cm<sup>-1</sup> (Figure 4.9) that can unequivocally be assigned to the symmetric stretching mode of the  $BO_3$  group; Hence, no band is left to be assigned to  $v_1$  of a potential  $BO_4$  group.

Ad (2a): Site occupancy refinements on crystals a and f provided no evidence for 9% vacancies at Al1, Al2 or Al3. In addition, all spectroscopic methods applied in this study yielded no indication of additional BO<sub>4</sub> or BO<sub>3</sub> with 9% occupancy.

The single, strong Raman band at 1016 cm<sup>-1</sup> (Figure 4.9) does not suggest an additional BO<sub>3</sub> group occupied to 9%. Within detection limits (2–3%), IR and Raman spectra (Figures 4.8 and 4.9 and Appendix A.1) are consistent with assumption of a single BO<sub>3</sub> group, considering both natural B isotopes. The symmetric stretching mode  $v_1$  is visible only as a very weak band at 1015 cm<sup>-1</sup> in IR spectra (IR active due to distortion from the ideal symmetry of the planar BO<sub>3</sub> group). The different B isotopes do not split this mode, as the central B atom is almost inert during this vibration. A similar position (1017/1019 cm<sup>-1</sup>) of this  $v_1$  mode was observed in Raman spectra of another aluminoborate with additional REEs by Xia et al. (1999). The antisymmetric stretching mode (doubly degenerate in case of ideal planar symmetry) is split into two modes by the lower site symmetry and further doubled by the two B isotopes. The resulting four bands are clearly visible between 1250 cm<sup>-1</sup> and 1450 cm<sup>-1</sup> in Raman spectra (Xia et al. 1999).

Ad (2b): The hypothesis that ca. 9% vacancies are distributed over several Al-polyhedra (Al1<sup>(VI)</sup>, Al2<sup>(V)</sup>, Al3<sup>(V)</sup>, and Al4<sup>(IV)</sup>) and for charge balance BO<sub>3</sub> triangles at nearby interstitial positions are occupied, can also be rejected based on the above spectroscopic arguments. Most convincing, results of LA ICP-MS measurements clearly show that the investigated crystals of samples *a* and *f* cannot have a composition of  $Al_{18}B_4O_{33}$ . Averaged LA ICP-MS data of 12.36 wt.-%  $B_2O_3$  with a standard deviation of 0.25 wt.-% result in  $Al_{4.99}B_{1.01}O_9 \leq Al_{4.97}B_{1.03}O_9 \leq Al_{4.95}B_{1.05}O_9$ . Due to reasons mentioned above (e.g. similarity in unit cell volumes), it is very likely that these values apply for all samples.

One of the striking results of structure refinements presented in this study is the significantly large  $U_{eq.}$  value for Al2 compared to all other Al sites. This observation is consistent with previous, less accurate structural data (Sokolova et al. 1978, Ihara et al. 1980, Garsche et al. 1991). Results of single-crystal structure refinements allow the interpretation of a small amount B at Al2 (2.7(5)% for sample *a* and 3.8(6)% for sample *f*). In corresponding refinements, displacement parameters for the Al2 site are no longer larger than those of the other Al sites in refinements with fully occupied cation sites (Appendix A.1). In this model B occupies the center of the tetrahedron formed by O1, 2 x O2 and O5. For samples a <sup>11</sup>B{<sup>1</sup>H} MAS-NMR simulations of  $BO_4$  with more than ca. 2% occupancy in addition to  $BO_3$  can be excluded (Figure 4.5). On the other hand, the average of the four shortest Al2-O distances is ca. 1.79 Å. It seems that this type of substitution is rather unlikely, considering the large size difference due to the characteristic tetrahedral B-O bond length of 1.476 Å.

An X-ray site occupation refinement is mainly sensitive on the number of scattering electrons. Thus alternatively to partial B occupancy, 2.1(4)% and 3.2(4)% vacancies at Al2 may also be successfully modeled for samples a and f, respectively. In this case, we assume interstitial B in three-fold coordination for charge balance. The Al2 polyhedron comprises 5 oxygen ligands, of which O1 has the lowest bond valence (Appendix A.1). Therefore, a new boron position must be close to the Al2 site and close to O1. Due to the low occupancy, no distinct new boron position could be found in difference Fourier-maps of single-crystal X-ray structure refinements (2% B is equal to an electron density of 0.10e<sup>-</sup>). Nevertheless, we assume that the new boron position is centered within one of the three faces of the Al2 polyhedron sharing O1 as apex. A similar substitution is known from the natural mullite-type borosilicate werdingite (Niven et al. 1991), in which BO<sub>3</sub> groups are correspondingly disordered with Al-tetrahedra. If the interstitial boron site in  $Al_5BO_9$  is statistically distributed among all three faces, it is below the detection limit of FTIR/Raman or single-crystal X-ray diffraction methods. Low concentrations of vacancies at Al2 with B in three-fold coordination distributed in the O1-O5-O2 (twice due to symmetry equivalent sites) and the O1-O2-O2 faces would locally distort the structure because O-O separations are much shorter in  $BO_3$  than in the irregular  $AlO_5$  polyhedron. Therefore, slightly different concentrations of vacancies at the Al2 position compensated by interstitial BO<sub>3</sub> polyhedra could explain the observed difference in cell parameters of samples a and f.

None of the  $Al_2O_3$ -rich mullite-type aluminoborates synthesized above 1100°C had  $Al_{18}B_4O_{33}$  composition. The exact stoichiometry is close to  $Al_5BO_9$ . Values derived from single-crystal diffraction data suggest  $Al_{5-x}B_{1+x}O_9$  with 0.021(6) < x < 0.038(6), which agrees with compositional data from LA ICP-MS yielding x = 0.03(2). Considering the historical background, it is assumed that the claimed  $Al_{18}B_4O_{33}$  stoichiometry is probably an artifact due to old inaccurate chemical analyses.

### Acknowledgements

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# Note

This text represents the contents of the article "Crystal-chemistry of mullite-type aluminoborates Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> and Al<sub>5</sub>BO<sub>9</sub>: A stoichiometry puzzle", as published in the Journal of Solid-State Chemistry 184 (2011) 70-80. The composition of mullite Al<sub>2</sub>(Al<sub>2+2x</sub>Si<sub>2-2x</sub>)O<sub>10-x</sub>, 0.2 < x < 0.5 has been changed to  $0.18 \le x \le 0.88$  in the introduction and in Table 4.4, "Reflections  $> 2\sigma(I)$ " was changed to "Reflections  $> 4\sigma(I)$ ". Some passages have been slightly modified to fit the layout of this thesis.

# **5** Thermal Expansion of Aluminoborates

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### 5.1 Abstract

Borates are known for low thermal expansion due to rigid behavior of  $BO_3$  and  $BO_4$  groups. Aluminoborates  $Al_5BO_9$ ,  $Al_4B_2O_9$ , grandidierite (Mg,Fe) $Al_3BSiO_9$  and jeremejevite  $Al_6B_5$ - $O_{15}(F,OH)_3$  were investigated by temperature dependent X-ray diffraction methods. Thermal expansion of all investigated structures is mainly due to bending and stretching of M-O-M angles and unusual thermal behavior of  $AlO_5$  and  $MgO_5$  polyhedra. Anisotropic expansion of the studied structures is compared with data of structurally related compounds.

### 5.2 Introduction

Crystal structures containing triangular  $BO_3$  and tetrahedral  $BO_4$  units are known for low or even negative thermal expansion behavior. Filatov and Bubnova (2008) conclude that (1)  $BO_3$ or  $BO_4$  units behave as rigid bodies, comparable to  $SiO_4$  tetrahedra (Hazen and Prewitt 1977, Tucker et al. 2000); (2)  $BO_3$  and  $BO_4$  units act as hinges, which is a reason for anisotropic thermal expansion; (3) The more cations with low bond strength a borate structure contains, the higher is its thermal expansion.

Thermal expansion of borate structures is therefore controlled by translational and vibrational motions of  $BO_3$  and  $BO_4$  groups combined with expansion and/or tilting of other polyhedra in the structure.

Temperature dependent behavior of two synthetic aluminoborates  $Al_5BO_9$  (Sokolova et al. 1978, Fisch et al. 2011 and refs. therein) and  $Al_4B_2O_9$  (Fischer et al. 2008b and refs. therein) and the minerals grandidierite (Mg,Fe) $Al_3BSiO_9$  (Stephenson and Moore 1968) and jeremejevite  $Al_6B_5O_{15}(F,OH)_3$  (Golovastikov et al. 1955, Foord et al. 1981) were investigated by high-temperature X-ray diffraction (XRD) methods.  $Al_5BO_9$ ,  $Al_4B_2O_9$  and grandidierite belong to the group of boron-mullites. They have mullite-type crystal structures characterized by unfolded, infinite edge-connected chains of  $AlO_6$  octahedra and crystallize in subgroups of space group P4/mbm (Fischer and Schneider 2008a). Jeremejevite crystallizes in space group  $P6_3/m$  and has no structural relation to the mullite-types. The compound was included in this study because it only contains Al and B cations, as in synthetic  $Al_5BO_9$  and  $Al_4B_2O_9$ . The four investigated structures contain chains of rather undistorted edge-connected  $AlO_6$  octahedra and planar coordinated  $BO_3$  triangles together with  $AlO_5$  and  $AlO_4$  ( $Al_5BO_9$ ),  $AlO_5$ ,  $AlO_4$  and  $BO_4$  ( $Al_4B_2O_9$ ), and  $AlO_5$ ,  $SiO_4$  and  $MgO_5$  (grandidierite). In jeremejevite, Al octahedra have five O and one (F,OH) ligand.

### 5.3 Sample Description

The  $Al_5BO_9$  single-crystal (0.13 $\cdot$ 0.07 $\cdot$ 0.4 mm<sup>3</sup>) was obtained from a mixture of  $9Al_2O_3:2B_2O_3$ in a  $K_2CO_3 + 3MoO_3$  melt (slowly cooled from 1373K) whereas the powder sample resulted from a solid-state reaction of  $9 \text{ Al}_2\text{O}_3$  with  $2 \text{ B}_2\text{O}_3$  heated at 1473K.  $\text{Al}_4\text{B}_2\text{O}_9$  powder was prepared according to Fischer et al. (2008b).

The grandidierite fragment  $(0.17 \cdot 0.34 \cdot 0.27 \text{ mm}^3)$  was cut from a raw, bluish gemstone from Madagascar. The jeremejevite crystal  $(0.14 \cdot 0.14 \cdot 0.07 \text{ mm}^3)$  was separated from a pale light-blue sample from Erongo region, Namibia. The remaining jeremejevite material was powdered and mixed with corundum as internal standard for high-temperature powder XRD measurements. Grandidierite could not be measured with powder XRD due to the small sample available. All single-crystals were squeezed into quartz glass capillaries and subsequently mounted on a goniometer head for high-temperature data collection.

## 5.4 X-ray Diffraction Methods

Single-crystal XRD data were collected with a Bruker Smart APEX2 CCD diffractometer (graphite monochromatized  $MoK\alpha$ ) at 298K and 873K. During high-temperature measurements, crystals were heated using a hot N<sub>2</sub>-gas blower, calibrated to known phase transitions. Data were collected with  $\omega$ - $\phi$ -scans and subsequently treated using Apex2 v. 2009-11.0 software package (Bruker-AXS 2009). The structural model by Fisch et al. (2011) was used for refinement of Al<sub>5</sub>BO<sub>9</sub> data whereas grandidierite and jeremejevite structures were solved by direct methods. Details of single-crystal XRD data treatment are given in Table 5.1. Atomic coordinates and anisotropic displacement parameters are listed in Appendix A.2.

According to Hazen and Prewitt (1977), average bond-expansion in polyhedra obeys the following rule:  $\langle \alpha \rangle = 32.9 \cdot (0.75 - z/\rho) \cdot 10^{-6} \text{K}^{-1}$ , where z is the cation charge and  $\rho$  the coordination number. If  $z/\rho \geq 0.75$  (e.g. AlO<sub>4</sub>, SiO<sub>4</sub>, BO<sub>3</sub>),  $\langle \alpha \rangle$  is predicted as 0K<sup>-1</sup>. Average  $\alpha$  values were calculated from mean polyhedral bond distances  $\langle d \rangle$  according to:  $\langle \alpha \rangle = (\langle d_{873K} \rangle - \langle d_{298K} \rangle)/\langle d_{298K} \rangle \cdot (873K - 298K)$ .

Bond-lengths data were not corrected for rigid body thermal motion effects (Johnson 1970) because the differences between corrected and uncorrected data were only slightly larger than uncorrected data within standard deviations. Average bond expansion in polyhedra calculated from uncorrected data sometimes resulted in negative thermal expansion for  $BO_3$  units, with errors almost as large as the value itself.

Negative expansion with increasing temperature for  $MO_x$  polyhedra is physically not feasible, thus, negative values are artifacts of not applied corrections for thermal rigid body librational effects and may be interpreted as zero expansion between 298K and 873K.

Powder XRD data were measured with a PANalytical X'Pert PRO MPD diffractometer (Cu radiation) equipped with an Anton-Paar HTK 1200 high-temperature goniometer attachment. Samples were measured at 298K, 323K and then in 50K steps up to 1273K (Al<sub>5</sub>BO<sub>9</sub>), 1323K (Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>) and 1073K (jeremejevite). Diffraction patterns were measured from 10° 20 to 70° 20 at 0.008°/step with 30-90 s/step. To eliminate outliers, lattice parameters were parameterized by a 2<sup>nd</sup> order polynomial: Length<sub>i</sub>(T) =  $x_{0i} + x_{1i} \cdot T + x_{2i} \cdot T^2$ . Polynomial parameters  $x_{0i}$ ,

Sample	$Al_5$	$Al_5BO_9$	Jeremejevite	ejevite	Grandidierite	idierite
Temperature	$298 \mathrm{K}^{\mathrm{a}}$	873K	298K	873K	298K	873K
Space group	$Cmc2_1$	$Cmc2_1$	$P6_3/m$	$P6_3/m$	Pnma	Pnma
a-axis	5.6689(6) Å	5.6825(6) Å	8.5526(3) Å	8.5840(16) Å	10.9967(1) Å	11.0193(1) Å
b-axis	15.0045(13) Å	15.0754(16) Å	8.5526(3) Å	8.5840(16) Å	5.7634(1) Å	5.7884(1) Å
c-axis	7.6897(11) Å	7.7121(7) Å	8.1793(4) Å	8.2187(15) Å	10.3321(1) Å	10.3434(1) Å
Volume	654.07(13) Å <sup>3</sup>	660.67(12) Å <sup>3</sup>	518.13(4) Å <sup>3</sup>	524.46(17) Å <sup>3</sup>	$653.05(3)$ Å $^3$	$659.74(3)$ Å $^3$
Z	4	4	2	2	4	4
Time p. frame	$10 \mathrm{~s}$	30  s	$40 \mathrm{~s}$	$40 \mathrm{s}$	$10 \mathrm{~s}$	$20 \mathrm{\ s}$
Refl. coll.	4919	1277	5491	2522	10569	14415
Max. 20	$69.91^{\circ}$	$65.64^{\circ}$	$69.21^{\circ}$	$60.86^{\circ}$	$63.20^{\circ}$	$89.18^{\circ}$
Index range $h$	-88	-77	-1211	-912	-1514	-1621
Index range $k$	-2323	-2120	-1113	-1111	-88	-1010
Index range $l$	-1212	-1010	-128	-1111	-1415	-1518
Unique refl.	1110	595	781	555	1105	2468
Refl. $> 4\sigma(I)$	1056	581	673	465	1118	2283
$R(\mathrm{int})$	0.0423	0.0414	0.0324	0.0270	0.0167	0.0162
$R(\sigma)$	0.0349	0.0529	0.0222	0.0251	0.0095	0.0110
L. S. param.	83	83	50	50	87	87
GoF	1.001	1.076	1.064	1.236	1.278	1.262
$R1,I>4 {\rm s}(I)$	0.0195	0.0242	0.0258	0.0308	0.0182	0.0207
R1, all data	0.0211	0.0247	0.0326	0.0434	0.0186	0.0230
w $R2$ on $F^2$	0.0351	0.0557	0.0627	0.1145	0.0563	0.0636
$\Delta \rho_{\rm max}$ , near	$0.23 e^{} \text{Å}^{-3}, B1$	$0.26 e^{} Å^{-3}$ , All	$0.49 e^{} Å^{-3}, O1$	$0.69 e^{} Å^{-3}, O3$	0.46 e <sup>-</sup> ·Å <sup>-3</sup> , O6	$0.53 e^{} A^{-3}, O3$
$\Delta \rho_{\min}$ , near	$-0.18 e^{-1} \cdot Å^{-3}$ , O4	-0.26 e <sup>-</sup> ·Å <sup>-3</sup> , Al1	-0.40 e <sup>-</sup> ·Å <sup>-3</sup> , Al1	-0.73 e <sup>-</sup> ·Å <sup>-3</sup> , O1	$-0.39 \text{ e}^{-3}$ , Si1	$-0.37 \text{ e}^{-3}$ , Si1

5 Thermal Expansion of Aluminoborates

 $x_{1i}$  and  $x_{2i}$  (*i* are lattice parameters *a*, *b* and *c*) were refined simultaneously for all diffraction patterns using Topas-Academic v. 4.1 (Coelho 2007a). Sample displacement was refined as a function of temperature due to its correlation with lattice parameters. From polynomials expressing temperature dependent lattice parameters, thermal expansion parameters  $\alpha(T)$ were derived according to the formula  $\alpha(T) = a(T)^{-1} \cdot \delta a(T) / \delta T$ . From  $\alpha(T)$ , the anisotropy factor *A* can be obtained as follows:  $A = (|\alpha_a - \alpha_b| + |\alpha_a - \alpha_c| + |\alpha_b - \alpha_c|) \cdot 10^6$  (Dreyer 1974, Schneider and Eberhard 1990). For hexagonal jeremejevite ( $\alpha_a = \alpha_b$ ),  $A = 2|\alpha_a - \alpha_c| \cdot 10^6$ .

# 5.5 Results and Discussion

#### 5.5.1 Thermal Evolution of Al<sub>5</sub>BO<sub>9</sub>

Main feature of the  $Al_5BO_9$  structure are mullite-type chains of edge-connected  $AlO_6$  octahedra extending parallel to the *a*-axis. Adjacent  $AlO_6$  chains are interconnected by two alternating segments parallel (010). One contains  $Al2O_5$  and  $AlO_4$  polyhedra whereas the other consists of  $Al3O_5$  and  $BO_3$  polyhedra, which build corner-linked chains ( $-BO_3$ - $Al3O_5$ - $BO_3$ -) parallel *b* linking adjacent octahedra. Both  $AlO_5$  polyhedra are more accurately described as  $AlO_{4+1}$  polyhedra as one Al-O bond is significantly longer than the other four. The latter segment containing the  $BO_3$  unit is less dense than the former one, resulting in a cavity parallel to the  $AlO_6$  chains (Figure 5.1).

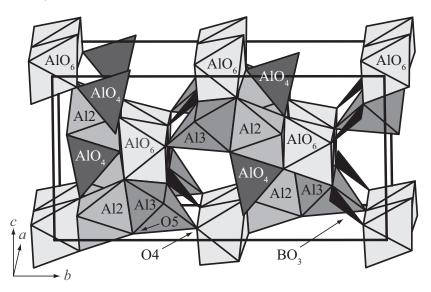


Figure 5.1: The structure of  $Al_5BO_9$ . The distance O4-O5 (marked) shows a large increase in the investigated temperature regime indicating that the corner linkages between  $Al3O_5$  and  $BO_3$  and the adjacent chain of octahedra are the weakest part of the structure.

Temperature dependent lattice parameters and  $\alpha(T)$  are shown in Tables 5.2 and 5.3, respectively. Between 298K and 1273K, the *b*-axis expands strongly, followed by the *c*- and the *a*-axis (the latter parallel to chains of AlO<sub>6</sub> octahedra). The volumetric expansion is +1.56%.

In the same temperature interval,  $\alpha_a$  increases most, followed by  $\alpha_c$  and  $\alpha_b$ , indicating that the shortest crystal axis (*a*-axis) has the steepest gradient, followed by the *c*-axis, whereas the longest crystal axis (*b*-axis) displays the lowest slope (Figure 5.2). As a result, thermal expansion of Al<sub>5</sub>BO<sub>9</sub> trends to more isotropic expansion behavior with increasing temperature, which is confirmed by a decrease of A from  $A_{298K} = 6.3$  to  $A_{1273K} = 2.4$ .

Inter-polyhedral angles remain the same within standard deviations in the investigated temperature interval, except the O5-Al3-O4 angle. This angle changes from 112.51(06)° to 113.56(13)°, resulting in an increase of the O4-O5 distance from 2.9236(17) Å to 2.9470(36) Å (Figure 5.1). Expansion of the AlO<sub>6</sub>, BO<sub>3</sub> and Al2O<sub>5</sub> polyhedra is in perfect agreement with the predicted values of Hazen and Prewitt (1977). However, for the Al3O<sub>5</sub> polyhedron,  $\langle \alpha \rangle = 8.9(2.5) \cdot 10^{-6} \text{K}^{-1}$  is about 180% of the expected value. In addition, the AlO<sub>4</sub> tetrahedron has  $\langle \alpha \rangle = 4.6(2.4) \cdot 10^{-6} \text{K}^{-1}$  instead of the expected value of 0K<sup>-1</sup> (Table 5.4). Thermal expansion of Al<sub>5</sub>BO<sub>9</sub> is therefore controlled by expansion of the Al3O<sub>5</sub> polyhedron and the AlO<sub>4</sub> tetrahedron.

One side of the  $-BO_3$ -Al3O<sub>5</sub>-BO<sub>3</sub>- chain linking two adjacent octahedral chains running parallel *a* is only corner-linked to AlO<sub>6</sub> polyhedra (Figure 5.1). As this is the only part of the structure where only corner-linkages exist, it is obvious that this joint is the weakest part of the structure, which agrees with Al3O<sub>5</sub> having the largest expansion of all polyhedra in the structure (Table 5.4).

Parameter	$x_0$	$x_1$	$x_2$
$Al_5BO_9$ for 298K $\leq$	$T \leq 1273 \mathrm{K}^{\mathrm{a}}$		
a(T) (Å)	5.66407	$1.44340 \cdot 10^{-5}$	$6.02392 \cdot 10^9$
b(T) (Å)	14.98650	$9.45645 \cdot 10^{-5}$	$9.24350 \cdot 10^{10}$
c(T) (Å)	7.68450	$3.61387 \cdot 10^{-5}$	$3.71557 \cdot 10^9$
V(T) (Å <sup>3</sup> )	652.299	$8.829 \cdot 10^{-3}$	$1.112 \cdot 10^{6}$
$Al_4B_2O_9$ for 298K $\leq$	$\leq T \leq 1323 \mathrm{K}^{\mathrm{a}}$		
a(T) (Å)	14.76128	$1.49152 \cdot 10^{-4}$	$-1.33666 \cdot 10^8$
b(T) (Å)	5.54025	$1.49592 \cdot 10^{-5}$	$4.33164 \cdot 10^9$
c(T) (Å)	15.05004	$8.43267 \cdot 10^{-5}$	$1.70444 \cdot 10^8$
V(T) (Å <sup>3</sup> )	1230.685	$2.263 \cdot 10^{-2}$	$1.417 \cdot 10^{6}$
Grandidierite for 29	$8 \mathrm{K} \leq T \leq 973 \mathrm{K}^{\mathrm{b}}$		
a(T) (Å)	10.9548	$6.2274 \cdot 10^{-5}$	$2.4334 \cdot 10^8$
b(T) (Å)	5.7557	$2.6772 \cdot 10^{-5}$	$1.5226 \cdot 10^8$
c(T) (Å)	10.3292	$5.7471 \cdot 10^{-6}$	$1.6289 \cdot 10^8$
V(T) (Å <sup>3</sup> )	651.29	$7.03 \cdot 10^{-3}$	$4.30 \cdot 10^{6}$
Jeremejevite for 298	$K \leq T \leq 1073 K^{a}$		
a(T) (Å)	8.53719	$6.51545 \cdot 10^{-5}$	$1.05757 \cdot 10^9$
c(T) (Å)	8.16183	$6.20481 \cdot 10^{-5}$	$-6.17363 \cdot 10^{10}$
V(T) (Å <sup>3</sup> )	515.169	$1.178 \cdot 10^{-2}$	$1.834 \cdot 10^{7}$

**Table 5.2:** Coefficients for the 2<sup>nd</sup> order polynomial length<sub>i</sub> $(T) = x_{0i} + x_{1i} \cdot T + x_{2i} \cdot T^2$  expressing lattice parameters *i* as function of temperature.

<sup>a</sup>from powder XRD, <sup>b</sup>from single-crystal XRD

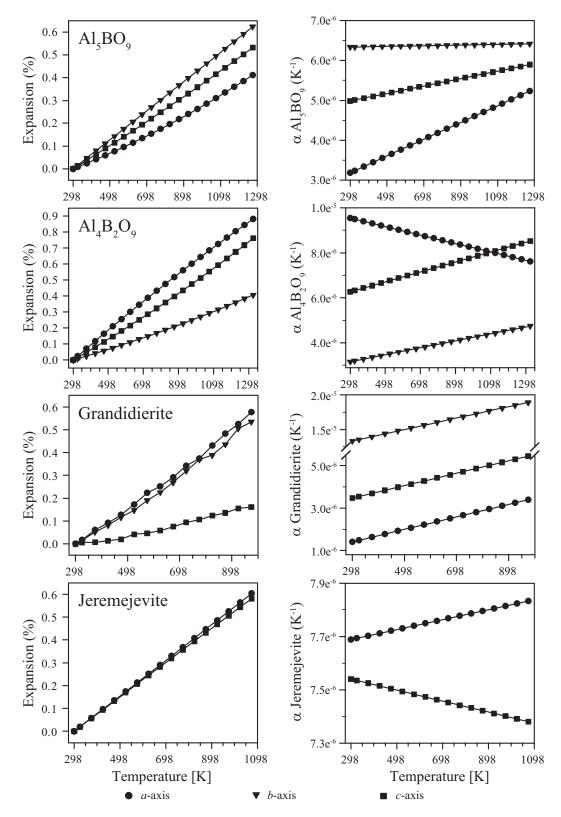


Figure 5.2: Normalized lattice parameters vs. temperature (left) and thermal expansion parameters  $\alpha$  as a function of temperature (right).

### 5.5.2 Thermal Evolution of Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>

Several ambiguous structure solutions for mullite-type  $Al_4B_2O_9$  exist. The most promising one is the monoclinic model (space group C2/m) of Fischer et al. (2008b) based on powder XRD, <sup>11</sup>B and <sup>27</sup>Al MAS-NMR data. This structure is related to the one of boralsilite (Peacor et al. 1999) but could not be completely solved from available data.

Three types of edge-connected chains of  $AlO_6$  octahedra extend parallel to the *b*-axis. The chains are linked among each other by three different polyhedral configurations. One of them is a complex arrangement of distorted  $AlO_5$  ( $AlO_{4+1}$ ),  $AlO_4$  polyhedra and  $BO_3$ triangles. The second type comprises  $AlO_{4+1}$  and  $BO_4$  tetrahedra which are disordered with  $AlO_4$  tetrahedra and  $BO_3$  triangles due to a partially occupied oxygen site. The third type contains  $BO_3$  and/or  $BO_4$  units. However, its exact configuration has not been resolved yet.

Lattice parameters as function of temperature and thermal expansion coefficients  $\alpha(T)$  are listed in Tables 5.2 and 5.3. The *a*-axis expands most, followed by the *c*- and the *b*-axis (the latter parallel to AlO<sub>6</sub> chains). The volumetric expansion is +2.06% and the monoclinic angle decreases from 90.8290(10)° at 298K to 90.7808(13)° at 1323K.  $\alpha_b$  of the shortest crystal axis increases most (in percent), followed by  $\alpha_c$  whereas  $\alpha_a$  decreases, meaning that the length-change of the *a*-axis diminishes with increasing temperature (Figure 5.2). The overall anisotropy of expansion decreases with temperature (*A* decreases from A<sub>298K</sub> = 12.6 to A<sub>1323K</sub> = 7.6). However, from room-temperature to about 1098K ( $\alpha_a = \alpha_c$ ), anisotropy of thermal expansion is most pronounced along the *a*-axis whereas at temperatures higher than 1098K, expansion is dominated along the *c*-axis.

Parameter	$a_0$	$a_1$	Parameter	$a_0$	$a_1$
$Al_5BO_9$ for 29	$98K \le T \le 12$	73K	$Al_4B_2O_9$ for 2	$298 \mathrm{K} \le T \le 13$	323K
$\alpha_a(T)$ (K <sup>-1</sup> )	$2.55439 \cdot 10^{-6}$	$2.10359 \cdot 10^{-9}$	$\alpha_a(T)$ (K <sup>-1</sup> )	$1.00906 \cdot 10^{-5}$	$-1.87160 \cdot 10^{-9}$
$\alpha_b(T)$ (K <sup>-1</sup> )	$6.31045 \cdot 10^{-6}$	$8.21191 \cdot 10^{-11}$	$\alpha_b(T)$ (K <sup>-1</sup> )	$2.70463 \cdot 10^{-6}$	$1.54378 \cdot 10^{-9}$
$\alpha_c(T)$ (K <sup>-1</sup> )	$4.70661 \cdot 10^{-6}$	$9.33647 \cdot 10^{-10}$	$\alpha_c(T)$ (K <sup>-1</sup> )	$5.61555 \cdot 10^{-6}$	$2.19838 \cdot 10^{-9}$
$\alpha_V(T)$ (K <sup>-1</sup> )	$1.357 \cdot 10^{-5}$	$3.114 \cdot 10^{-9}$	$\alpha_V(T)$ (K <sup>-1</sup> )	$1.842 \cdot 10^{-5}$	$1.870 \cdot 10^{-9}$
Grandidierite	for 298K $\leq~T$	$\leq 973 K$	Jeremejevite f	for 298K $\leq T$	$\leq 1073 \mathrm{K}$
$\alpha_a(T)$ (K <sup>-1</sup> )	$5.3387 \cdot 10^{-7}$	$2.9700 \cdot 10^{-9}$	$\alpha_a(T)$ (K <sup>-1</sup> )	$7.63280 \cdot 10^{-6}$	$1.86261 \cdot 10^{-10}$
$\alpha_b(T)$ (K <sup>-1</sup> )	$1.0853 \cdot 10^{-5}$	$8.2929 \cdot 10^{-9}$	$\alpha_c(T)$ (K <sup>-1</sup> )	$7.60138 \cdot 10^{-6}$	$-2.06187 \cdot 10^{-10}$
$\alpha_c(T)$ (K <sup>-1</sup> )	$2.5960 \cdot 10^{-6}$	$2.9337 \cdot 10^{-9}$	$\alpha_V(T)$ (K <sup>-1</sup> )	$2.286 \cdot 10^{-5}$	$1.730 \cdot 10^{-10}$
$\alpha_V(T)$ (K <sup>-1</sup> )	$1.09 \cdot 10^{-5}$	$1.27 \cdot 10^{-5}$			

**Table 5.3:** Coefficients for polynomials  $\alpha(T) = a_{0i} + a_{1i} \cdot T$  expressing thermal expansion as a function of temperature (*i* corresponds to lattice parameters).

### 5.5.3 Thermal Evolution of Grandidierite (Mg,Fe<sup>2+</sup>)Al<sub>3</sub>BSiO<sub>9</sub>

Grandidierite  $(Mg,Fe)Al_3BSiO_9$  is the Mg-rich member of the grandidierite-ominelite series (Dzikowski et al. 2007, Hiroi et al. 2001) and is of particular interest because it contains both Al and (Mg,Fe) in five-fold coordination (MacKenzie and Meinhold 1997, Dzikowski

et al. 2007, Farges 2001). Structure refinement of single-crystal XRD room-temperature data indicated an almost pure Mg endmember with 97.3(3)% Mg and 2.7(3)% Fe for the selected sample.

The grandidierite structure (space group Pnma) contains two different chains of edgeconnected AlO<sub>6</sub> octahedra parallel to the *b*-axis. The Al1O<sub>6</sub> chain is rather straight whereas chains of Al2O<sub>6</sub> octahedra are strongly folded (Figure 5.3). Parallel (100), adjacent Al octahedra are directly connected by corner-linked chains of  $-BO_3$ -AlO<sub>5</sub>-BO<sub>3</sub>- extending along *b*. Parallel (001), the chains of octahedra are joined by corner-linked chains of  $-SiO_4$ -MgO<sub>5</sub>-SiO<sub>4</sub>- extending parallel *b*. Along *b*, channels between adjacent chains of octahedra are either occupied by arrangements of SiO<sub>4</sub>, AlO<sub>5</sub> and MgO<sub>5</sub> polyhedra, or are locally empty if BO<sub>3</sub> polyhedra link octahedra parallel to (100). AlO<sub>5</sub> and MgO<sub>5</sub> polyhedra are regular bipyramids.

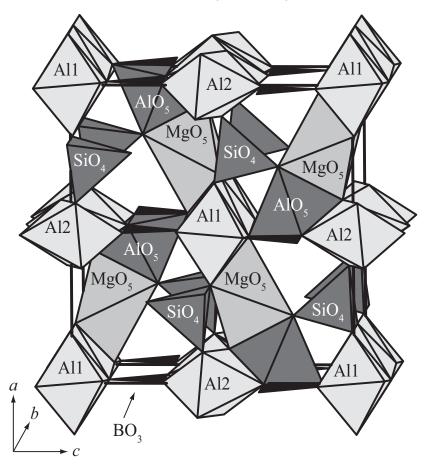


Figure 5.3: The crystal structure of grandidierite in space group Pnma.

Polynomials expressing temperature dependence of lattice parameters and  $\alpha(T)$  are shown in Tables 5.2 and 5.3. Between 298K and 973K, expansion of the *a*-axis is dominant, followed by the *b*-axis (parallel to the AlO<sub>6</sub> chains) and the *c*-axis. The volumetric expansion is +1.28%. Temperature dependence of expansion parameters is  $\alpha_a > \alpha_c > \alpha_b$  (Figure 5.2). Expansion of grandidierite is highly anisotropic. A increases from A<sub>298K</sub> = 23.8 to A<sub>873K</sub> = 31.0. This high anisotropy results from  $\alpha_b$  being a factor of ten larger than  $\alpha_a$  and  $\alpha_c$ . The increase in anisotropy is therefore due to the strong expansion along the *b*-axis.

Polyhedral bond lengths and average thermal expansion values obtained from the singlecrystal structural refinements are listed in Table 5.4. Average thermal expansion of the Al1O<sub>6</sub>, AlO<sub>5</sub> and SiO<sub>4</sub> polyhedra are in perfect agreement with values of Hazen and Prewitt (1977), Table 5.4. Values for Al2O<sub>6</sub> and BO<sub>3</sub> fit to predicted values within two standard deviations. The expansion of the MgO<sub>5</sub> polyhedron is too low ( $<\alpha > = 8.3(1.0) \cdot 10^{-6}$ K<sup>-1</sup> instead of 11.5  $\cdot 10^{-6}$ K<sup>-1</sup>).

Polyhedron	298K (Å)	873K (Å)	$< \alpha > (K^{-1})$
Al <sub>5</sub> BO <sub>9</sub> :			
AlO <sub>6</sub>	1.9004(9)	1.909(2)	$7.7(1.8) \cdot 10^{-6}$
$Al2O_5$	1.8642(11)	1.869(2)	$4.5(2.3) \cdot 10^{-6}$
$Al3O_5$	1.8822(12)	1.892(3)	$8.9(2.5) \cdot 10^{-6}$
AlO <sub>4</sub>	1.7472(11)	1.752(2)	$4.6(2.4) \cdot 10^{-6}$
$BO_3$	1.3769(14)	1.377(3)	$-0.7(3.8) \cdot 10^{-6}$
Grandidierite:			
Al1O <sub>6</sub>	1.8981(7)	1.9076(4)	$8.2(0.8) \cdot 10^{-6}$
Al2O <sub>6</sub>	1.9094(8)	1.9171(4)	$7.1(0.8) \cdot 10^{-6}$
AlO <sub>5</sub>	1.8548(11)	1.8604(6)	$5.3(1.2) \cdot 10^{-6}$
$MgO_5$	2.0311(19)	2.0407(6)	$8.3(1.0) \cdot 10^{-6}$
$\mathrm{SiO}_4$	1.6290(10)	1.6289(6)	$-0.2(1.2) \cdot 10^{-6}$
$BO_3$	1.3656(14)	1.3631(8)	$-3.2(2.0) \cdot 10^{-6}$
Jeremejevite:			
$AlO_5(F,OH)^a$	1.8861(8)	1.8961(16)	$9.2(2.0) \cdot 10^{-6}$
$B1O_3$	1.3855(13)	1.3851(40)	$-0.5(4.4) \cdot 10^{-6}$
$B2O_3$	1.3646(7)	1.3623(19)	$-2.9(2.6) \cdot 10^{-6}$
Expected values for $<\alpha>$ :			
$< \alpha > AlO_6:$	$8.23 \cdot 10^{-6} \mathrm{K}^{-1}$	$< \alpha > MgO_5:$	$11.5 \cdot 10^{-6} \mathrm{K}^{-1}$
$< \alpha > AlO_5$ :	$4.94 \cdot 10^{-6} \mathrm{K}^{-1}$	$< \alpha > BO_3$ :	$0 K^{-1}$
$<\alpha>\operatorname{AlO}_4$ :	0K <sup>-1</sup>	$<\alpha>$ SiO4:	0K <sup>-1</sup>

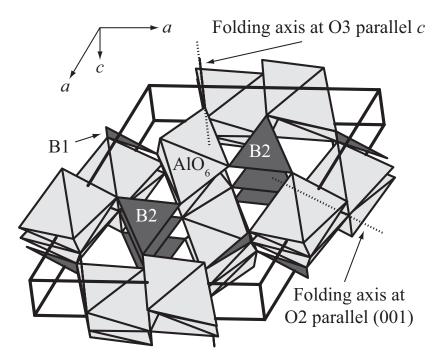
**Table 5.4:** Average polyhedral bond lengths at 298K and 873K, their average thermal expansion coefficient  $\langle \alpha \rangle$  and expected  $\langle \alpha \rangle$  values according to Hazen and Prewitt (1977).

<sup>a</sup>treated as  $AlO_6$ 

The most significant changes in inter-polyhedral angles occur at Al2-O5-Al2 (stretching from 100.48(5)° to 101.10(3)°) and at Al2-O4-Al2 (bending from 94.01(5)° to 93.40(3)°), which results from stretching of the folded chain of Al2O<sub>6</sub> octahedra. Further angular changes occur at Mg-O2-Al3 (124.32(6)° to 124.69(3)°), Mg-O5-Al2 (129.00(3)° to 128.73(1)°), Si1-O4-Al2 (132.91(2)° to 133.28(1)°) and at B1-O7-Al3 (130.82(8)° to 131.22(5)°). Except unfolding of the Al2O<sub>6</sub> chain, the largest structural changes are observed in the configuration of SiO<sub>4</sub>, AlO<sub>5</sub> and MgO<sub>5</sub> polyhedra within the channels along *b*. Expansion of this channel is limited to expansion of the AlO<sub>6</sub> octahedra at its corners because the channel walls are held together by non-expanding BO<sub>3</sub> units parallel to the (100) plane and by a face of the nonexpanding SiO<sub>4</sub> tetrahedra parallel to the (001) plane. MgO<sub>5</sub> has the lowest bond strength in the structure  $(z/\rho = 0.4)$  making it the most preferable buffer to compensate expansion of other polyhedra in the structure.

#### 5.5.4 Thermal Evolution of Jeremejevite Al<sub>6</sub>B<sub>5</sub>O<sub>15</sub>(F,OH)<sub>3</sub>

Jeremejevite  $Al_6B_5O_{15}(F,OH)_3$  was the first known aluminoborate mineral (Mallard 1887). The unit-cell volume of 518.13(4) Å<sup>3</sup> (Table 5.1) indicates that our sample is close to a F-endmember of jeremejevite, which is a common composition for natural jeremejevite (Foord et al. 1981).



**Figure 5.4:** The crystal structure of jeremejevite. Dashed lines indicate axes at which M-O-M angular changes are observed. B1 and B2 indicate  $B1O_3$  and  $B2O_3$  triangles, respectively.

Viewed along the *c*-axis, the jeremejevite structure (Figure 5.4) is characterized by two types of channels. An empty hexagonal channel is located at 0, 0, *c*, whereas the other is triangular and located at  $\frac{2}{3}$ ,  $\frac{1}{3}$ , *c* and  $\frac{1}{3}$ ,  $\frac{2}{3}$ , *c*. The triangular channels are enclosed by chains of edge sharing AlO<sub>5</sub>F octahedra. These chains are not straight as in mullite-type structures but consist of staggered arrangements of edge connected dimers of AlO<sub>5</sub>F octahedra expanding along *c*. B1O<sub>3</sub> triangles (face parallel to the *c*-direction) support the chains along *c*. B2O<sub>3</sub> units (face parallel to the (001) plane) within the triangular channel connect adjacent octahedra apexes at the same *c*-height. Along *c*, every third octahedral apex pointing towards the center of a triangular channel consists of F instead of O and is thus not connected to a BO<sub>3</sub> unit.

Polynomials expressing temperature dependent lattice parameters are shown in Table 5.2 and polynomials expressing  $\alpha(T)$  in Table 5.3. Between 298K and 1073K, all axes expand similarly. The volumetric expansion is +1.8%. At 298K  $\alpha_a$  is slightly larger than  $\alpha_c$  and with increasing temperature  $\alpha_a$  expands and  $\alpha_c$  decreases slightly (Figure 5.2) resulting in an increase of the anisotropy factor A from  $A_{298K} = 0.3$  to  $A_{1073K} = 0.9$ . Jeremejevite has the lowest A of all investigated structures. Expansion of BO<sub>3</sub> triangles and AlO<sub>5</sub>F (Table 5.4) is in good agreement with the expected values of Hazen and Prewitt (1977).

In the investigated temperature regime, structural changes in jeremejevite are controlled by inter-polyhedral folding at Al1-O3-Al1 from  $140.93(2)^{\circ}$  to  $141.11(6)^{\circ}$  and at Al1-O2-Al1 from  $124.19(2)^{\circ}$  to  $123.86(6)^{\circ}$  (Al1-O2-Al1 is the angle between two connected dimers of octahedra). Angular changes are a result of AlO<sub>6</sub> expansion, which is attenuated by nonexpanding BO<sub>3</sub> units connected to them. Larger expansion of AlO<sub>5</sub>F octahedra (compared to zero expansion of BO<sub>3</sub> units) can only be compensated by rotation. Dimers of octahedra are rotated at the edge joining two dimers (Figure 5.4), resulting in decrease of the Al1-O2-Al1 angle. In addition, at the linking corner of two adjacent octahedral chains the Al1-O3-Al1angle expands.

### 5.6 Discussion

Thermal behavior of the investigated boron-mullite compounds can be compared to those of the Al<sub>2</sub>SiO<sub>5</sub> polymorphs andalusite and sillimanite and mullite (Al<sub>2</sub>(Al<sub>2+2x</sub>Si<sub>2-2x</sub>)O<sub>10-x</sub>,  $0.18 \leq x \leq 0.88$ ). All three have mullite-type structures with AlO<sub>6</sub> chains, interlinked by single chains of alternating Si- and Al-tetrahedra (sillimanite), double-chains of AlO<sub>5</sub> and SiO<sub>4</sub> polyhedra (andalusite) and by disordered arrangements of Al and Si in mullite (Fischer and Schneider 2005).

Winter and Ghose (1979) investigated thermal expansion of sillimanite and andalusite. In sillimanite, from 298K to 1273K, lattice parameter expansion with temperature is b > c > a (c is parallel to the chains of octahedra) and A decreases from 11.3 to 10.8. The most distinctive expansion along the b-axis is due to anisotropic expansion of octahedra. Even though the mullite-type chain is held together by chains of rigid non-expanding SiO<sub>4</sub>/AlO<sub>4</sub> tetrahedra, octahedral rotation into the channels is responsible for expansion c > a. The expansion of andalusite is most expressed along the a-axis, followed by the b- and c-axis (mullite-type chains parallel to c). Average expansion of AlO<sub>6</sub> is slightly higher than expected ( $<\alpha > = 12.6(3.2) \cdot 10^{-6} \text{K}^{-1}$  instead of  $8.2 \cdot 10^{-6} \text{K}^{-1}$ ). Anisotropic expansion of AlO<sub>6</sub> units is has its strongest contribution to the a-axis whereas expansion along the c-axis is constrained by the double-chain of rigid AlO<sub>5</sub>/SiO<sub>4</sub> polyhedra. Due to the double-chains, no open channels exist and AlO<sub>6</sub> rotation is limited leading to the lowest expansion along the c-axis. Andalusite shows high and increasing anisotropy of thermal expansion ( $A_{298K} = 20.6$ ,  $A_{1273K} = 21.75$ ).

In 3:2 mullite, cell dimensions increase in the sequence b > c > a (mullite-type chains parallel to c). Anisotropy factors are rather low, A = 6.2 between 298K and 1173K (Schneider and Eberhard 1990) and A = 3.8 between 298K and 1273K (Brunauer et al. 2001b). Based on neutron powder diffraction data, structural changes in mullite with increasing temperature are due to rotation and increasing distortion of Al-octahedra and shifts in coordinates and varying Si and Al populations between mullite-type chains (Brunauer et al. 2001a).

In all investigated structures,  $BO_3$  units (and  $SiO_4$  in grandidierite) behave as rigid bodies and therefore act as rigid linkages (hinges) between other polyhedra, confirming the statement of Filatov and Bubnova (2008). Between 298K and 973K,  $Al_5BO_9$  has the lowest volumetric expansion, followed by grandidierite,  $Al_4B_2O_9$  and jeremejevite.

The smallest volumetric expansion of  $Al_5BO_9$  can be explained by its structural relation to sillimanite: Expansion of polyhedra is balanced by the empty channels parallel to the mullite-type chains. As one wall of the channel is formed by the O4-O5 edge of  $Al_{3}O_5$ , the channel is rather flexible. However, in contrast to sillimanite, mullite-type chains are only close to two channels instead of four, and therefore, rotation of octahedra is limited. To explain the smallest expansion behavior along the *a*-axis, the relation to the andalusite structure becomes relevant: Expansion along the *a*-axis is limited by the chain of  $-BO_3$ - $Al_{3}O_5$ - $BO_3$ consisting of rigid  $BO_3$  units (rigid  $AlO_4$  and  $SiO_4$  in andalusite) parallel to the mullite-type chains (Wada et al. 1993). Combined with limited rotation of octahedra, expansion along the *a*-axis (parallel to the mullite-type chains) has the strongest restrictions. For structures containing  $AlO_4$  and  $AlO_6$  polyhedra Hazen and Prewitt (1977) state that structural changes with increasing pressure correspond to those with decreasing temperature. Elastic behavior of  $Al_5BO_9$  was investigated by Gatta et al. (2010) and the behavior of lattice parameters with increasing pressure (compression is highest along the *b*-axis, followed by the *c*- and *a*-axis) is in agreement with our results for expansion of lattice parameters with temperature.

In grandidierite, all polyhedra except MgO<sub>5</sub> expand as expected according to Hazen and Prewitt (1977). The MgO<sub>5</sub> polyhedron buffers thermal expansion due to its low bond strength  $(z/\rho = 0.4)$ . Furthermore, slightly anisotropic expansion of the Al2-octahedra combined with unfolding of the corresponding chain is responsible for b > c expansion (mullite-type chains parallel b). All polyhedra in jeremejevite behave as rigid bodies. No flexible parts in the structure can compensate expansion and the structure reacts to rising temperature with folding and tilting of polyhedral arrangements.

The observed sequence of volumetric expansion may seem to contradict the conclusion of Filatov and Bubnova (2008) stating that thermal expansion of borates decreases with the number of cations with low bond strength (according to which grandidierite should undergo the largest expansion). It has to be considered that all of the structures investigated in this study are of similar chemical composition. In other words,  $AlO_4$  behaves not much different to  $SiO_4$ , as do  $MgO_5$  and  $AlO_5$ . Expansion properties of the investigated compounds are therefore mainly controlled by structural flexibility characteristic of each individual structure, i.e. their ability to buffer expansion by soft links between polyhedra rather than by differences in bond-strengths.

# 6 Thermal Stability of Jeremejevite

### 6.1 Introduction

Investigating strain related optical properties of crystals, Foord and Cunningham (1978) noticed that natural jeremejevite  $(Al_6B_5O_{15}(F,OH)_3, \text{very F-rich in nature, Foord et al. 1981})$  from Swakopmund, Namibia, changes from blue to an opaque white compound at 800°C.

In a study on synthesis, stability and breakdown products of the synthetic (OH)-endmember of jeremejevite, Stachowiak and Schreyer (1998) observed  $Al_4B_2O_9$ ,  $Al_6B_8O_{21}$ ,  $Al_3BO_6$  (Capponi et al. 1972), AlBO<sub>3</sub> (Capponi et al. 1972, Bither 1973, Vegas et al. 1977) and corundum as secondary phases in hydrothermal synthesis experiments.  $Al_4B_2O_9$  was also found as inclusions in (OH)-jeremejevite crystals. The type and amount of secondary phases depended on synthesis conditions (pressure, temperature).  $Al_6B_8O_{21}$  was previously reported in jeremejevite synthesis experiments by Reynaud (1977). Synthetic (OH)-jeremejevite remains stable up to fluid pressures of 50 kbar with a maximum temperature of ca. 750°C at 35 kbar (Schreyer and Werding 1997). Most interesting is the conclusion of Stachowiak and Schreyer (1998) that (OH)-jeremejevite breaks down to  $Al_4B_2O_9$  at ca. 625°C and 3 kbar. With increasing pressure,  $Al_4B_2O_9$  is reported to decompose to  $Al_3BO_6$  and  $AlBO_3$ . From extrapolated data of Stachowiak and Schreyer (1998), breakdown of the synthetic (OH)-endmember of jeremejevite can be predicted at ca. 500°C and ambient pressure:

$$2\operatorname{Al}_6B_5O_{15}(OH)_3 \xrightarrow{500^\circ C} 3\operatorname{Al}_4B_2O_9 + 2\operatorname{B}_2O_3 + 3\operatorname{H}_2O_3 + 3\operatorname{$$

During the investigation of thermal expansion of aluminoborates (Chapter 5), temperaturedependent jeremejevite lattice parameters were refined up to 800°C, however, data were collected up to 1100°C. Above ca. 800°C jeremejevite decomposed to an aluminoborate phase.

### 6.2 Experimental

Two natural blueish jeremejevite samples from Erongo region Namibia were investigated with powder X-ray diffraction and FTIR spectroscopy.

One sample (jeremejevite#1, characterized in Chapter 5) was ground, mixed with corundum as internal standard and subsequently measured with high-temperature powder X-ray diffraction up to 1100°C. After cooling to room temperature, the sample was re-measured with powder XRD and Fourier transformed infrared spectroscopy. The second sample (jeremejevite#2) was heated at 925°C for one hour and then heated at 950°C for two hours in air atmosphere, whereas another fragment of the same sample was heated at 1250°C. Both specimens turned white and were measured with powder X-ray diffraction and FTIR spectroscopy at room temperature.

Unit cell volumes of 518.13(4) Å<sup>3</sup> (#1, Chapter 5) and 518.91(6)Å<sup>3</sup> (#2) indicate that both jeremejevites are close to F-endmembers (Foord et al. 1981).

Powder XRD data were measured with a PANalytical X'Pert PRO MPD diffractometer (Cu radiation) equipped with an Anton-Paar HTK 1200 high-temperature goniometer attachment. Temperature-dependent data were acquired at 25°C, 50°C and then each 50°C up to 1100°C.

Diffraction patterns were collected from  $10^{\circ} 2\theta$  to  $70^{\circ} 2\theta$  at  $0.008^{\circ}$ /step with 40 s/step. Lattice parameters of the phases were derived from Rietveld refinements performed with Topas-Academic v. 4.1 (Coelho 2007a), using the structural models of Al<sub>5</sub>BO<sub>9</sub> (Chapter 4) Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> (Fischer et al. 2008b).

Fourier transformed infrared powder spectra were collected with a Perkin Elmer Spectrum One FTIR spectrometer equipped with a Perkin Elmer universal ATR accessory (Zn-Se + diamond crystal). The beam path included a mid-infrared light source, a KBr beam splitter and a TGS detector. Sample and background spectra were averaged from 50 scans measured  $550 \text{ cm}^{-1}$  to  $1500 \text{ cm}^{-1}$  at 4 cm<sup>-1</sup> resolution.

### 6.3 Results

A selection of temperature-dependent powder XRD patterns of jeremejevite#1 is shown in Figure 6.1. Between 25°C and 800°C, the pattern of jeremejevite does not significantly change, except of minor shifts due to thermal expansion (Section 5.5.4 in Chapter chap:HiTemp). At 850°C, the powder X-ray diffraction pattern is similar to  $Al_4B_2O_9$  (Fischer et al. 2008b). This phase retained stable up to ca. 1000°C when it transformed to a phase with an  $Al_5BO_9$ -like powder pattern (Chapter 3). Rietveld refined lattice parameters of in-situ high-temperature powder XRD experiment of jeremejevite#1 and of jeremejevite#2, measured at room temperature, are given in Table 6.1.

FTIR spectra of the jeremejevite#2 sample are shown in Figure 6.2. For comparison, spectra of  $Al_5BO_9$  and  $Al_4B_2O_9$  are also shown (measured with the same spectrometer with the same setup and settings).

### 6.4 Discussion

Unit cell dimensions of the phase measured at 950°C (jeremejevite#1) do not significantly differ from those of Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> measured at 950°C (a = 14.9178(3) Å, b = 5.56376(5) Å, c = 15.1724(3) Å, V = 1259.180(19) Å<sup>3</sup>,  $\beta = 90.7864(9)^{\circ}$ ) and those of the 1050°C phase can be compared to Al<sub>5</sub>BO<sub>9</sub> measured at 1000°C (a = 5.69221(3) Å, b = 15.10838(11) Å, 7.73652(5) Å, 665.342(8) Å<sup>3</sup>).

		· · ·					
	$T_{\rm M}$ (°C)	$T_{\rm H}$ (°C)	a-axis (Å)	b-axis (Å)	c-axis(Å)	Volume (Å)	$\beta$ (°)
#1	850	850	14.919(3)	5.5622(8)	15.143(4)	1256.5(4)	90.639(8)
#1	1050	1050	5.6832(4)	15.1077(15)	7.7285(8)	663.57(11)	
#1	25	1100	5.6626(4)	15.0060(13)	7.6855(7)	653.06(9)	
#2	25	850	14.820(8)	5.5498(17)	15.098(6)	1241.6(9)	90.938(8)
#2	25	1250	5.6651(3)	15.0147(9)	7.6891(5)	654.03(7)	

**Table 6.1:** Refined lattice parameters of jeremejevite in-situ measured by high-temperature powder X-ray diffraction (#1) and previously heated jeremejevite measured at room temperature (#2).

 $T_{\rm M}$  = temperature during measurement,  $T_{\rm H}$  maximum heating temperature of the sample

Room-temperature lattice parameters of the Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> and Al<sub>5</sub>BO<sub>9</sub> phases obtained by heating jeremejevite#2 at 950°C and 1100°C, respectively, are in the range of literature-data of Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> and Al<sub>5</sub>BO<sub>9</sub>. According to Fischer et al. (2008b), Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> lattice parameters are: a = 14.8056(7) Å, b = 5.5413(2) Å, c = 15.0531(6) Å, V = 1234.83 Å<sup>3</sup>,  $\beta = 90.913(2)^{\circ}$ . Al<sub>5</sub>BO<sub>9</sub> lattice parameters (Chapter 4) are a = 5.6686(2) Å, b = 15.0060(9) Å, c = 7.6892(7)Å, V = 654.07(6) Å<sup>3</sup>).

FTIR spectra of jeremejevite#2 heated at 950°C and  $Al_4B_2O_9$  are very similar, however, peaks are better resolved in the  $Al_4B_2O_9$  reference spectrum. The only difference in the spectra of jeremejevite#2 heated at 1250°C and  $Al_5BO_9$  can be found at ca. 850 cm<sup>-1</sup>. According to MacKenzie (2005), bands in the region of 830 cm<sup>-1</sup> to 909 cm<sup>-1</sup> are ascribed to  $AlO_4$ -stretch vibrations.

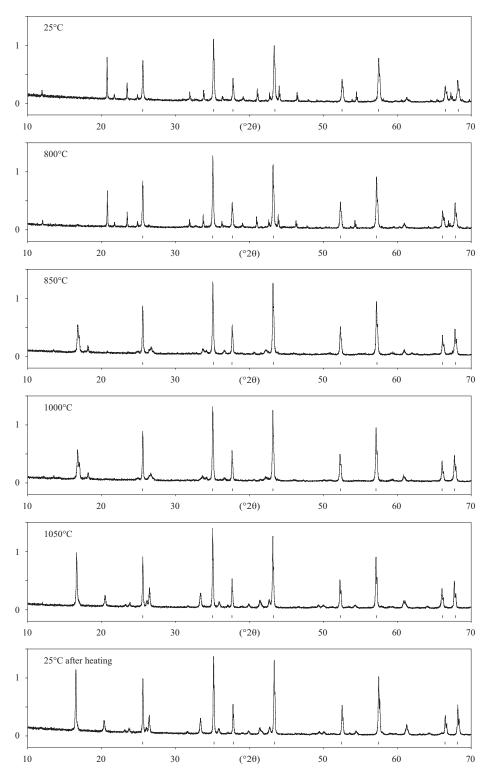
In-situ high-temperature powder X-ray diffraction experiments indicate that natural jeremejevite decomposes at temperatures between 800°C and 850°C. Agreement of lattice parameters and corresponding FTIR spectra yield  $Al_4B_2O_9$  as product. For an F endmember of jeremejevite, the reaction of  $Al_4B_2O_9$  formation from jeremejevite at high temperature can be described as:

$$2\operatorname{Al}_6B_5\operatorname{O}_{15}F_3 \xrightarrow{800^\circ\mathrm{C}-850^\circ\mathrm{C}} 3\operatorname{Al}_4B_2\operatorname{O}_9 + B_2\operatorname{O}_3 + 2\operatorname{BF}_3\uparrow$$

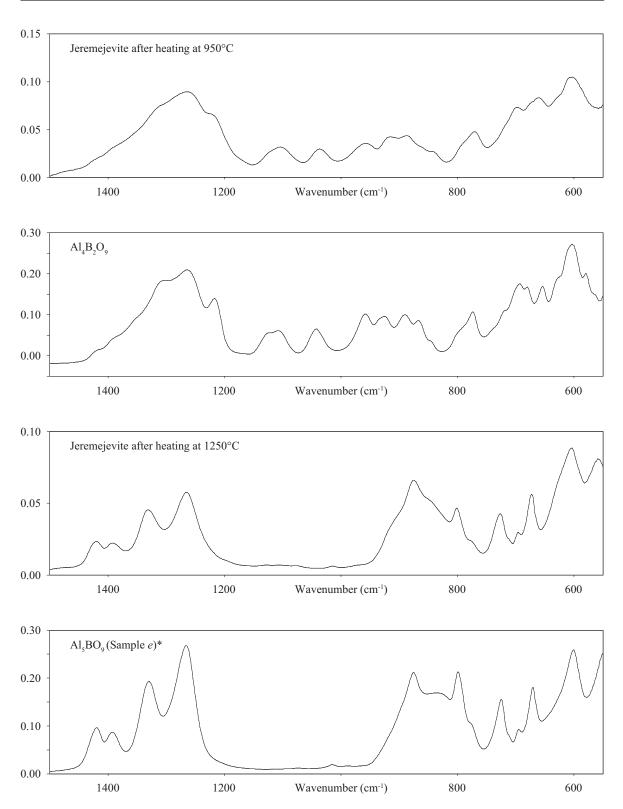
 $BF_3$  is a volatile gas even at room temperature and will escape. At the formation conditions of  $Al_4B_2O_9$  (> 850°C),  $B_2O_3$  may either be present as melt or evaporate. A minor concentration of residual  $B_2O_3$  may cause the differences in FTIR spectra between reference  $Al_5BO_9$ and  $Al_4B_2O_9$  compounds and those produced by decomposition of jeremejevite. Residual  $B_2O_3$  hydrates under ambient conditions to X-ray amorphous  $2H_3BO_3$  and can therefore not be detected in diffraction patterns.

Formation of  $Al_5BO_9$  from  $Al_4B_2O_9$  at ca. 1050 °C is already known and has been extensively discussed in literature (Chapter 3, Scholze 1956, Gielisse and Forster 1962, Ray 1992, Mazza et al. 1992, Duoy 2005 and others).

The significant difference in breakdown temperature between natural (ca. 800°C–850°C) and synthetic (OH)-jeremejevite (ca. 500°C, Stachowiak and Schreyer 1998) can thus be ascribed to the chemical difference (F or OH) of the compounds.



**Figure 6.1:** Temperature-dependent powder XRD pattern of jeremejevite mixed with ca. 70 wt.-% corundum as internal standard. The pattern measured at 25°C shows diffraction peaks of jeremejevite and corundum (*hkl*-ticks). At 800°C, jeremejevite is still stable. Diffraction peaks of the pattern at 850°C fit the structure of  $Al_4B_2O_9$  according to Fischer et al. (2008b). Peaks of the pattern measured at 1050°C and at 25°C after heating fit the  $Al_5BO_9$  structure. Intensity scale is counts  $\cdot 10^3$ .



**Figure 6.2:** Top to bottom: FTIR spectra of jeremejevite after heating at 950°C,  $Al_4B_2O_9$  (prepared according to Fischer et al. 2008b), jeremejevite measured after heating at 1250°C and an FTIR spectrum of  $Al_5BO_9$  (sample *e* in Chapter 4). Peaks between 1450 cm<sup>-1</sup> and 1250 cm<sup>-1</sup> are assigned to  $BO_3$  vibrations, whereas peaks in the region of 1200 cm<sup>-1</sup> to 950 cm<sup>-1</sup> are assigned to  $BO_4$  stretching vibrations. Intensity scale is ATR absorbance.

# 7 Stability at High Pressure, Elastic Behavior and Pressure-Induced Structural Evolution of "Al₅BO<sub>9</sub>", a Mullite-Type Ceramic Material

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# 7.1 Abstract

Elastic behavior and pressure-induced structural evolution of synthetic boron-mullite "Al<sub>5</sub>BO<sub>9</sub>" (a = 5.678(2) Å, b = 15.015(4) Å and c = 7.700(3) Å, space group  $Cmc2_1, Z = 4$ ) were investigated up to 7.4 GPa by in-situ single-crystal X-ray diffraction with a diamond anvil cell under hydrostatic conditions. No phase-transition or anomalous compressional behavior occurred within the investigated P-range. Fitting the P-V data with a truncated secondorder (in energy) Birch-Murnaghan Equation-of-State (BM-EoS), using the data weighted by the uncertainties in P and V, we obtained:  $V_0 = 656.4(3)$  Å<sup>3</sup> and  $K_{\rm T0} = 165(7)$  GPa ( $\beta_{\rm V0}$ = 0.0061(3) GPa<sup>-1</sup>). The evolution of the Eulerian finite strain vs. normalized stress ( $f_{\rm E}$ - $F_{\rm E}$ plot) leads to an almost horizontal trend, showing that a truncated second-order BM-EoS is appropriate to describe the elastic behavior of " $Al_5BO_9$ " within the investigated *P*-range. The weighted linear regression through the data points gives:  $F_{\rm E}(0) = 159(11)$  GPa. Axial compressibility coefficients yielded:  $\beta_a = 1.4(2) \cdot 10^{-3}$  GPa<sup>-1</sup>,  $\beta_b = 3.4(4) \cdot 10^{-3}$  GPa<sup>-1</sup>, and  $\beta_c = 1.7(3) \cdot 10^{-3}$  GPa<sup>-1</sup> ( $\beta_a:\beta_b:\beta_c = 1:2.43:1.21$ ). The highest compressibilities observed in this study within (100) can be ascribed to the presence of voids represented by 5-membered rings of polyhedra: Al1-Al3-Al4-Al1-Al3, which allow accommodating the effect of pressure by polyhedral tilting. Polyhedral tilting around the voids also explains the higher compressibility along [010] than along [001]. The stiffer crystallographic direction observed here might be controlled by the infinite chains of edge-sharing octahedra running along [100], which act as "pillars", making the structure less compressible along the a-axis than along the b- and caxis. Along [100], compression can only be accommodated by deformation of the edge-sharing octahedra (and/or by compression of the Al-O bond lengths), as no polyhedral tilting can occur. In addition, a comparative elastic analysis among the mullite-type materials is carried out.

# 7.2 Introduction

Compounds with chemical composition belonging to the ternary system  $Al_2O_3-B_2O_3-SiO_2$  are a class of ceramic materials investigated in particular for industrial applications. Their hightemperature stability, low-thermal expansion coupled with high creep resistance, low electric conductivity, high chemical stability and low-density are leading to an increasing number of outstanding applications of such materials: As construction and engineering ceramics, refractory linings due to their high resistance to corrosion, optically translucent ceramics for high-temperature furnace windows, fire-protecting lining in nuclear plants due to the low density and the capability of absorbing neutrons, substrates for catalytic convertors, and electronic devices (Fischer and Schneider 2005; 2008a, Li and Chang 2006, Peng et al. 2006, Tang et al. 2006, Zhang et al. 2006, Tao et al. 2007, Wei et al. 2007). Among those, mullite  $(Al_2(Al_{2+2x}Si_{2-2x})O_{10-x}$  with  $\approx 0.2 < x < \approx 0.9$  and  $B_2O_3 = 0$ ) (Saalfeld and Guse 1981, Angel et al. 1991, Schneider and Komarneni 2005) is doubtless of great importance. Several studies have been devoted to the stability fields, solubility and crystal-chemistry of the compounds belonging to the ternary system  $Al_2O_3$ - $B_2O_3$ -SiO\_2 (Baumann and Moore 1942, Letort 1952, Dietzel and Scholze 1955, Scholze 1956, Gielisse and Forster 1962, Kim and Hummel 1962, Capponi et al. 1972, Reynaud 1977, Sokolova et al. 1978, Werding and Schreyer 1984; 1992; 1996, Rymon-Lipinski et al. 1985, Mazza et al. 1992, Peacor et al. 1999, Fischer and Schneider 2005; 2008a, Buick et al. 2008, Fischer et al. 2008b, Grew et al. 2008, Griesser et al. 2008). In their recent review paper on crystal-chemistry of boroaluminosilicates with mullite-type structures, Fischer and Schneider (2008a) described the structural relationship in the  $Al_2O_3$ - $B_2O_3$ -SiO\_2 compounds, showing that all boron-mullites can be derived from a hypothetical aristotype with topological symmetry P4/mbm, in which the main building block is represented by chains of edge-sharing MO<sub>6</sub> octahedra running along [001].

Thermo-elastic behavior and pressure (P) and temperature (T) structural evolution of aluminosilicates along the joint  $Al_2O_3$ -SiO<sub>2</sub> have been extensively investigated by Brillouin spectroscopy (Vaughan and Weidner 1978) and by in-situ X-ray/neutron single-crystal/powder diffraction (Brace et al. 1969, Winter and Ghose 1979, Ralph et al. 1984, Schneider and Eberhard 1990, Comodi et al. 1997, Yang et al. 1997a;b, Brunauer et al. 2001b, Friedrich et al. 2004, Burt et al. 2006, Gatta et al. 2006b). The P/T-induced main deformation mechanisms were described on the basis of in-situ HP/HT-structure refinements. For  $2 Al_2O_3$ :SiO<sub>2</sub> mullite, a full description of the elastic properties (i.e. stiffness coefficients ( $c_{ij}$ ), elastic compliances ( $s_{ij}$ ), Young's moduli ( $E_{ij}$ ), Poisson's ratios ( $v_{ij}$ ), bulk modulus (K) and shear modulus (G)) was reported by Hildmann et al. (2001) by single-crystal resonant ultrasound spectroscopy (RUS); Temperature derivatives of the elastic constants between -170°C and 1400°C were later derived by Schreuer et al. (2006) by single-crystal RUS. The elastic constants and their T-derivatives of 2.5Al<sub>2</sub>O<sub>3</sub>:SiO<sub>2</sub> mullite have been reported by Kriven et al. (1999) and Palko et al. (2002) by Brillouin spectroscopy. In contrast, no elastic data are available for boroaluminate or boroaluminosilicates.

Aim of the present study was to investigate the elastic behavior and *P*-induced structural evolution of a boroaluminate by means of in-situ single-crystal X-ray diffraction, with a diamond anvil cell, in order to elucidate the role of boron on the elastic features of mullite-type materials. We selected the "Al<sub>5</sub>BO<sub>9</sub>" compound for this first H*P*-experiment. Two slightly different compositions of the investigated compound are present in literature: Baumann and Moore (1942), Ihara et al. (1980) and Garsche et al. (1991) interpreted the material as  $9Al_2O_3:2B_2O_3$  (Al<sub>4.91</sub>B<sub>1.09</sub>O<sub>9</sub>, Z = 4, commonly cited as  $Al_{18}B_4O_{33}$ ), whereas Sokolova et al. (1978) concluded the compound to consist of  $10Al_2O_3:2B_2O_3$  (Al<sub>5</sub>BO<sub>9</sub>, Z = 4). In  $Al_{18}B_4O_{33}$ , 9% of aluminum tetrahedra are substituted by four-fold coordinated boron. In  $Al_5BO_9$ , boron is only present on a three-fold coordinated site. In this study, test single-crystal X-ray refinements with boron constraint to the three-fold coordinated site were superior to refinements with excess boron according to  $Al_{4.91}B_{1.09}O_9$ . For this reason, the investigated compound is here referred to as "Al<sub>5</sub>BO<sub>9</sub>".

Sokolova et al. (1978) described the structure in space group  $Cmc2_1$  (a = 5.6673(7) Å, b = 15.011(2) Å, c = 7.693(1) Å), consisting of mullite-type octahedral chains (in this set running along [100]), linked by edge-sharing AlO<sub>5</sub> bipyramids alternating with AlO<sub>4</sub> tetrahedra and BO<sub>3</sub> triangular units (Figure 7.1). In their classification of mullite-type materials, Fischer and Schneider (2008a) assigned this compound to the "MUL-VIII.33,  $Bb2_1m$ : A9B2" group, together with Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub>, Al<sub>9</sub>BSi<sub>2</sub>O<sub>19</sub> (boromullite) and Al<sub>16.6</sub>Cr<sub>1.4</sub>B<sub>4</sub>O<sub>33</sub>. Due to the structural homologies with mullite, the authors considered the terms "boron-mullite" and "B-mullite" as appropriate for "Al<sub>5</sub>BO<sub>9</sub>". Garsche et al. (1991) showed that Cr<sup>3+</sup> can replace Al<sup>3+</sup> at the octahedral site in Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> up to about 10 wt.-% Cr<sub>2</sub>O<sub>3</sub>. A further Al<sub>5</sub>BO<sub>9</sub> structural model was reported by Mazza et al. (1992), with a pseudo-tetragonal unit-cell ( $a \approx b \approx 7.6$  Å and c  $\approx 2.8$  Å) in space group *Pbam*. This compound is supposed to be metastable and transforms to equilibrium phase upon heating Fischer and Schneider (2008a).

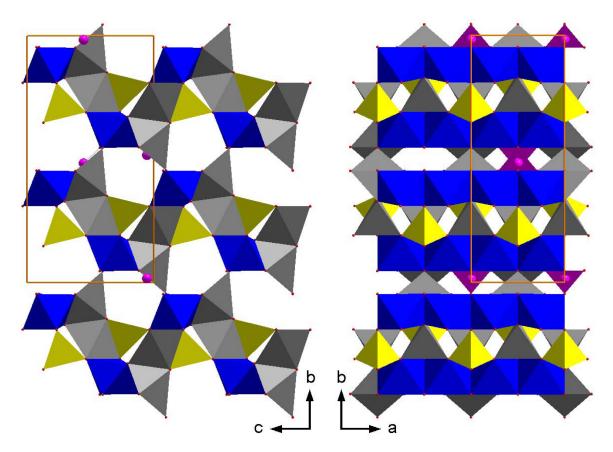


Figure 7.1: The crystal-structure of "Al<sub>5</sub>BO<sub>9</sub>" viewed down [100] (left) and down [001] (right) in  $Cmc2_1$ . Octahedra are represented in blue, distorted bipyramidal polyhedra in gray, tetrahedra in yellow and triangular BO<sub>3</sub> units in purple (purple spheres represent the B-sites). The edge-sharing octahedral chains running along [100] are well evident (right), as well as the distorted 5-membered rings of polyhedra forming the voids (left).

### 7.3 Experimental Methods

"Al<sub>5</sub>BO<sub>9</sub>" crystals were synthesized by slow cooling of a starting mixture in a flux consisting of 0.50 g K<sub>2</sub>CO<sub>3</sub> and 1.56 g MoO<sub>3</sub> corresponding to  $1 \text{ K}_2\text{CO}_3 + 3 \text{ MoO}_3 \rightarrow \text{ K}_2\text{Mo}_3\text{O}_{10} + \text{CO}_2\uparrow$ . The starting mixture was composed of Al<sub>2</sub>O<sub>3</sub> and B<sub>2</sub>O<sub>3</sub> mixed in a molar ratio of 9:2 (0.03 g B<sub>2</sub>O<sub>3</sub> and 0.19 g Al<sub>2</sub>O<sub>3</sub>). All compounds were ground for 10 min in an agate mortar and placed in a lid-covered platinum crucible. The mixture was heated to 800°C at a gradient of 100°C/h and then further heated to 1100°C at a rate of 50°C/h. After 2h at 1100°C, the melt was slowly cooled down to 600°C at 10°C/h.

Separation of the whitish, half-transparent crystals from yellowish flux was done by soaking the melt in hot water. Yield was about 0.1 g of radially grown aggregates of needle-like crystals with a maximum size of about  $0.2 \cdot 0.2 \cdot 0.8$  mm<sup>3</sup>.

One platy crystal  $(180.140.60 \ \mu m^3)$ , free of defects or twinning under the transmitting polarized light microscope, was selected for X-ray diffraction experiments. Diffraction data were first collected at room conditions with an Oxford Diffraction-Xcalibur diffractometer equipped with CCD, using graphite monochromatized Mo $K\alpha$ -radiation, operated at 50 kV and 40 mA. A combination of  $\omega/\phi$  scans was used in order to maximize the reciprocal space coverage and redundancy, with a scan width of 0.4° and an exposure time of 30 s/frame (Table 7.1). The distance between the crystal and the detector was 80 mm. 7265 Bragg reflections were collected in the range 2 < 2 $\theta$  < 70° 2 $\theta$ , of which 1302 were unique and 926 with I > $4\sigma(I)$  (Table 7.1). The diffraction pattern was fully indexed with an orthorhombic lattice with a = 5.678(2) Å, b = 15.015(4) Å and c = 7.700(3) Å, in agreement with the experimental findings of Sokolova et al. (1978) for  $Al_5BO_9$ . Integrated intensities were then corrected for Lorentz-polarization (Lp) and for absorption effects (by Gaussian integration based upon the shape and dimensions of the crystal), using the CrysAlis package (Oxford Diffraction Ltd. 2005). After corrections, the discrepancy factor among symmetry-related reflections (Laue class mmm) was R(int) = 0.074 (Table 7.1) and reflection conditions were consistent with space group  $Cmc_{21}$ , as reported by Sokolova et al. (1978). The anisotropic structural refinement was then performed using the SHELX-97 software (Sheldrick 1997), starting from atomic coordinates of Sokolova et al. (1978). The refined Flack parameter (Sheldrick 1997) was 0 within  $1\sigma(x)$ . Neutral atomic scattering factors of B, Al and O from the International Tables for Crystallography (Wilson and Prince 1999) were used. No peak larger than +0.82/-0.61 $e^{-}/Å^{3}$  was present in the final difference-Fourier synthesis and the variance-covariance matrix showed no significant correlation between refined parameters. Further details pertaining to the structural refinement at ambient conditions are reported in Tables 7.1 and 7.2–7.13.

An ETH-type diamond anvil cell (DAC) (Miletich et al. 2000) was used to perform the in-situ high-pressure experiment. 250  $\mu$ m thick T301 steel foil was used as gasket, which was pre-indented to a thickness of about 110  $\mu$ m before drilling a 300  $\mu$ m hole by spark-erosion. The crystal of "Al<sub>5</sub>BO<sub>9</sub>" already measured at ambient conditions was placed into the gasket hole together with a single-crystal of quartz for pressure calibration (Angel et al. 1997).

	Table 7	Table 7.1: Details pertaining to the data collections and refinements of "Al <sub>5</sub> BO <sub>9</sub> " at different pressures	rtaining to th	e data collect	tions and refi	nements of "/	$Al_5BO_9$ " at di	ifferent pressu	ires.	
P (GPa)	0.0001	$P_0: 0.0001^{a}$	$P_1: 0.15(5)$	$P_2: 0.91(4)$	$P_3: 1.99(6)$	$P_4: 3.32(6)$	$P_5: 4.78(6)$	$P_6: 5.83(5)$	$P_7: 5.99(5)$	$P_8: 6.45(6)$
Diffractometer					Xcalib	Xcalibur CCD				
X-ray radiation	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$	$MoK\alpha$
Scan type	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$	$\phi/\omega$
Scan width ( $^{\circ}/\mathrm{fr.}$ )	0.4	0.4	0.4	0.4	0.4	0.4	0.4	0.4	0.4	0.4
Exposure $(s/fr.)$	30	30	30	30	30	30	30	30	30	30
Temperature (K)	298	298	298	298	298	298	298	298	298	298
Space Group	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$	$Cmc2_1$
a-axis (Å)	5.678(2)	5.680(3)	5.677(3)	5.668(4)	5.662(4)	5.655(4)	5.648(4)	5.639(4)	5.641(2)	5.639(4)
b-axis (Å)	15.015(4)	15.02(60)	15.020(50)	14.990(30)	14.940(30)	14.860(30)	14.790(20)	14.760(40)	14.760(30)	14.730(30)
c-axis (Å)	7.700(3)	7.698(2)	7.694(4)	7.683(6)	7.669(6)	7.652(7)	7.644(7)	7.635(17)	7.626(5)	7.618(7)
Volume	656.5(4)	657(3)	656(2)	652.8(15)	648.7(15)	643.0(15)	638.5(11)	635(2)	635.0(14)	632.8(15)
Z	4	4	4	4	4	4	4	4	4	4
Maximum 2 $\theta$ (°)	70.02	77.99	78.02	78.40	78.02	77.07	77.60	75.61	78.09	77.37
Measured refl.	7265	1672	1421	1570	1580	1659	1641	1613	1609	1279
Unique refl.	1302	411	503	603	585	600	555	538	534	463
Refl. $> 4\sigma(I)$	926	311	350	361	357	401	402	383	385	307
$R(\mathrm{int})$	0.0739	0.1011	0.1007	0.1315	0.1163	0.1260	0.1131	0.1199	0.0952	0.1009
L.S. parameters	82	31	30	30	30	30	30	30	30	30
$R1,I>4 {\rm s}(I)$	0.0444	0.0753	0.0845	0.0866	0.0919	0.1098	0.1089	0.1020	0.1182	0.1187
${ m w}R2$	0.0821	0.1109	0.1181	0.1270	0.1268	0.1417	0.1305	0.1313	0.1437	0.1320
GooF	1.001	1.274	1.084	1.008	1.041	1.220	1.252	1.215	1.367	1.223
<sup>a</sup> with the crystal in the DAC without <i>P</i> -medium $R(\text{int}) = \sum  F_{\text{obs}}^2 - \langle F_{\text{obs}}^2 \rangle  / \sum (F_{\text{obs}}^2), R1 = \sum ( F_{\text{obs}}  -  F_{\text{calc}} ) / \sum  F_{\text{obs}} , wR2 = \sqrt{\sum (w(F_{\text{obs}}^2 - F_{\text{calc}}^2)^2) / \sum (w(F_{\text{obs}}^2)^2)},$	the DAC v $\langle F_{\rm obs}^2 \rangle   / \Sigma$	vithout <i>P</i> -media $(F_{\rm obs}^2), R1 = \sum$	um $  F_{\rm obs}  -  F_{\rm calc}  $	$ \rangle/\sum F_{\rm obs} , v$	$vR2 = \sqrt{\sum(w)}$	$(F_{\rm obs}^2 - F_{\rm calc}^2)^2$	$(w(F_{\rm obs}^2))/\sum (w(F_{\rm obs}^2))$	$\frac{2}{2}$		
$w = 1/(\sigma^2(F_{ m obs}^2) + (0.02P)^2), \ P = ({ m Max}(F_{ m obs}^2,0) + 2F_{ m calc}^2)/3$	$(0.02P)^2), i$	$^{\scriptscriptstyle D} = ({ m Max}(F_{ m obs}^2,$	$(0) + 2F_{ m calc}^2)/3$							

 $\overline{94}$ 

A methanol:ethanol mixture (4:1) was used as hydrostatic pressure-transmitting medium (Angel et al. 2007). Intensity data collections at 0.0001 GPa (crystal in DAC without any pressure medium,  $P_0$ ), 0.15(5) ( $P_1$ ), 0.91(5) ( $P_2$ ), 1.99(6) ( $P_3$ ), 3.32(6) ( $P_4$ ), 4.78(6) ( $P_5$ ), 5.83(5) (P<sub>6</sub>), 5.99(5) (P<sub>7</sub>) and 6.45(6) (P<sub>8</sub>) GPa (Table 7.1) were performed adopting the same experimental set-up and data collection protocol used with the crystal in air (Table 7.1). At any given pressure, integrated intensity data were corrected for Lp and absorption effects due to the crystal and the DAC using the ABSORB 5.2 computer program (Burnham 1966, Angel 2002). No violation of reflection conditions dictated by space group  $Cmc2_1$  was observed within the investigated *P*-range. The structure refinements were conducted using soft geometrical restraints aimed to restrain Al-O and B-O distances to those obtained at room-pressure in air, with an estimated standard deviation of  $\pm 0.04$  Å. This improved the stability of the HP-refinements, as soft restrains act as if they were additional experimental observations (Sheldrick 1997, Gatta et al. 2006a; 2008). In order to reduce the number of refined variables, isotropic displacement parameters were refined by grouping all of the Al-sites and all of the O-sites. Refined atomic positions and displacement parameters are reported in Tables 7.2–7.13. Bond distances and angles are listed in Appendix A.3. At 7.4 GPa, the gasket hole collapsed and the "Al<sub>5</sub>BO<sub>9</sub>" crystal was incidentally broken. Unit-cell constants measured at ambient conditions after decompression using a small fragment of the "Al<sub>5</sub>BO<sub>9</sub>" crystal  $(60.40.30 \ \mu m)$  recovered from the gasket hole showed that *P*-induced structural changes up to 7.4 GPa are completely reversible.

### 7.4 Elastic Behavior

The monotonic variation of " $Al_5BO_9$ " unit-cell parameters with pressure is shown in Figure 7.2. No evidence of phase-transition or change in the compressional behavior were observed within the investigated pressure-range. The elastic behavior of " $Al_5BO_9$ " is here described with a truncated second-order Birch-Murnaghan Equation-of-State (II-BM-EoS) (Birch 1947). This EoS is based upon the assumption that the high-pressure strain energy in a solid can be expressed as a Taylor series in the Eulerian finite strain, defined as

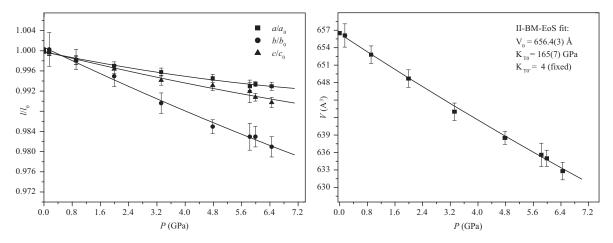
$$f_{\mathrm{E}} = rac{1}{2} \cdot \sqrt[3]{\left(rac{V_0}{V}
ight)^2} - 1$$

where  $V_0$  and V represent the unit-cell volume at ambient and HP conditions, respectively. Expansion in the Eulerian strain polynomial has the following form:

$$P(f_{\rm E}) = 3 \cdot K_{\rm T0} \cdot f_{\rm E} \cdot \sqrt[5]{(1+2 \cdot f_{\rm E})^2} \cdot \left\{ 1 + \frac{3}{2} \cdot (K'-4) \cdot f_{\rm E} + \frac{3}{2} \cdot \left[ K_{\rm T0} \cdot K'' + (K'-4) \cdot (K'-3) + \frac{35}{9} \right] \cdot f_{\rm E}^2 + \dots \right\}$$

where  $K_{T0}$  represents the bulk modulus  $(K_{T0} = -V_0(\delta P/\delta V)_{P=0} = 1/\beta V_0$ , where  $\beta V_0$  is the volume compressibility coefficient at ambient conditions), K' and K'' represent its pressure

derivatives  $(K' = \delta K_0 / \delta P; K'' = \delta^2 K_0 / \delta P^2)$ . Fitting the *P*-*V* data with a truncated secondorder (in energy) BM-EoS with the EOS-FIT 5.2 program (Angel 2001), using the data weighted by the uncertainties in *P* and *V*, we obtain:  $V_0 = 656.4(3)$  Å<sup>3</sup>,  $K_{T0} = 165(7)$  GPa  $(\beta_{V0} = 0.0061(3)$  GPa<sup>-1</sup>), and K' = 4 (fixed). A fitting with a third-order BM-EoS leads to a strong correlation between the refined parameters (especially  $K_{T0}$  and K'), with worse fitting statistic parameters.



**Figure 7.2:** Evolution of the unit-cell parameters of " $Al_5BO_9$ " with pressure. For the unit-cell lengths, the solid lines represent polynomial regression curves through the data points. For the unit-cell volume, the second-order Birch-Murnaghan Equation-of-State fit is shown.

The evolution of Eulerian finite strain vs. normalized stress  $(F_{\rm E} = P/[3f_{\rm E}\sqrt[5]{(1+2f_{\rm E})^2}])$ (Angel 2000) is shown in Figure 7.3. The weighted linear regression through the data points leads to:  $F_{\rm E}(0) = 159(11)$  GPa. The almost horizontal regression function confirms that the use of a truncated second-order BM-EoS is appropriate to describe the elastic behavior of "Al<sub>5</sub>BO<sub>9</sub>" within the *P*-range investigated. The value of the normalized stress extrapolated at ambient conditions ( $F_{\rm E}(0)$ ) and the bulk modulus ( $K_{\rm T0}$ ) refined on the basis of the BM-EoS are in good agreement.

The axial compressibility coefficients  $(\beta_j = 1/l_{0j} \cdot (\delta l_j / \delta P))$ , where  $l_{0j}$  (j = a, b, c) is the length of the unit-cell edges under room conditions) were calculated by weighted polynomial regressions trough the data points, yielding to:

$$egin{aligned} a/a_0 &= 1-0.0014(2){\cdot}P + 6(3){\cdot}10^{-5}{\cdot}P^2 \ (R^2 = 0.9906), \ b/b_0 &= 1-0.0034(4){\cdot}P + 5(6){\cdot}10^{-5}{\cdot}P^2 \ (R^2 = 0.6647), \ c/c_0 &= 1-0.0017(3){\cdot}P + 4(5){\cdot}10^{-5}{\cdot}P^2 \ (R^2 = 0.9855), \end{aligned}$$

with  $\beta_a = 1.4(2) \cdot 10^{-3}$  GPa<sup>-1</sup>,  $\beta_b = 3.4(4) \cdot 10^{-3}$  GPa<sup>-1</sup>,  $\beta_c = 1.7(3) \cdot 10^{-3}$  GPa<sup>-1</sup> ( $\beta_a:\beta_b:\beta_c = 1:2.43:1.21$ ).

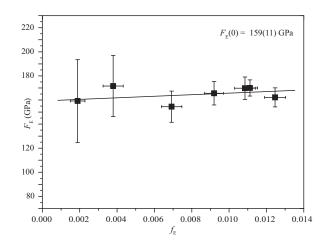


Figure 7.3: Eulerian finite  $(f_E)$  strain vs. normalized stress  $(F_E)$  plot. The weighted linear regression through the data points is shown.

### 7.5 Structural Evolution with Pressure

The single-crystal structural refinement of  ${}^{\circ}Al_{5}BO_{0}$  at room conditions confirms the structure model previously reported by Sokolova et al. (1978). Among the boroaluminosilicates, the topology of the " $Al_5BO_9$ " structure is unique, although several structural homologies can be found with the mullite and mullite-type materials (e.g. sillimanite). In particular, the main building units are represented by edge-sharing octahedral chains linked by edge-sharing  $AlO_5$  bipyramids alternating with  $AlO_4$  tetrahedra, forming small cavities that host boron coordinated by three framework oxygens ( $BO_3$  triangular units), as shown in Figure 7.1. In the structural model of Sokolova et al. (1978) used in this study (i.e. space group  $Cmc2_1$ with  $a \approx 5.67$  Å,  $b \approx 15.01$  Å,  $c \approx 7.69$  Å), octahedral chains run parallel to [100]. Atomic positions, bond-distances and angles refined in this study at room conditions agree with those previously reported (Tables 7.2–7.13). Bond-distances and angles show that octahedra, forming the [100]-chains, and the tetrahedra are not regular (octahedron:  $\langle Al1-O \rangle_{oct} = 1.903$  Å and  $\Delta$ (Al1-O)<sub>max</sub> = 0.116 Å; Tetrahedron:  $\langle$ Al4-O $\rangle$ <sub>tet</sub> = 1.753 Å and  $\Delta$ (Al4-O)<sub>max</sub> = 0.041 Å). The two bipyramidal units are strongly distorted (i.e.  $\Delta$ (Al2-O)<sub>max</sub> = 0.418 Å and  $\Delta$ (Al3- $O_{\text{max}} = 0.516 \text{ Å}$ ). The BO<sub>3</sub> triangular unit is almost regular, with B-O<sub>max</sub> = 1.382(4) Å and B-O<sub>min</sub> = 1.372(6) Å (i.e.  $\Delta$ (B-O)<sub>max</sub> = 0.010 Å).

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
Al1	0.0001	0.2519(2)	0.11653(6)	0.3336(2)	0.0069(2)
$8b^{c}$	$P_0^{a}$	0.2513(4)	0.1155(8)	0.3337(4)	0.0072(4)
	$P_1$	0.2521(3)	0.1171(6)	0.3343(4)	0.0079(4)
	$P_2$	0.2521(4)	0.1167(4)	0.3304(5)	0.0082(4)
	$P_3$	0.2513(4)	0.1164(4)	0.3326(6)	0.0086(4)
	$P_4$	0.2528(5)	0.1174(4)	0.3325(5)	0.0090(5)
	$P_5$	0.2520(5)	0.1171(4)	0.3328(6)	0.0086(4)
	$P_6$	0.2518(5)	0.1171(4)	0.3333(6)	0.0081(4)
	$P_7$	0.2524(6)	0.1172(5)	0.3325(7)	0.0084(5)
	$P_8$	0.2517(7)	0.1160(6)	0.3326(7)	0.0077(5)

Table 7.2: Atomic positions and isotropic displacement parameters of Al1 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Tab	le 7	.3:	Atom	ic p	positions	and	isotropic	disp	lacement	parameters	of	Al2	$\operatorname{at}$	different	pressure.	
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Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Ų)
Al2	0.0001	0.0	0.25515(9)	0.5199(2)	0.0096(4)
$4a^{c}$	$P_0^{a}$	0.0	0.2590(8)	0.5208(5)	0.0072(4)
	$P_1$	0.0	0.2547(6)	0.5216(6)	0.0079(4)
	$P_2$	0.0	0.2552(5)	0.5216(6)	0.0082(4)
	$P_3$	0.0	0.2561(4)	0.5198(6)	0.0086(4)
	$P_4$	0.0	0.2556(5)	0.5207(7)	0.0090(5)
	$P_5$	0.0	0.2546(5)	0.5214(7)	0.0086(4)
	$P_6$	0.0	0.2553(5)	0.5200(7)	0.0081(4)
	$P_7$	0.0	0.2554(5)	0.5210(8)	0.0084(5)
	$P_8$	0.0	0.2551(6)	0.5199(8)	0.0077(5)
a		a	1:		

<sup>a</sup>with the crystal in the DAC without pressure medium

 $^{b}U_{eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

$ \begin{array}{ c c c c c c c c } \hline {\rm Site} & P({\rm GPa}) & x & y & z & U_{\rm eq.}/U_{\rm iso}{}^{\rm b}({\rm \AA}^2) \\ \hline {\rm Al3} & 0.0001 & 0.0 & 0.44410(9) & 0.5156(2) & 0.0075(3) \\ \hline {\rm 4a}^{\rm c} & P_0{}^{\rm a} & 0.0 & 0.44418(8) & 0.5151(5) & 0.0072(4) \\ \hline & P_1 & 0.0 & 0.4456(6) & 0.5140(5) & 0.0079(4) \\ \hline & P_2 & 0.0 & 0.4457(5) & 0.5134(6) & 0.0082(4) \\ \hline & P_3 & 0.0 & 0.4457(5) & 0.5152(6) & 0.0086(4) \\ \hline & P_4 & 0.0 & 0.4453(5) & 0.5188(7) & 0.0090(5) \\ \hline & P_5 & 0.0 & 0.4445(5) & 0.5162(8) & 0.0086(4) \\ \hline & P_6 & 0.0 & 0.4443(5) & 0.5170(8) & 0.0081(4) \\ \hline & P_7 & 0.0 & 0.4458(6) & 0.5189(8) & 0.0084(5) \\ \hline & P_8 & 0.0 & 0.4465(6) & 0.5202(9) & 0.0077(5) \\ \hline \end{array} $		-		1 1	-	1
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Ų)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	Al3	0.0001	0.0	0.44410(9)	0.5156(2)	0.0075(3)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$4a^{c}$	$P_0^{a}$	0.0	0.4444(8)	0.5151(5)	0.0072(4)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		$P_1$	0.0	0.4456(6)	0.5140(5)	0.0079(4)
$\begin{array}{c ccccc} P_4 & 0.0 & 0.4453(5) & 0.5188(7) & 0.0090(5) \\ P_5 & 0.0 & 0.4445(5) & 0.5162(8) & 0.0086(4) \\ P_6 & 0.0 & 0.4443(5) & 0.5170(8) & 0.0081(4) \\ P_7 & 0.0 & 0.4458(6) & 0.5189(8) & 0.0084(5) \\ \end{array}$		$P_2$	0.0	0.4457(5)	0.5134(6)	0.0082(4)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		$P_3$	0.0	0.4457(5)	0.5152(6)	0.0086(4)
$\begin{array}{cccc} P_6 & 0.0 & 0.4443(5) & 0.5170(8) & 0.0081(4) \\ P_7 & 0.0 & 0.4458(6) & 0.5189(8) & 0.0084(5) \end{array}$		$P_4$	0.0	0.4453(5)	0.5188(7)	0.0090(5)
$P_7$ 0.0 0.4458(6) 0.5189(8) 0.0084(5)		$P_5$	0.0	0.4445(5)		0.0086(4)
		$P_{6}$	0.0	0.4443(5)	0.5170(8)	0.0081(4)
$P_8$ 0.0 0.4465(6) 0.5202(9) 0.0077(5)		$P_7$	0.0		. ,	
		$P_8$	0.0	0.4465(6)	0.5202(9)	0.0077(5)

Table 7.4: Atomic positions and isotropic displacement parameters of Al3 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements

<sup>c</sup>Wyckoff position

Site	P (GPa)	x	y	Z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
Al4	0.0001	0.0	0.29682(9)	0.1677(2)	0.0084(4)
$4a^{c}$	$P_0^{\mathbf{a}}$	0.0	0.2966(7)	0.1676(5)	0.0072(4)
	$P_1$	0.0	0.2966(6)	0.1675(5)	0.0079(4)
	$P_2$	0.0	0.2954(5)	0.1667(6)	0.0082(4)
	$P_3$	0.0	0.2980(5)	0.1677(6)	0.0086(4)
	$P_4$	0.0	0.2965(5)	0.1677(6)	0.0090(5)
	$P_5$	0.0	0.2969(5)	0.1667(7)	0.0086(4)
	$P_6$	0.0	0.2973(5)	0.1658(7)	0.0081(4)
	$P_7$	0.0	0.2976(6)	0.1658(8)	0.0084(5)
	$P_8$	0.0	0.2979(7)	0.1677(8)	0.0077(5)

Table 7.5: Atomic positions and isotropic displacement parameters of Al4 at different pressure.

<sup>a</sup> with the crystal in the DAC without pressure medium

<sup>b</sup>  $U_{eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Table 7.6: Atomic positions and isotropic displacement parameters of B1 at different pressure.

Site	P (GPa)	x	y	z	$U_{\rm eq.}/U_{\rm iso}{}^{\rm b}$ (Ų)
B1	0.0001	0.0	0.0164(4)	0.0542(7)	0.0112(14)
$4a^{c}$	$P_0^{\mathbf{a}}$	0.0	0.0201(17)	0.0549(13)	0.004(3)
	$P_1$	0.0	0.0176(15)	0.0542(13)	0.0079(4)
	$P_2$	0.0	0.0193(14)	0.0515(16)	0.0082(4)
	$P_3$	0.0	0.0209(14)	0.0482(17)	0.0086(4)
	$P_4$	0.0	0.0166(15)	0.0513(18)	0.0090(5)
	$P_5$	0.0	0.0177(15)	0.0555(19)	0.0086(4)
	$P_{6}$	0.0	0.0176(15)	0.0529(19)	0.0081(4)
	$P_7$	0.0	0.0186(17)	0.052(2)	0.0084(5)
	$P_8$	0.0	0.0173(17)	0.048(2)	0.0077(5)
a		a			

<sup>a</sup>with the crystal in the DAC without pressure medium

<sup>b</sup>  $U_{eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Ų)
01	0.0001	0.0	0.3285(2)	0.7049(5)	0.0090(8)
$4a^{c}$	$P_0^{a}$	0.0	0.3289(16)	0.700(1)	0.0082(6)
	$P_1$	0.0	0.3266(12)	0.7027(10)	0.0085(6)
	$P_2$	0.0	0.3262(10)	0.7016(11)	0.0068(7)
	$P_3$	0.0	0.3290(9)	0.6994(12)	0.0082(7)
	$P_4$	0.0	0.3265(10)	0.7008(14)	0.0083(8)
	$P_5$	0.0	0.3285(10)	0.7049(14)	0.0079(7)
	$P_{6}$	0.0	0.3308(10)	0.7014(14)	0.0067(7)
	$P_7$	0.0	0.3254(12)	0.7026(16)	0.0065(8)
	$P_8$	0.0	0.3238(13)	0.6990(16)	0.0036(8)

Table 7.7: Atomic positions and isotropic displacement parameters of O1 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
O2	0.0001	0.2558(4)	0.19016(15)	0.5345(3)	0.0073(6)
$8b^{c}$	$P_0^{a}$	0.2564(7)	0.1914(13)	0.5330(7)	0.0082(6)
	$P_1$	0.2566(8)	0.1882(10)	0.5345(7)	0.0085(6)
	$P_2$	0.2566(9)	0.1918(7)	0.5339(7)	0.0068(7)
	$P_3$	0.2567(9)	0.1908(7)	0.5335(8)	0.0082(7)
	$P_4$	0.2590(11)	0.1895(7)	0.5343(8)	0.0083(8)
	$P_5$	0.2566(10)	0.1905(7)	0.5308(9)	0.0079(7)
	$P_6$	0.2549(10)	0.1903(7)	0.5311(9)	0.0067(7)
	$P_7$	0.2582(12)	0.1891(8)	0.5337(9)	0.0065(8)
	$P_8$	0.2578(12)	0.1901(9)	0.5321(9)	0.0036(8)

Table 7.8: Atomic positions and isotropic displacement parameters of O2 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Table 7.9: Atomic	positions and	isotropic d	isplacement	parameters of O3 a	at different pressure.

	-			-	-
Site	P (GPa)	x	y	z	$U_{\rm eq.}/U_{\rm iso}{}^{\rm b}$ (Ų)
O3	0.0001	0.0	0.4541(2)	0.9045(5)	0.0086(8)
$4a^{c}$	$P_0^{\mathbf{a}}$	0.0	0.4535(15)	0.9073(10)	0.0082(6)
	$P_1$	0.0	0.4534(11)	0.9060(9)	0.0085(6)
	$P_2$	0.0	0.454(1)	0.9086(13)	0.0068(7)
	$P_3$	0.0	0.4554(8)	0.9062(14)	0.0082(7)
	$P_4$	0.0	0.4554(9)	0.9033(15)	0.0083(8)
	$P_5$	0.0	0.4536(9)	0.9100(16)	0.0079(7)
	$P_6$	0.0	0.4550(9)	0.9101(17)	0.0067(7)
	$P_7$	0.0	0.4561(11)	0.9008(17)	0.0065(8)
	$P_8$	0.0	0.4553(12)	0.9021(18)	0.0036(8)
a	. 1: .1 D.	a	11		

<sup>a</sup>with the crystal in the DAC without pressure medium

 $^{b}U_{eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

$\begin{array}{c c c c c c c c c c c c c c c c c c c $		-			-	-
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	O4	0.0001	0.0	0.0465(2)	0.4250(5)	0.0096(8)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$4a^{c}$	$P_0^{a}$	0.0	0.0434(14)	0.4261(11)	0.0082(6)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		$P_1$	0.0	0.0457(10)	0.4245(10)	0.0085(6)
$\begin{array}{c ccccc} P_4 & 0.0 & 0.0468(8) & 0.4194(16) & 0.0083(8) \\ P_5 & 0.0 & 0.0445(9) & 0.4232(16) & 0.0079(7) \\ P_6 & 0.0 & 0.0463(8) & 0.4246(16) & 0.0067(7) \\ P_7 & 0.0 & 0.0458(10) & 0.4204(18) & 0.0065(8) \\ \end{array}$		$P_2$	0.0	0.0477(8)	0.4260(13)	0.0068(7)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		$P_3$	0.0	0.0465(8)	0.4243(13)	0.0082(7)
$\begin{array}{cccc} P_6 & 0.0 & 0.0463(8) & 0.4246(16) & 0.0067(7) \\ P_7 & 0.0 & 0.0458(10) & 0.4204(18) & 0.0065(8) \end{array}$		$P_4$	0.0	0.0468(8)	0.4194(16)	0.0083(8)
$P_7$ 0.0 0.0458(10) 0.4204(18) 0.0065(8)		$P_5$	0.0	0.0445(9)	0.4232(16)	0.0079(7)
		$P_{6}$	0.0	0.0463(8)	0.4246(16)	0.0067(7)
$P_8$ 0.0 0.0480(11) 0.4180(19) 0.0036(8)		$P_7$	0.0			
		$P_8$	0.0	0.0480(11)	0.4180(19)	0.0036(8)

Table 7.10: Atomic positions and isotropic displacement parameters of O4 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements

<sup>c</sup>Wyckoff position

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
O5	0.0001	0.0	0.3522(2)	0.3716(5)	0.0090(8)
$4a^{c}$	$P_0^{a}$	0.0	0.348(2)	0.3721(9)	0.0082(6)
	$P_1$	0.0	0.3475(14)	0.3706(8)	0.0085(6)
	$P_2$	0.0	0.3503(11)	0.3738(10)	0.0068(7)
	$P_3$	0.0	0.3504(10)	0.3722(10)	0.0082(7)
	$P_4$	0.0	0.3516(11)	0.3692(11)	0.0083(8)
	$P_5$	0.0	0.3521(12)	0.3672(12)	0.0079(7)
	$P_{6}$	0.0	0.3534(12)	0.3712(12)	0.0067(7)
	$P_7$	0.0	0.3505(13)	0.3663(13)	0.0065(8)
	$P_8$	0.0	0.3530(16)	0.3698(13)	0.0036(8)

Table 7.11: Atomic positions and isotropic displacement parameters of O5 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Table 7.12: Atomic positions and isotropic displacement parameters of O6 at different pressure.

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Å <sup>2</sup> )
O6	0.0001	0.0	0.1911(2)	0.2671(5)	0.0076(8)
$4a^{c}$	$P_0^{a}$	0.0	0.1908(17)	0.2687(10)	0.0082(6)
	$P_1$	0.0	0.1912(13)	0.2668(10)	0.0085(6)
	$P_2$	0.0	0.1933(10)	0.2667(11)	0.0068(7)
	$P_3$	0.0	0.1903(10)	0.2654(11)	0.0082(7)
	$P_4$	0.0	0.1898(11)	0.2626(13)	0.0083(8)
	$P_5$	0.0	0.1907(11)	0.2674(13)	0.0079(7)
	$P_{6}$	0.0	0.1904(11)	0.2658(13)	0.0067(7)
	$P_7$	0.0	0.1910(12)	0.2641(15)	0.0065(8)
	$P_8$	0.0	0.1933(13)	0.2628(16)	0.0036(8)

<sup>a</sup>with the crystal in the DAC without pressure medium

<sup>b</sup>  $U_{eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

Site	P (GPa)	x	y	z	$U_{ m eq.}/U_{ m iso}{}^{ m b}$ (Ų)
07	0.0001	0.2110(5)	0.04838(16)	0.1183(3)	0.0097(6)
$8b^{c}$	$P_0^{a}$	0.2113(9)	0.0540(13)	0.1187(6)	0.0082(6)
	$P_1$	0.2105(8)	0.050(1)	0.1193(6)	0.0085(6)
	$P_2$	0.2106(8)	0.0491(8)	0.1201(7)	0.0068(7)
	$P_3$	0.2130(8)	0.0479(7)	0.1192(9)	0.0082(7)
	$P_4$	0.2112(9)	0.0472(7)	0.1216(8)	0.0083(8)
	$P_5$	0.2124(9)	0.0494(7)	0.1172(10)	0.0079(7)
	$P_{6}$	0.2124(10)	0.0489(7)	0.1188(10)	0.0067(7)
	$P_7$	0.2119(10)	0.0479(8)	0.1220(9)	0.0065(8)
	$P_8$	0.2117(10)	0.0478(10)	0.1223(9)	0.0036(8)

Table 7.13: Atomic positions and isotropic displacement parameters of O7 at different pressure.

<sup>a</sup>with the crystal in the DAC without pressure medium

 ${}^{\rm b}U_{\rm eq.}$  from anisotropic refinements with the crystal in air (0.0001 GPa),  $U_{\rm iso}$  for isotropic refinements <sup>c</sup>Wyckoff position

The refinements at HP-conditions show only minor change of the "Al<sub>5</sub>BO<sub>9</sub>" structure (Tables 7.2–7.13 and Appendix A.3). The variation of the polyhedral bond-distances and angles is not larger than  $2(\sigma)$  within the *P*-range investigated. In other words, between 0.0001 and 6.5 GPa, polyhedra behave as rigid units. A corresponding behavior is also found for the BO<sub>3</sub>unit (Appendix A.3). Polyhedral distortions at 6.5 GPa are comparable to those observed at room-conditions (at 6.45 GPa:  $\Delta$ (Al1-O)<sub>max</sub> = 0.12 Å,  $\Delta$ (Al2-O)<sub>max</sub> = 0.46 Å,  $\Delta$ (Al3-O)<sub>max</sub> = 0.56 Å,  $\Delta$ (Al4-O)<sub>max</sub> = 0.04 Å and  $\Delta$ (B-O)<sub>max</sub> = 0.01 Å). However, inter-polyhedral angles and distances vary significantly with pressure, which are interpreted as the main compressional mechanisms governing the unit-cell elastic anisotropy. In particular, if the structure is viewed down [100], voids with pentagonal shape, formed by 5-membered rings of polyhedra (i.e. Al1-Al3-Al4-Al1-Al3, Figure 7.4) occur. The decrease of the O7-O5, O7-O2 and O3-O2 distances between 0.0001 and 6.45 GPa (-5.2%, -4.6% and -2.7%, respectively) is significantly larger than the O3-O3 and O3-O5 ones (-1.6% and -1.1%) (Figures 7.4 and 7.5). The P-induced compression along O7-O5 and O7-O2 contributes more to the shortening of the b-axis, whereas the P-induced decrease of O3-O3 and O3-O2 distances controls compression along the c-axis (Figure 7.4). Shortening of the O7-O5, O7-O2, O3-O2, O3-O3 and O3-O5 distances is ascribable to polyhedral tilting, rather than an intra-polyhedral compression, as shown by the evolution of the O-O-O angles inscribed in the pentagonal distorted rings. In particular, the O2-O3-O7 angle decreases from  $85.8(2)^{\circ}$  to  $82.5(2)^{\circ}$  (i.e.  $\Delta$ (O2-O3-O7)  $\approx 3.6^{\circ}$ ) and O7-O3-O5 from  $88.2(4)^{\circ}$  to  $87.0(3)^{\circ}$  (i.e.  $\Delta$ (O7-O3-O5)  $\approx 1.2^{\circ}$ ), whereas the variation of O3-O5-O2, O5-O2-O3 and O3-O7-O3 are  $< 1^{\circ}$ . These P-induced mechanisms lead to a preferred compression along [010].

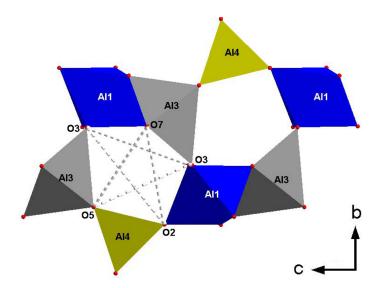


Figure 7.4: Configuration of the voids represented by 5-membered rings of polyhedra (i.e. Al1-Al3-Al4-Al1-Al3). The "diameters" of the voids (i.e. O7-O5, O7-O2, O3-O3, O3-O2 and O3-O5) are shown as dashed lines.

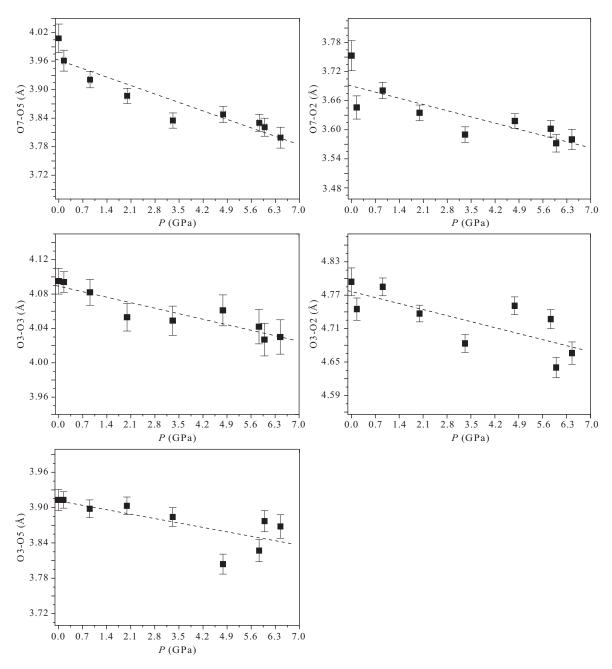


Figure 7.5: Evolution of the O7-O5, O7-O2, O3-O3, O3-O2 and O3-O5 distances (i.e. "diameters" of the voids represented by the 5-membered ring of polyhedra). Dashed lines represent the weighted linear regressions through the data points.

### 7.6 Discussion and Conclusions

The present in-situ high-pressure experiment shows that " $Al_5BO_9$ " remains crystalline at least up to 7.4 GPa, and any *P*-induced structural change is completely reversible. No phase-transition or anomalous elastic behavior have been observed within the *P*-range investigated, proving the high stability in elastic regime of this ceramic material at least up to 7.4 GPa.

The elastic analysis shows that "Al<sub>5</sub>BO<sub>9</sub>" is a stiff material. Its isothermal bulk modulus is similar to the one of mullite (adiabatic bulk modulus of 2:1-mullite:  $K_{\rm S} = 169.1$  GPa, Schreuer et al. 2006;  $K_{\rm S} = 2$ :1-mullite: 169.2 GPa, Hildmann et al. 2001;  $K_{\rm S} = 2.5$ :1-mullite: 173.6.0 GPa, Palko et al. 2002;  $K_{\rm S} = 2.5$ :1-mullite: 171.6 GPa, Kriven et al. 1999; Values recalculated by Schreuer et al. 2006 as average of Voigt and Reuss model) and sillimanite ( $K_{\rm S} = 171.4$ GPa, Vaughan and Weidner 1978;  $K_{T0} = 171(1)$  GPa, Yang et al. 1997b, and  $K_{T0} = 164(1)$ GPa, Burt et al. 2006). The elastic anisotropy of mullite-type materials observed in previous experiments is strikingly high (e.g. 2:1-mullite:  $(s_{11}+s_{12}+s_{13}):(s_{21}+s_{22}+s_{23}):(s_{31}+s_{32}+s_{33})$  $=\beta_a:\beta_b:\beta_c = 1.40:1.86:1$ , Hildmann et al. 2001; Sillimanite:  $\beta_a:\beta_b:\beta_c = 1.22:1.63:1$ , based on single-crystal unit-cell parameters measured between 0.0001 and 5.3 GPa by Yang et al. 1997b, and  $(3K_{\rm a})^{-1}:(3K_{\rm b})^{-1}:(3K_{\rm c})^{-1} = \beta_{\rm a}:\beta_{\rm b}:\beta_{\rm c} = 1.82:2.63:1$ , based on single-crystal unit-cell parameters measured between 0.0001 and 8.5 GPa by Burt et al. 2006). All the aforementioned studies showed that mullite-type materials are stiffer along the c-axis than along the a- or b-axis. For a comparative elastic analysis of the mullite-type materials extended to "Al<sub>5</sub>BO<sub>9</sub>", we have to consider an origin shift by  $\mathbf{t} = (0, 0, \frac{1}{3})$  and a transformation matrix by  $\mathbf{T}$ = (**b**, **c**, **a**), according to Fischer and Schneider (2008a), from  $Cmc2_1$  setting used in this study and in Sokolova et al. (1978) to  $Bb2_1m$  setting. In this light, there is a full agreement between the stiffer direction found in this study (i.e. [100]) and that found in mullite-type materials previously investigated (i.e. [001]). The higher compressibilities observed in this study within (100) can be ascribed to the presence of voids, which allow accommodating the effect of pressure by polyhedral tilting. Polyhedral tilting around the aforementioned 5membered rings (Al1-Al3-Al4-Al1-Al3) also explains the higher compressibility along [010] than along [001]. The stiffer crystallographic direction observed here might be controlled by the infinite chains of edge-sharing octahedra running along [100], which act as "pillars", making the structure less compressible along the a-axis than along the b- and c-axis. The reason is that along [100] compression can be accommodated only by deformation of the edge-sharing octahedra (and/or by compression of the Al-O bond lengths), as no polyhedral tilting can occur. In response to the applied pressure, any structure is expected to react first by tilting the polyhedra, then by distorting the polyhedra, and finally by decreasing the bond distances. This hierarchy is due to the fact that the first mechanism is energetically less costly than the other two (Gatta 2010).

Following the comparative analysis on the elastic behavior of mullite-type materials, we observed a similar elastic behavior between " $Al_5BO_9$ " and mullite(s) or sillimanite along the direction of the octahedral chains, which is the less compressible direction, but a different behavior occurs in the plane perpendicular to the chains. Considering the aforementioned metrical relationship with other mullite-type materials, we expected that " $Al_5BO_9$ " were more compressible along [001] than along [010]. However, we believe that this different elastic behavior on (100) is due to the different structural configuration on (100) in " $Al_5BO_9$ " compared to other mullite-type materials. In particular, the presence of distorted bipyramids in

"Al<sub>5</sub>BO<sub>9</sub>", which act as bridging-units of the [100]-octahedral chains and are not present in mullite or sillimanite, might explain the different elastic behavior on (100).

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## Note

This text represents the contents of the article "Stability at high pressure, elastic behavior and pressure-induced structural evolution of " $Al_5BO_9$ ", a mullite-type ceramic material", as published in Physics and Chemistry of Minerals 37(4) (2010) 227-236. The unit cell volume has been added to Table 7.1 and some passages have been slightly modified to fit the layout of this thesis.

# 8 Solid Solution Experiments

#### 8.1 Introduction

Since Dietzel and Scholze (1955) reported compounds with mullite-like properties consisting of  $B_2O_3$ ,  $SiO_2$  and  $Al_2O_3$ , solid solutions in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$  were assumed. The chemical variability of natural boromullite samples investigated by Buick et al. (2008) is close to compositions of synthetic samples of Dietzel and Scholze (1955) and may thus support this assumption.

 $Al_{6-x}B_xO_9$  solid solutions were investigated by Mazza et al. (1992)  $(1 \le x \le 3)$  and MacKenzie et al. (2007)  $(1 \le x \le 4)$ . Both reported decreasing lattice parameters with increasing x. From FTIR spectra, Mazza et al. (1992) showed that four-fold coordinated boron is incorporated into the mullite structure (for x = 2) and MacKenzie et al. (2007) identified  $BO_4$  in boron-mullites with  $x \ge 1.2$  by <sup>11</sup>B MAS NMR spectroscopy.

Recently, Griesser et al. (2008) investigated boron incorporation into the mullite structure in samples synthesized at 950°C and 1300°C. Starting materials were of compositions between 3:2 mullite and  $Al_6B_4O_{15}$  and between 7:3 mullite and a hypothetical compound with 70:30 mol-%  $Al_2O_3$ :  $B_2O_3$ .  $Al_2O_3$  ratio was kept constant and  $B_2O_3$  and SiO<sub>2</sub> were varied according to a hypothetical  $B \rightarrow Si$  substitution. Samples with less than 20 mol-%  $B_2O_3$  were identified as Al-Si-B-mullites (powder XRD) and lattice parameters decreased with increasing boron content. The compositional ranges of starting materials resulting in samples identified as Al-Si-B-mullites (based on lattice parameters) are indicated with the two thick gray lines in Figure 8.1. 950°C samples with  $B_2O_3$  contents above 15–20 mol-% were reported as  $Al_4B_2O_9$ (Mazza et al. 1992) with excess SiO<sub>2</sub>. Samples with the same composition but heated at 1300°C were identified as  $Al_{18}B_4O_{33}$  (Scholze 1956, Garsche et al. 1991) with excess SiO<sub>2</sub>. BO<sub>4</sub> incorporation in  $Al_4B_2O_9$  samples was confirmed by FTIR spectroscopy.

In a recent study, Zhang et al. (2010) investigated boron doped (max. 6.14 wt.-%  $B_2O_3$ ) mullites. Starting from 3:2 mullite, the Al/Si ratio was kept constant at 3/1 and the B/Al ratio was varied from 0–0.4/3 (shown with the thin gray line in Figure 8.1). Samples heated at 1000°C, are reported to consist of mullite,  $Al_4B_2O_9$  (Mazza et al. 1992) and a "spinel phase". At 1400°C,  $Al_{18}B_4O_{33}$  and mullite were identified with X-ray powder diffraction. With increasing  $B_2O_3$  content, Zhang et al. (2010) report decreasing lattice parameters.

Grew et al. (2008) investigated chemical variability of compounds between  $Al_8Si_2B_2O_{19}$ and  $Al_{18}B_4O_{33}$  and for compositions close to boralsilite ( $Al_{16}B_6Si_2O_{37}$ ). Results of chemical analyses of their own synthesis products and from literature indicated multite-type boron compounds with different composition between  $Al_8B_2Si_2O_{19}$  and  $Al_{18}B_4O_{33}$  and close to the ideal composition of boralsilite.

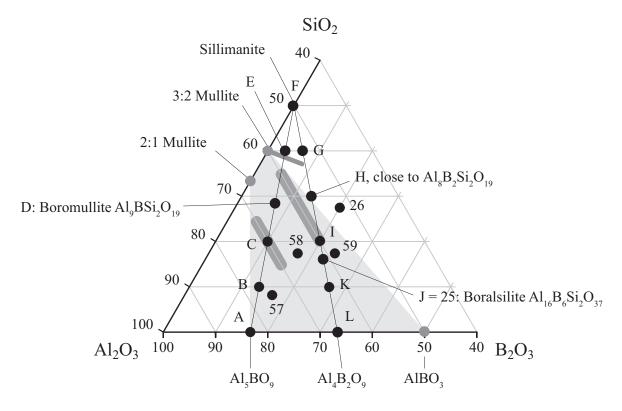


Figure 8.1: Compositions of starting materials A–L in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$  are indicated with black dots. The boron-mullite stability field is plotted as gray area (Werding and Schreyer 1996, compare Section 2.2 in Chapter 2). Incorporation of boron into the mullite crystal structure as investigated by Griesser et al. (2008) is indicated with the two thick gray lines. Compositions of boron-doped mullites according to Zhang et al. (2010) plot on the thin gray line.

Chemically slightly different samples of natural boromullite (Buick et al. 2008) and synthetic compounds of similar composition (Dietzel and Scholze 1955) plot on a line between  $Al_5BO_9$  and sillimanite. Furthermore, the Al/(Si+B) ratio on a line from sillimanite to  $Al_4B_2O_9$  (via  $Al_8B_2Si_2O_{19}$  and boralsilite  $Al_{16}B_6Si_2O_{37}$ ) is constant (Fischer and Schneider 2008a). This may be interpreted as indicator for solid solutions. Unfortunately, except of the study by Griesser et al. (2008), only very restricted areas of the boron-mullite compositional field were systematically investigated.

Hence, 12 samples with starting compositions on the two lines (Figure 8.1) were synthesized by sol-gel routes at 950°C and 1250°C. In addition, five samples with different compositions in or close to the boron-mullite stability field were prepared by a solid-state reaction route.

Boron incorporation into mullite-type crystal structures in the system  $Al_2O_3$ - $B_2O_3$ - $SiO_2$ is most likely due to the following exchange mechanisms: B can either replace Si in  $SiO_4$ tetrahedra or replace Al in  $AlO_4$  tetrahedra. Further, B can be incorporated as a  $BO_3$  unit in place of an  $AlO_4$  or  $SiO_4$  tetrahedron.  $B^{3+} \rightarrow Si^{4+}$  substitutions are not charge neutral, thus, oxygen vacancies will be involved. In boron-mullites  $(Al_2O_3-B_2O_3)$ , B can be present in the mullite-type structure due to  $BO_4 \rightarrow AlO_4$  and/or  $BO_3 \rightarrow AlO_4$  substitutions.

Any  $BO_x \rightarrow (Al,Si)O_4$  substitution in a crystal structure will result in a reduction of unit cell volume due to the smaller space required by  $BO_3$  and  $BO_4$  units compared to  $(Al,Si)O_4$ tetrahedra.  $BO_3$  were distinguished from  $BO_4$  groups with FTIR investigations, whereas lattice parameters were derived from refined powder X-ray diffraction patterns.

Sample	Theoretical stoichiometry	$Al_2O_3 \ (mol-\%)$	$B_2O_3 \pmod{\%}$	$SiO_2 (mol-\%)$
А	$Al_5BO_9$	83.33	16.67	0.00
В		76.55	13.35	10.00
С		70.00	10.00	20.00
D	Boromullite $Al_9Si_2BO_{19}$	64.29	7.14	28.57
Ε		56.57	3.35	40.00
F	Sillimanite $Al_2SiO_5$	50.00	0.00	50.00
G		53.37	5.59	40.00
Η	close to $Al_8Si_2B_2O_{19}$	56.61	13.38	30.00
Ι		60.00	20.00	20.00
J	Boralsilite $Al_{16}B_6Si_2O_{37}$	61.54	23.08	15.38
Κ		63.31	26.69	10.00
$\mathbf{L}$	$Al_4B_2O_9$	66.67	33.33	0.00
25	Boralsilite $Al_{16}Si_2B_6O_{37}$	61.54	23.08	15.38
26	$Al_{16}Si_4B_6O_{41}$	53.33	20.00	26.67
57	$Al_{4.5}Si_{0.5}BO_{9.25}$	75.00	16.67	8.33
58	$Al_4SiBO_{9.5}$	66.67	16.67	16.67
59	$Al_{3.5}Si_{1.5}BO_{9.75}$	58.33	25.00	16.67

Table 8.1: Composition of starting materials of samples A–L (sol-gel route) and #25, #26 and #57-59 (solid-state route).

### 8.2 Experimental

Samples A–L were produced from sol-gel precursors consisting of aluminum nitrate nonahydrate (Al(NO<sub>3</sub>)<sub>3</sub> · 9 H<sub>2</sub>O), tetra ethyl orthosilicate (C<sub>8</sub>H<sub>20</sub>O<sub>4</sub>Si) and boric acid (H<sub>3</sub>BO<sub>3</sub>). Each precursor was heated at 950°C and at 1250°C, respectively. Samples #25, #26 and #57–59 were produced by a solid-state reaction from corundum (Al<sub>2</sub>O<sub>3</sub>), dehydrated silica gel (SiO<sub>2</sub>) and B<sub>2</sub>O<sub>3</sub> at 1200°C (#25 and #26) or 1300°C (#57–59). Sample compositions are shown in Figure 8.1 and listed Table 8.1. Further details on sample preparation are given in Section 3.2.2 in Chapter 3.

After grinding, all samples were measured with a PANalytical X'pert Pro MPD powder diffractometer (Cu $K\alpha$ ) and phases were identified with the ICDD PDF-2 2004 database. Samples A–F were measured with exactly the same conditions (5°–75° 2 $\theta$ , 40 s/step and 0.008°/step). Powder XRD pattern of all samples are shown in Figures 8.5 to 8.9.

Fourier transformed infrared powder spectra of 950°C samples A–L (Figures 8.10 and 8.11) were measured on a Perkin Elmer Spectrum One FTIR spectrometer with a Perkin

Elmer universal ATR accessory (Zn-Se + diamond crystal). The beam path included a MIR light source, a KBr beam splitter and a TGS detector. Sample and background spectra were averaged from 50 scans measured 550 cm<sup>-1</sup> to 1500 cm<sup>-1</sup> at 4 cm<sup>-1</sup> resolution. A sample of pure corundum, and a mixture of 50:50 mol-% sillimanite were also measured on this spectrometer (Figure 8.14).

FTIR spectra of 1250°C samples A–L (Figures 8.12 and 8.13) were acquired from 300 cm<sup>-1</sup> to 2000 cm<sup>-1</sup> on a Bruker Tensor 27 FTIR spectrometer equipped with a globar MIR light source, a KBr beam splitter, a DLaTGS detector and a a Harrick MVP 2 diamond ATR accessory. Sample and background spectra were averaged from 100 scans at 4 cm<sup>-1</sup> resolution. Reference spectra of  $Al_5BO_9$  and  $Al_4B_2O_9$  (Chapter 4) were measured on this spectrometer (Figure 8.14).

Spectra were interpreted based on literature data (band assignation according to Griesser et al. 2008) and by comparison to spectra of reference samples of  $Al_4B_2O_9$ ,  $Al_5BO_9$ , corundum and the 50:50 mol-% mixture of sillimanite and  $Al_5BO_9$ .

#### 8.3 Results

Lattice parameters of 950°C samples A–L were Pawley refined starting with mullite-like lattice parameters a = 7.5 Å, b = 7.6 Å, c = 2.8 Å in space group *Pbam*. Refined lattice parameters are given in Table 8.2 and shown in Figure 8.2. Samples A–H showed a broad peak at ca. 20° 20 which may fit a cubic phase with  $a \approx 7.887$  Å (space group  $Fd\bar{3}m$ , e.g.  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>, Chakravorty 1993, Gutierrez et al. 2002, Duoy 2005 and Cividanes et al. 2010). According to Chakravorty (1993) and Chakraborty and Das (2003), a Si-Al spinel-type phase with a similar X-ray powder pattern exists. This phase may incorporate minor amounts of Si and is referred to as "spinel-type" phase here. Quartz could be identified in sample I and may be present in samples J and K. X-ray diffraction patterns of the Pawley refinements are shown in Figures 8.5 and 8.6.

Samples A–L heated at 1250°C and samples #25, #26 (1300°C) and #57-59 (1200°C) fit to ICDD PDF-2 database entries of  $Al_5BO_9/Al_{18}B_4O_{33}$  and corundum (A–E and G–L). In sample F, mullite, corundum and cristobalite were identified. Cristobalite is normally not formed below 1470°C, however, impurities may lead to metastable cristobalite formation at lower temperatures (Smallman and Bishop 1999).

Choosing a suitable structural model for Rietveld refinements is tricky. Based on peak positions an  $Al_5BO_9$ - or boromullite-like phase in space group  $Cmc2_1$  can also be fitted with a mullite-type phase in *Pbam*. In the latter case, a single peak at ca. 20° 20 will not be fitted. A phase in  $Cmc2_1$  with similar lattice parameters as  $Al_5BO_9$  has two peaks at ca. 16.5° 20 and 16.7° 20 ((021) and (110)). Mullite in *Pbam* has only one peak at ca. 16.4° 20 (110). The occurrence of the two peaks at ca. 16.5° 20 (unambiguously present in patterns of samples A–D and H–L; It may also be assumed in E and G) and the unassigned peak at ca. 20° 20 were taken as evidence that a boromullite-like phase in space group  $Cmc2_1$  (Buick et al. 2008) but not a mullite-like phase in *Pbam* is present. Lattice parameters of this phase are listed in Table 8.3 and plotted in Figure 8.3.

The percentage of  $Al_5BO_9$ - and sillimanite-like modules was refined by introducing a new parameter "wt. BM". For  $Al_5BO_9$ , wt. BM = 1, for a sillimanite-like structure (in space group  $Cmc2_1$  with a doubled *b*-axis), wt. BM = 0, whereas 0.5 presents ideal boromullite ( $Al_5BO_9$ :Al4Si2O10). From simulated powder XRD patterns of  $Al_5BO_9$  and boromullite, refinement of the wt. BM parameter proved to be stable and reproducible. Values of the wt. BM parameter are shown in Figure 8.4 and listed in Table 8.4. Refining wt. BM vastly improved fitting of intensities, however, simultaneous refinement of atomic positions was not possible.

By applying the Rietveld method, excess phases could be quantified (Table 8.4 and Figure 8.3). Corundum quantification of sample E yielded 4.4 wt.-% more excess corundum than  $Al_2O_3$  in the starting mixture (refining wt. BM is not the reason for this result). Thus, the presence of X-ray amorphous phases is assumed. The different intensities of the diffraction patterns (Figures 8.7 and 8.8) can also be explained with varying amount of amorphous phases.

In all  $B_2O_3$ -containing 950°C samples (A–E, G–L),  $BO_3$  groups are present based on IR bands between 1200 cm<sup>-1</sup> and 1400 cm<sup>-1</sup>.  $BO_4$  characteristic absorptions are expected between ca. 1030 cm<sup>-1</sup> and 1150 cm<sup>-1</sup>.  $SiO_4$  vibrations have IR bands in the same region, thus a distinction in samples containing both  $SiO_2$  and  $B_2O_3$  in the starting materials is difficult. However,  $BO_4$  groups are apparently present in samples I–L (containing 20 or more mol-%  $B_2O_3$ ), as the two IR bands at ca. 1030 cm<sup>-1</sup> and 1150 cm<sup>-1</sup> increase with decreasing  $SiO_2$  content.

The region below ca. 900 cm<sup>-1</sup> is mainly ascribed to  $AlO_4$  and  $AlO_6$  vibrations. Resolution of the spectra in this region differ significantly. Whenever the spinel-type phase is clearly identifiable in powder XRD patterns, resolution is poor (Figures 8.7 and 8.8).

BO<sub>3</sub> IR vibrations are also observed for in 1250°C samples A–L (except F, which does not contain B<sub>2</sub>O<sub>3</sub>). In samples with SiO<sub>4</sub> contents below ca. 40 mol-% SiO<sub>2</sub>, an IR band at 1015 cm<sup>-1</sup> fits to the symmetric stretching mode  $v_1$  of a BO<sub>3</sub> group (discussed in Section 4.6 in Chapter 4). The spectrum of sample F (mullite, cristobalite and corundum) differs from the other spectra by a broad IR band at ca. 1040 cm<sup>-1</sup> to 1200 cm<sup>-1</sup>. AlO<sub>x</sub> vibrations are generally found at lower wavenumbers (reference spectrum of corundum in Figure 8.14, and Griesser et al. 2008) and the spectrum of the Al<sub>5</sub>BO<sub>9</sub>-sillimanite mixture shows a similar, but less broad IR band in the same range. Intensity of bands in the same region increases with composition towards sample F and decreases towards sample L. For the 1250°C samples, this region is ascribed to SiO<sub>4</sub> vibrations.

### 8.4 Discussion

#### 950°C samples A–L

According to results of FTIR spectra and powder XRD data, 950°C samples can be divided into four groups:

**Samples A**<sub>950</sub>–**C**<sub>950</sub>. Only a mullite-like phase and a broad bump at ca. 20° 20 can be identified from XRD patterns of samples A–C. This peak may fit quartz but as this peak is also present in the sample without SiO<sub>2</sub>, the peak is assigned to the spinel-type phase (Chakravorty 1993, Chakraborty and Das 2003). Hence, it is not possible to make any assumption on Al<sub>2</sub>O<sub>3</sub>, B<sub>2</sub>O<sub>3</sub> and SiO<sub>2</sub> distribution.

With chemical compositions towards sillimanite (sample F), lattice parameters of samples A–C increase (Table 8.2 and Figure 8.2). Starting materials of samples B and C (towards sillimanite composition) contain more  $SiO_2$  at the expense of  $Al_2O_3$  and  $B_2O_3$ . It is thus likely that unit cell parameters increase due to incorporation of Si into the multie structure.

Despite increasing lattice parameters, the region of  $SiO_4$  vibrations (around ca. 1100 cm<sup>-1</sup>) is empty in FTIR spectra of samples A–C (Figure 8.10), or IR bands with low intensity cannot be identified due to poor resolution of the spectra. Bands between 1200 cm<sup>-1</sup> and 1400 cm<sup>-1</sup> indicate that boron is three-fold coordinated. The strong increase of IR absorptions below 700 cm<sup>-1</sup> may be due to the spinel-type phase, as this absorption is similar for all samples except I–L, where the spinel-type phase cannot be found in diffraction patterns.

Sample	a-axis (Å)	<i>b</i> -axis (Å)	c-axis (Å)	Volume (Å <sup>3</sup> )
А	7.551(8)	7.6655(7)	2.8395(2)	164.36(3)
В	7.563(8)	7.6657(7)	2.8430(2)	164.83(3)
С	7.568(8)	7.6658(7)	2.85095(19)	165.40(2)
D	7.523(3)	7.651(3)	2.8182(7)	162.19(9)
E	7.4961(16)	7.659(3)	2.8203(7)	161.92(8)
F	7.577(17)	7.781(17)	2.888(4)	170.3(6)
G	7.554(2)	7.610(2)	2.8292(5)	163.70(7)
Η	7.5494(8)	7.6590(7)	2.82845(18)	163.54(2)
Ι	7.4985(8)	7.6438(6)	2.80875(17)	160.99(2)
J	7.4883(11)	7.6175(10)	2.7995(3)	159.69(3)
Κ	7.4978(16)	7.6243(14)	2.8002(3)	160.07(5)
$L^{a}$	14.984(17)	5.569(4)	15.41(2)	1285(3)

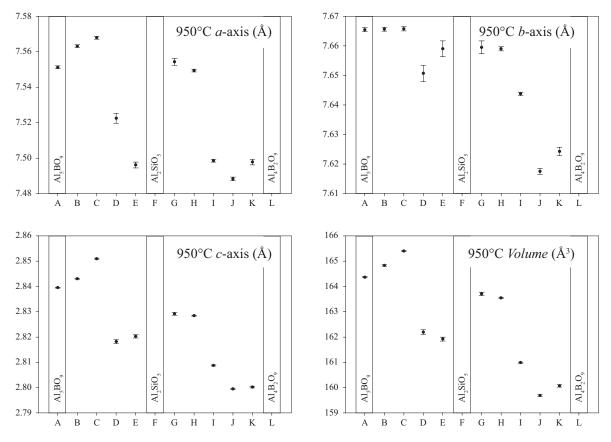
**Table 8.2:** Pawley refined lattice parameters of the mullite-like phase (space group Pbam) of samples A–L heat treated at 950°C.

<sup>a</sup>space group C2/m, monoclinic angle  $\beta = 90.98(7)^{\circ}$  (Fischer et al. 2008b)

Samples  $D_{950}$  and  $E_{950}$ . X-ray diffraction patterns of samples D and E show two strong broad peaks at ca. 46° 2 $\theta$  and 67° 2 $\theta$ . Position and intensities of those peaks are compatible with the spinel-type phase (Chakravorty 1993, Chakraborty and Das 2003). Lattice param-

eters of the mullite-type phases in samples D and E are significantly smaller than those of samples A–C. The significantly lower intensity in the diffraction patterns (Figure 8.5) may be an indicator for at least one X-ray amorphous phase (SiO<sub>2</sub>?). The presence of the excess spinel-type and possibly an amorphous phase may inhibit further Si incorporation into the mullite-type phase. Thus, it is be assumed that the mullite-type phase is in a non-equilibrium state at this composition and temperature.

IR absorptions in the region of 1200 cm<sup>-1</sup> and 1400 cm<sup>-1</sup> (BO<sub>3</sub>) are strongly reduced in sample E, which is in agreement with the fewer amount of  $B_2O_3$  in the starting mixture. In the spectrum of the same sample, a new peak appearing at ca. 1015 cm<sup>-1</sup> may be ascribed to vibrations of SiO<sub>4</sub> (Griesser et al. 2008) or to v<sub>1</sub> of a BO<sub>3</sub> group.



**Figure 8.2:** Pawley-refined lattice parameters of mullite-like phases in space group *Pbam* of samples A–E and G–K. Lattice parameters A–C increase with composition towards F, whereas they decrease with composition from G towards K. Outliers D and E may be due to non-equilibrium conditions (see text). Data from samples F and L are not shown because diffraction patterns are not compatible with a unit cell in space group *Pbam*. For samples A, F and L, starting compositions are indicated.

**Samples G<sub>950</sub>–K<sub>950</sub>.** Lattice parameters of samples G–K decrease with increasing  $B_2O_3$  content of the starting materials (Table 8.2 and Figure 8.2). In the powder XRD pattern of sample G, the peaks at ca. 46° 20 and 67° 20 indicate the presence of the spinel-type phase, whereas in sample H, this phase can only be assumed due to the broad bump at ca. 20° 20.

In sample I, powder XRD peaks match the structure of quartz, whereas in samples J and K, quartz can only be assumed (peak at  $20.8^{\circ} 2\theta$ ).

Besides rather unresolved BO<sub>3</sub> IR bands, a new band appears in FTIR spectra between ca. 1075 cm<sup>-1</sup> and 1150 cm<sup>-1</sup>. This region is compatible with SiO<sub>4</sub> absorptions (Griesser et al. 2008) and has also been observed in FTIR spectra of sample F (see below). From the broad band at ca. 1050 cm<sup>-1</sup> in spectra of samples I, J and K, the presence of BO<sub>4</sub> groups is evident. A second BO<sub>4</sub> band at ca. 1120 cm<sup>-1</sup> interferes with SiO<sub>4</sub> bands (between 1093 cm<sup>-1</sup> and 1131 cm<sup>-1</sup>, Griesser et al. 2008) when BO<sub>4</sub> tetrahedra are present and thus complicates identification. However, from FTIR spectra, it may be assumed that absorptions between ca. 1030 cm<sup>-1</sup> and 1150 cm<sup>-1</sup> are predominantly due to SiO<sub>4</sub> in samples G and H, whereas they are due to BO<sub>4</sub> in samples J and K. BO<sub>4</sub> groups in samples with more than ca. 20 mol-% B<sub>2</sub>O<sub>3</sub> (represented here with I, J, K and L) have also been reported by Griesser et al. (2008).

Samples  $F_{950}$  and  $L_{950}$ . The powder XRD pattern of sample F indicates only poorly crystalline phases. Nonetheless, the spinel-type phase is present. The small broad peak at 26.6° 20, fits the (101) peak of quartz, hence, quartz was included in the refinement. The FTIR spectrum of this sample mainly differs from spectra of the other samples in the region between 1200 cm<sup>-1</sup> to 1400<sup>-1</sup> (as no boron is present) and by the broad, poorly resolved bump between 1000 cm<sup>-1</sup> and ca. 1250 cm<sup>-1</sup>. According to Griesser et al. (2008), no SiO<sub>4</sub> IR bands are present below 980 cm<sup>-1</sup>. Between 980 cm<sup>-1</sup> and 550 cm<sup>-1</sup>, only AlO<sub>4</sub> and AlO<sub>6</sub> absorptions are present. Thus, the broad bump can be ascribed to quartz, whereas the slope below is due to the spinel-type phase.

Diffraction peaks of Sample L fit to a structure in space group C2/m with lattice parameters (Table 8.2) similar to those of  $Al_4B_2O_9$  according to Fischer et al. (2008b). From the FTIR spectrum (Figure 8.11),  $BO_4$  IR bands are clearly visible and the pattern is similar to the FTIR spectrum of the reference  $Al_4B_2O_9$  sample (Chapter 4), shown in Figure 8.14. The differences in the two spectra may be ascribed to the different sample preparation (more  $B_2O_3$  in the starting material of reference  $Al_4B_2O_9$  and to the different FTIR spectrometers used.

Griesser et al. (2008) reported this phase with mullite-like lattice parameters in space group *Pbam*. If the powder XRD pattern of sample L is indexed with a mullite-like phase in *Pbam* instead of C2/m, not all peaks are considered. The sample found by Griesser et al. (2008) may thus have a different structure than sample L synthesized in this study.

#### 1250°C samples A–L

From Rietveld-refined lattice parameters, the wt. BM parameter and excess corundum (Tables 8.3 and 8.4 and Figure 8.3), results of the samples heated at 1250°C can be divided in two groups (A–E and G–L) with a different trend each. In the plots, data for sample F are not included because the powder XRD pattern is not compatible with a phase in space group

 $Cmc2_1$ . Instead, a multie-like phase (*Pbam*), corundum and cristobalite were found. This sample was subsequently heated at 1275°C for 100h but the ratios of multie, cristobalite and corundum did not significantly change. The almost inert behavior of this sample is ascribed to preferential formation of corundum and cristobalite at the cost of multie. Once cristobalite and corundum have crystallized, they react rather sluggish at 1275°C, thus indicating nonequilibrium conditions.

**Trend from Al\_5BO\_9 (A<sub>1250</sub>) to sillimanite (F<sub>1250</sub>) composition.** In Rietveld refinements of samples A–E, the wt. BM parameter increases with increasing composition of the starting materials towards the sillimanite composition (sample F). At the same time, no significant trend in the *a* and *b* lattice parameters and the cell volume can be identified. However, the *c*-axis decreases with composition towards F, which is in agreement with Buick et al. (2008). Towards E, samples consist of significant amounts of excess corundum. Rietveld quantifications of sample E indicate even more excess corundum than  $Al_2O_3$  in the starting mixture. Not all of the material will thus be crystalline or incorporated into the *Cmc*2<sub>1</sub>-phase and actual composition of this phase is certainly different from the starting mixture.

Sample	a-axis (Å)	b-axis (Å)	c-axis (Å)	Volume (Å <sup>3</sup> )
А	5.66830(11)	15.0228(4)	7.69732(17)	655.45(2)
В	5.66977(16)	15.0310(6)	7.6948(3)	655.77(3)
$\mathbf{C}$	5.66943(18)	15.0316(6)	7.6929(3)	655.59(4)
D	5.6638(3)	15.0281(8)	7.6872(4)	654.30(6)
Ε	5.6663(3)	15.03356(7)	7.6879(6)	654.89(8)
$\mathbf{F}^{\mathbf{a}}$	7.5652(3)	7.6885(3)	2.88304(9)	167.704(14)
G	5.6796(3)	15.05822(5)	7.6827(5)	657.05(7)
Η	5.66500(16)	15.0240(5)	7.6892(2)	654.42(4)
Ι	5.66323(14)	15.0162(5)	7.6909(2)	654.03(3)
J	5.66159(16)	15.0161(6)	7.6899(3)	653.76(4)
Κ	5.66008(18)	15.0168(6)	7.6868(3)	653.35(4)
$\mathbf{L}$	5.66843(9)	15.0107(3)	7.69057(14)	654.37(2)
25	5.6610(2)	15.0218(07)	7.6829(3)	653.90(5)
26	5.6675(2)	15.0263(07)	7.6788(4)	653.93(5)
57	5.6873(3)	15.0491(10)	7.6823(5)	657.51(7)
58	5.6880(4)	15.0471(13)	7.6851(7)	657.75(10)
59	5.6793(3)	15.0424(09)	7.6804(5)	656.13(7)

**Table 8.3:** Rietveld refined lattice parameters of samples A–L heated at  $1250^{\circ}$ C and samples #25 and #26 (1300°C) and #57–59 (1200°C) in space group  $Cmc2_1$ .

<sup>a</sup>space group *Pbam* 

FTIR spectra of samples A–E show the typical absorptions ascribed to BO<sub>3</sub> units (1200 cm<sup>-1</sup> and 1400 cm<sup>-1</sup>, Figure 8.12). The low intensity band at 1015 cm<sup>-1</sup> (most expressed in the FTIR spectrum of sample A) fits the  $v_1$  vibration of the BO<sub>3</sub> group. In similarly prepared samples of Griesser et al. (2008), no BO<sub>4</sub> was found. Hence, it is not likely that the the low-intensity bands appearing between 1200 cm<sup>-1</sup> and 1020 cm<sup>-1</sup> result from BO<sub>4</sub> groups. As

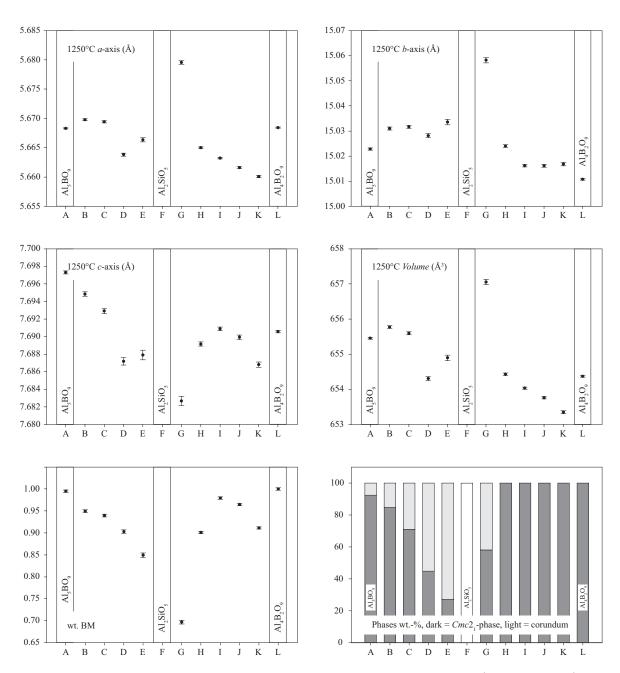


Figure 8.3: Lattice parameters from Rietveld refinements of the  $Cmc2_1$ -phase (middle and top) and the wt. BM parameter (bottom left) of samples A–L. Data from sample F are not shown because lattice parameters are not compatible with a unit cell of a phase in space group  $Cmc2_1$ . The wt. BM factor is the weight of  $Al_5BO_9$  and sillimanite modules in a mullite-type phase in  $Cmc21_1$ . A weight of 1 corresponds to  $Al_5BO_9$ , whereas a weight of 0.5 equals boromullite. Quantification (Rietveld wt.-%) of the  $Cmc2_1$ -phase and corundum is shown in the bottom right diagram. Starting material compositions of A, F and L are indicated.

mentioned before,  $SiO_4$  bands may also occur in this region. By comparing the spectra to the one of sample F (mullite, cristobalite, corundum), and to the spectrum of the 50:50 mol-% mixture of  $Al_5BO_9$  and sillimanite (Figure 8.14), it becomes obvious that the low-intensity

bands in this region in spectra of samples B-E are due to  $SiO_4$  groups, although intensity does not increase with increasing  $SiO_2$ . Below 1000 cm<sup>-1</sup>, IR band positions are in agreement with the reference spectra of pure  $Al_5BO_9$  and the sillimanite- $Al_5BO_9$  mixture mixture (Figure 8.14). Intensities differ significantly with increasing composition towards sample F (most expressed around ca. 550 cm<sup>-1</sup>, where a high-intensity band exists in the references spectrum of corundum (Figure 8.14).

Sample	wt% $Cmc2_1$ phase	wt% $\mathrm{Al}_2\mathrm{O}_3$	wt. BM <sup>a</sup>
А	92.4(3)	7.6(3)	0.995(3)
В	84.7(4)	15.3(4)	0.950(3)
С	71.0(4)	29.0(4)	0.939(3)
D	44.8(3)	55.2(3)	0.903(4)
E	27.0(3)	73.0(3)	0.849(5)
$\mathbf{F}$	$n/a^b$	$n/a^b$	$n/a^b$
G	58.0(4)	42.0(4)	0.697(4)
Η	100.0	0.0	0.901(3)
Ι	100.0	0.0	0.979(3)
J	100.0	0.0	0.965(3)
К	100.0	0.0	0.911(3)
$\mathbf{L}$	100.0	0.0	1.000(3)
25	100.0	0.0	0.881(4)
26	98.6(4)	1.4(4)	0.780(3)
57	95.7(6)	4.3(6)	0.667(3)
58	96(2)	4(2)	0.616(4)
59	91.9(6)	$n/a^c$	0.599(4)

Table 8.4: Additional data from Rietveld refinements for samples A–L heated at 1250°C and samples #25, #26 and #57-59.

<sup>a</sup>wt. BM = structural weight between boromullite in  $Cmc2_1$  (wt. BM = 0.5) and  $Al_5BO_9$  (wt. BM = 1)  $^{b}68.1(4)$  wt.-% Mullite in Pbam + 26.5(4) wt.-% corundum + 5.4(4) wt.-% cristobalite  $^{\circ}2.63(6)$  wt.-% corundum and 5.48(6) wt.-% cristobalite

Trend from sillimanite (F<sub>1250</sub>) to Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub> (L<sub>1250</sub>) composition. Considering the distribution of wt. BM, lattice parameters and missing excess corundum in samples H–L, a gap between sample G and samples H–L is obvious. FTIR spectra of sample G indicate SiO<sub>4</sub> vibrations (from ca. 1050 cm<sup>-1</sup> to 1200 cm<sup>-1</sup>), however, no distinct Si-containing phase could be found in XRD data. Probably, this  $Cmc_{21}$ -phase contains significantly more Si than all the others, which would make sense due to the larger lattice parameters.

In samples H–K, the amount of  $Al_5BO_q$ -modules in the  $Cmc2_1$  phase increases, combined with a decrease of the a-axis and the cell volume. These results are compatible with trends reported by Buick et al. (2008), but no clear trend is evident from evolution of the b- and caxis. In XRD patterns (Figure 8.8), no other phase than the  $Cmc2_1$ -phase could be identified.

IR bands of  $BO_3$  groups are present in FTIR spectra of samples H–L (including the  $v_1$ band at 1015 cm<sup>-1</sup>, Figure 8.13). SiO<sub>4</sub> vibrational bands (between 1200 cm<sup>-1</sup> and 1020 cm<sup>-1</sup>, assigned based on FTIR spectra of sample F and the 50:50 mol-% mixture of sillimanite and  $Al_5BO_9$ ) are less pronounced with lower  $SiO_2$  content of the starting mixtures. In samples I-L, a new band appears at 1200 cm<sup>1</sup>, which could not be identified.

Previous investigations on synthesis of boron-mullites  $(Al_2O_3-B_2O_3)$  have shown that compounds with  $Al_5BO_9$  compositions crystallize at temperatures above ca. 1100°C from starting materials representing  $Al_4B_2O_9$ . The higher amount of  $B_2O_3$  in the starting materials can lead to larger "crystallites" in the resulting  $Al_5BO_9$  compound. Due to a very similar FTIR spectrum of sample L and the reference spectrum of  $Al_5BO_9$  (Figures 8.13 and 8.14), a wt. BM parameter of 1.000(3) (Table 8.4) and similar cell volume to the one of  $Al_5BO_9$ , sample L is interpreted as compound with a composition of or close to  $Al_5BO_9$ .

#### 1200°C–1300°C samples #25, #26 and #57–59.

Lattice parameters and excess corundum of samples #25, #26 and #57–59 (Table 8.3 and Figure 8.4) show no trend with the compositions of the starting material (increasing sample numbers do not represent continuous change in composition of starting materials, Figure 8.1 and Table 8.1). Nonetheless, unit cell volumes of samples #25 and #26 are close to the one of  $Al_5BO_9$ . As no  $SiO_2$  phase can be identified in the powder patterns of those samples, Si must either be incorporated in the structure or be present as X-ray amorphous phase. In the powder XRD pattern of sample #26 (Figure 8.9), a broad bump between ca. 18° 20 and 25° 20 may indicate an almost X-ray amorphous phase (SiO<sub>2</sub>?). The same can be observed in the difference plots of sample #57 and #58. A single weak peak at ca. 21.7° 20 might fit cristobalite.

Due to the inconsistent behavior of lattice parameters and the wt. BM parameter, it can be assumed that the synthesis route via solid-state reaction is not suitable for the incorporation of Si into a mullite-type structure in space group  $Cmc2_1$ , most likely due to the sluggish reaction behavior of Al<sub>2</sub>O<sub>3</sub> and the lack of a fluxing agent (e.g. B<sub>2</sub>O<sub>3</sub> in sufficient quantities).

#### Comparison with literature data

Trends of the 950°C samples are in accordance with findings of Griesser et al. (2008), although not the same area of the boron-mullite stability field (Werding and Schreyer 1996) was covered with the composition of the starting materials. Decreasing lattice parameters with increasing  $B_2O_3$  content and incorporation of four-fold coordinated boron into the structure for samples produced with more than ca. 20 mol-%  $B_2O_3$  in the starting mixture can be confirmed. However, excess SiO<sub>2</sub> has not been found in samples with  $BO_4$  in the mullite-like structure. The spinel-type phase has not been reported by Griesser et al. (2008) and it is unknown if X-ray amorphous phases were considered.

For samples A and L (both consisting solely of  $Al_2O_3$  and  $B_2O_3$ ), results of Mazza et al. (1992) and MacKenzie et al. (2007), reporting decreasing lattice parameters and  $BO_4$  incorporation into the structure, can be confirmed (for better comparison with literature data, indexing the powder XRD pattern of sample L with a mullite-like structure in space group Pbam leads to a cell volume of 157.59(7) Å<sup>3</sup>).

Combining our data with those of Griesser et al. (2008) and MacKenzie et al. (2007), may indicate that all samples synthesized at 950°C with a maximum of 20 mol-% SiO<sub>2</sub> and more than 20 mol-%  $B_2O_3$  in the starting mixtures contain four-fold coordinated boron. As the chemical difference between sample composition and the composition of the starting materials is not precisely known (this study, Griesser et al. 2008 and MacKenzie et al. 2007), suggestion of a stability field for BO<sub>4</sub>-containing mullite-type structures may not be accurate and is only mentioned as suggestion for further studies.

No studies exist in which  $Al_2O_3$ - $B_2O_3$ - $SiO_2$ -phases synthesized at 1250°C were investigated for variable structural incorporation of  $Al_5BO_9$  and sillimanite-like modules. Furthermore, our data of the 1250°C samples are significantly different to those of 1300°C samples by Griesser et al. (2008), because in this study, mullite (with mullite-like lattice parameters in space group *Pbam*) was only found in sample F containing no  $B_2O_3$  in the precursor. Unfortunately, Griesser et al. (2008) do not mention how they identified the phases reported as "mullite" and if they found unidentified peaks in their powder XRD data. Further, Griesser et al. (2008) did also not indicate which  $SiO_2$  phase they identified in the powder XRD patterns. This is a crucial point. If a diffraction pattern of a  $Cmc2_1$ -phase like boromullite or  $Al_5BO_9$  is fitted with a mullite-like unit cell in space group *Pbam*, the unassigned peaks can be misleadingly interpreted as quartz, although with slightly too large lattice parameters.

As previously mentioned for the 950°C samples, based on FTIR and powder X-ray diffraction data, it is unknown how much the composition of the synthesized mullite-type phases is represented by the composition of the starting materials. High amounts of excess corundum clearly indicate that this difference is rather large. Once corundum has crystallized from the initially amorphous precursor material, it is almost inactive and hinders further crystallization of mullite-type compounds.

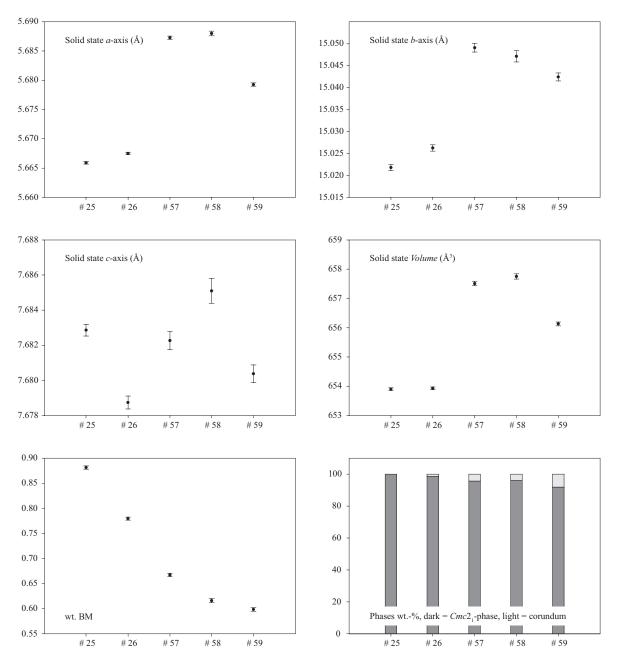
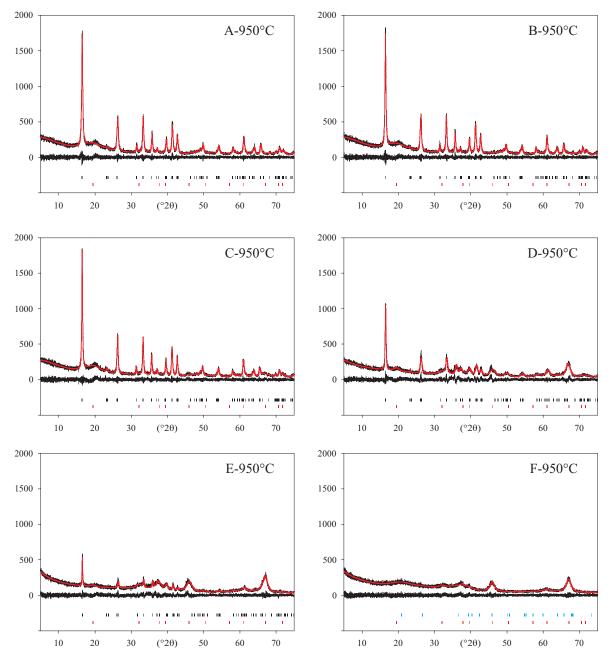
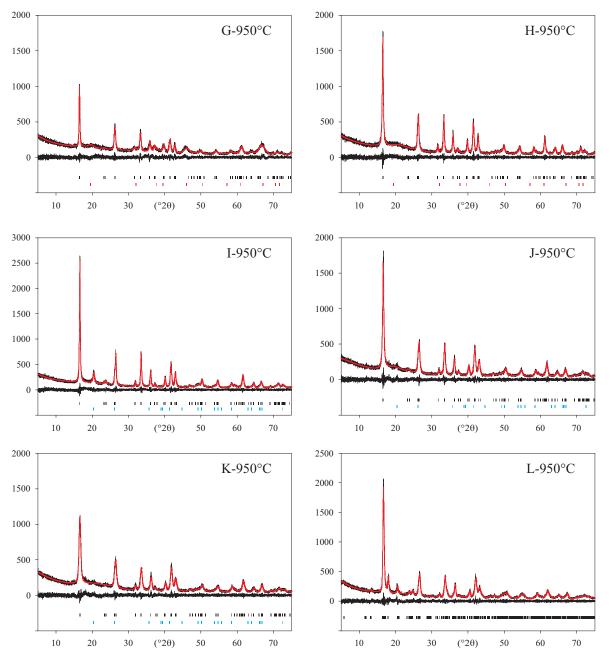


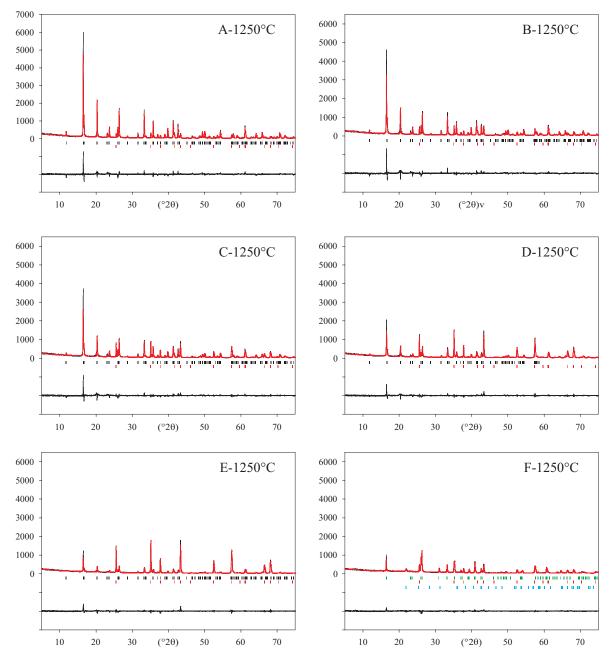
Figure 8.4: Rietveld-refinement results of samples #25, #26 and #57-59. Sample nomenclature does not represent continuous change in composition of the starting mixtures, thus, the decrease of the wt. BM parameter is not in agreement with continuous change of starting compositions. Lattice parameters of the  $Cmc2_1$ -phase are shown top and middle left and right. The wt. BM factor (bottom left) is the fraction of  $Al_5BO_9$  and sillimanite modules in this mullite-type phase. A weight of 1 corresponds to  $Al_5BO_9$ , whereas a weight of 0.5 equals boromullite. Quantification of the  $Cmc2_1$ -phase and corundum is shown in the bottom right diagram. For sample #59, the quantification the light-gray area represents summed wt.-% of corundum and cristobalite.



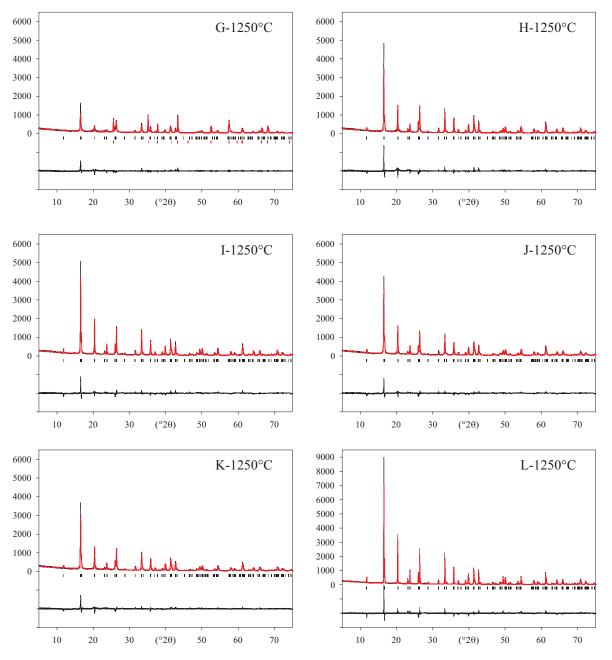
**Figure 8.5:** Pawley-refined XRD patterns of 950°C samples A–F. hkl-ticks of phases with mullite-like lattice parameters in space group *Pbam* are black, hkl-ticks of the cubic spinel-type phase are red, hkl-ticks of quartz are blue. Intensity scale is counts.



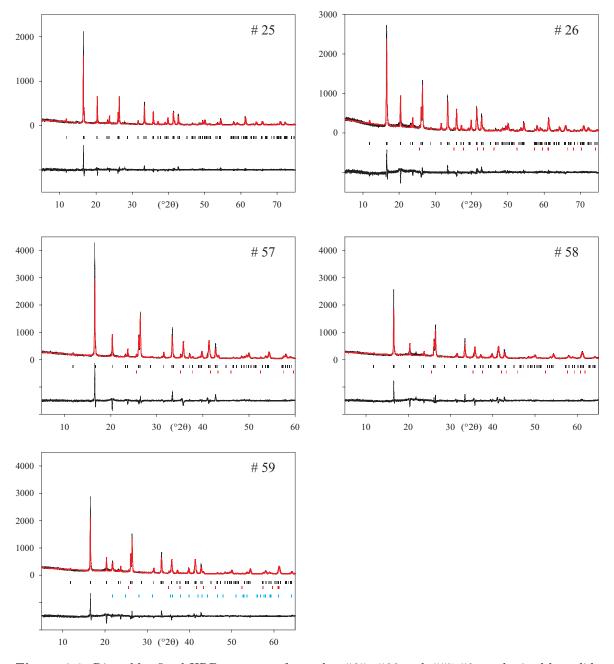
**Figure 8.6:** Pawley-refined XRD patterns of  $950^{\circ}$ C samples G–L. *hkl*-ticks of phases with mullite-like lattice parameters in space group *Pbam* are black, *hkl*-ticks of the cubic spinel-type phase are red, *hkl*-ticks of quartz are blue. Intensity scale is counts.



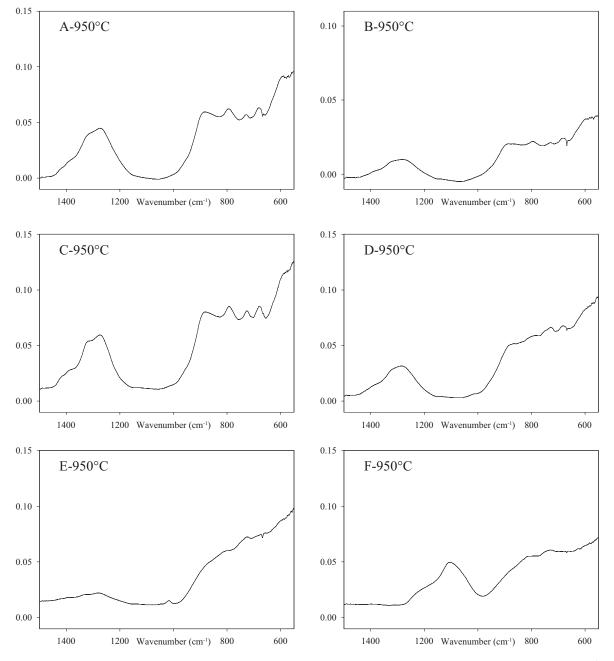
**Figure 8.7:** Rietveld-refined XRD patterns of  $1250^{\circ}$ C samples A–F. *hkl*-ticks of phases refined in space group  $Cmc2_1$  with variable ratio of  $Al_5BO_9$  and sillimanite modules are black. Corundum *hkl*-ticks are red, *hkl*-ticks of cristobalite are blue and *hkl*-ticks of the phase identified as mullite (sample F) in *Pbam* are green. Intensity scale is counts.



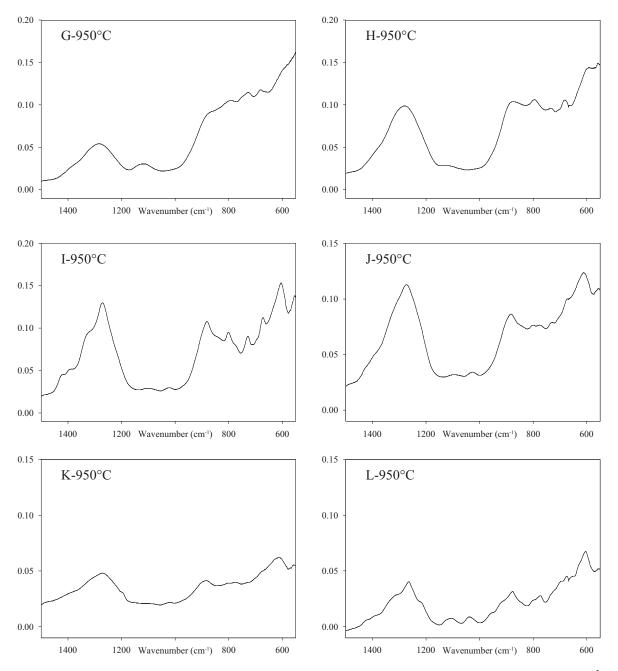
**Figure 8.8:** Rietveld-refined XRD patterns of 1250°C samples G–L. *hkl*-ticks of phases refined in space group  $Cmc2_1$  with variable ratio of  $Al_5BO_9$  and sillimanite modules are black. Corundum *hkl*-ticks are red. Intensity scale is counts.



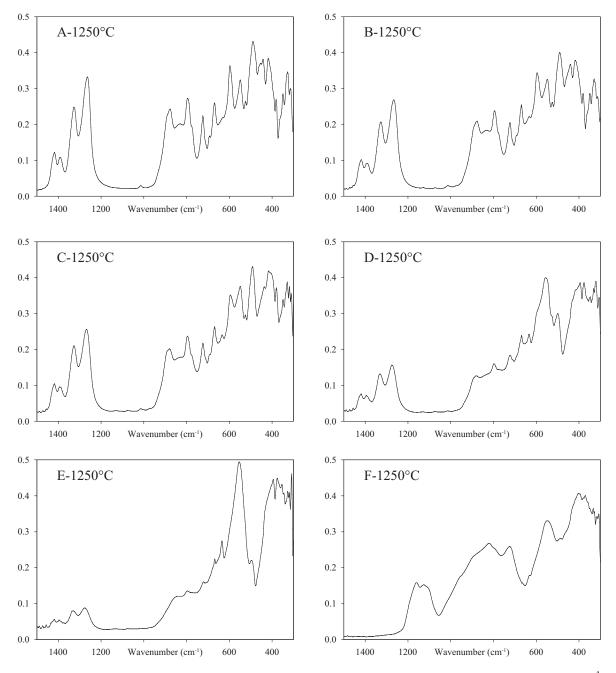
**Figure 8.9:** Rietveld-refined XRD patterns of samples #25, #26 and #57-59 synthesized by solidstate reactions at 1200°C or 1300°C. *hkl*-ticks of phases refined in space group  $Cmc2_1$  with variable ratio of  $Al_5BO_9$  and sillimanite modules are black. Corundum *hkl*-ticks are red, cristobalite *hkl*-ticks are blue. Diffraction patterns were not collected with the same conditions, thus, differing intensities are not meaningful. Intensity scale is counts.



**Figure 8.10:** ATR-FTIR spectra of samples A–F heated at 950°C. Absorptions between 1400 cm<sup>-1</sup> and 1200 cm<sup>-1</sup> are ascribed to BO<sub>3</sub> vibrations, those between ca. 1200 cm<sup>-1</sup> and 1000 cm<sup>-1</sup> to SiO<sub>4</sub> and IR bands below ca. 900 <sup>-1</sup> to AlO<sub>x</sub> groups. Intensity scale is ATR absorbance.



**Figure 8.11:** ATR-FTIR spectra of samples G–L heated at 950°C. Absorptions between 1400 cm<sup>-1</sup> and 1200 cm<sup>-1</sup> are ascribed to BO<sub>3</sub> vibrations and those below ca. 900<sup>-1</sup> to AlO<sub>x</sub> groups. IR bands between 1200 cm<sup>-1</sup> and 1000 cm<sup>-1</sup> (assigned to SiO<sub>4</sub> vibrations in Figure 8.10) decrease with increasing  $B_2O_3$  content of the starting materials. The two resolved bands observed at ca. 1050 cm<sup>-1</sup> and 1130 cm<sup>-1</sup> dominant in samples I–L are ascribed to BO<sub>4</sub> units. Intensity scale is ATR absorbance.



**Figure 8.12:** ATR-FTIR spectra of samples A–F heated at 1250°C. IR bands between 1400 cm<sup>-1</sup> and 1200 cm<sup>-1</sup> are ascribed to BO<sub>3</sub> vibrations and those between ca. 1200 cm<sup>-1</sup> and 1000 cm<sup>-1</sup> to SiO<sub>4</sub> (most expressed in sample F consisting of mullite, cristobalite and corundum). Bands below ca. 900 cm<sup>-1</sup> are ascribed to AlO<sub>x</sub> groups. The predominant broad peak at 600 cm<sup>-1</sup> (samples D and E) is likely correlated with the high amount of excess corundum. Intensity scale is ATR absorbance.

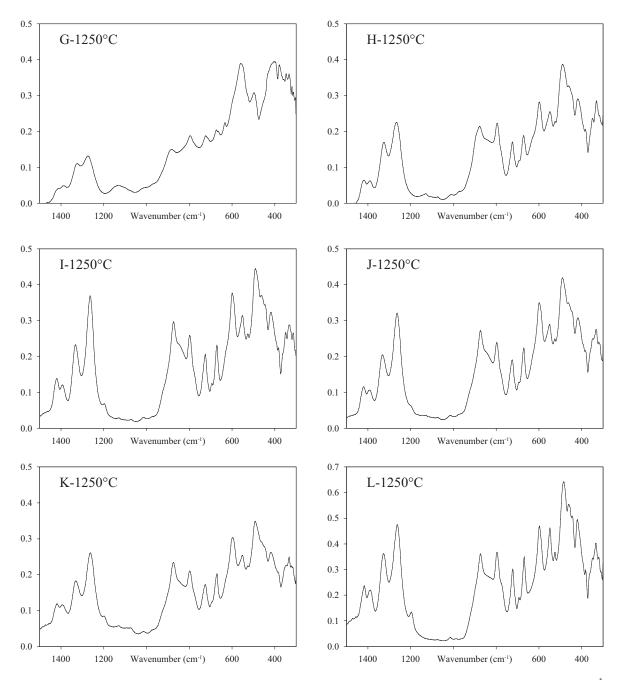


Figure 8.13: ATR-FTIR spectra of samples G–L heated at  $1250^{\circ}$ C. IR bands between 1400 cm<sup>-1</sup> and 1200 cm<sup>-1</sup> are ascribed to BO<sub>3</sub> vibrations. From a sample consisting of mullite, cristobalite and corundum, bands between ca. 1200 cm<sup>-1</sup> and 1000 cm<sup>-1</sup> are ascribed to SiO<sub>4</sub> vibrations. They are less pronounced with decreasing SiO<sub>2</sub> content of starting materials. Bands below ca. 900 <sup>-1</sup> are ascribed to AlO<sub>x</sub> groups. The predominant broad peak at 600 cm<sup>-1</sup> in sample G is likely correlated with the high amount of excess corundum. A peak appearing at ca. 1200 cm<sup>-1</sup> could not be identified. Intensity scale is ATR absorbance.

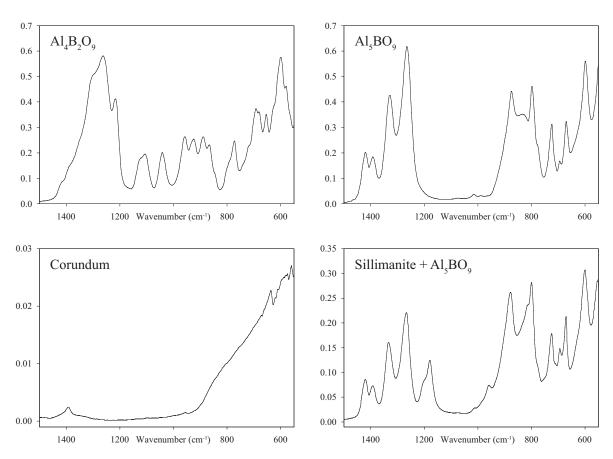


Figure 8.14: ATR-FTIR spectra of reference samples  $Al_5BO_9$ ,  $Al_4B_2O_9$ , corundum and a mixture of 50:50 mol-% sillimanite and  $Al_5BO_9$ . From the spectrum of  $Al_4BO_9$ , compared to  $Al_5BO_9$ , the two additional IR bands at ca. 1050 cm<sup>-1</sup> and 1150 cm<sup>-1</sup> can be assigned to  $BO_4$  vibrations. The spectrum of the mixture of 50:50 mol-% sillimanite and  $Al_5BO_9$  differs mainly from the spectrum of pure  $Al_5BO_9$  by the broad band below 1200 cm<sup>-1</sup>, which can be ascribed to  $SiO_4$  vibrations. Corundum absorptions are most expressed below ca. 900 cm<sup>-1</sup> and will interfere in samples with excess corundum. Intensity scale is ATR absorbance.

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#### A.1 Supplementary Data for Chapter 4

### Single-crystal XRD data of sample *a* refined with fully occupied positions according to refinement strategy 1

**Table A1.1:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample a without resolution restriction and fully occupied positions (refinement strategy 1).

	Diffractometer: Enraf Nonius CAD4						
X-ray radiation	Mo <i>K</i> α (0.71073 Å)	Reflections rejected	92				
X-ray power	50 kV, 40 mA	Reflections used	5636				
Crystal size	$0.2 \cdot 0.13 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.93^{\circ}$				
Measurement time	Max. $120 \text{ sec/step}$	Index range $h$	-9 9				
Absorption corr.	$\psi$ scans	Index range $k$	-24 24				
Resolution	No restriction	Index range $l$	-12 12				
Sample name	a  (ABO-01-2009,  # 01)	Unique reflections	1552				
Temperature	298K	$ ext{Reflections} > 4 \mathbf{\sigma}(I)$	1423				
Space group	$Cmc2_1$	R(int)	0.0373				
a-axis	5.6682(2) Å	$R(\sigma)$	0.0265				
b-axis	15.0060(9) Å	Refined parameters	82				
c-axis	7.6892(4) Å	GoF	1.183				
Volume	654.07(6) Å <sup>3</sup>	$R1,I>4{ m s}(I)$	0.0200				
Z	4	R1, all data	0.0263				
Density $\rho$	$2.942  { m g/cm^3}$	w $R2$ on $F^2$	0.0510				
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.54~{\rm e}^{\rm -}{\rm \cdot}{\rm \AA}^{\rm -3}$ near Al2				
Reflections collected	5728	$\Delta  ho_{\min}$	-0.29 e <sup>-</sup> ·Å <sup>-3</sup> near Al1				

**Table A1.2:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample a without resolution restriction and fully occupied positions (refinement strategy 1).

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24814(7)	0.38355(3)	0.0	1.0	8b	0.00412(7)
Al2	0.0	0.25529(4)	0.68662(10)	1.0	4a	0.00550(10)
Al3	0.0	0.44414(4)	0.68202(11)	1.0	4a	0.00484(10)
Al4	0.0	0.29692(4)	0.33410(11)	1.0	4a	0.00449(10)
B1	0.5	0.51575(13)	0.2216(3)	1.0	4a	0.0051(3)
01	0.0	0.32926(9)	0.87072(19)	1.0	4a	0.0068(2)
O2	0.24304(16)	0.30970(7)	0.20133(15)	1.0	8b	0.00671(17)
O3	0.5	0.45357(9)	0.09233(19)	1.0	4a	0.0059(2)
O4	0.0	0.54594(9)	0.57205(19)	1.0	4a	0.0061(2)
O5	0.0	0.35189(9)	0.53795(18)	1.0	4a	0.0074(2)
O6	0.0	0.19136(9)	0.4349(2)	1.0	4a	0.0050(2)
07	0.28876(17)	0.45168(7)	0.78556(14)	1.0	8b	0.00681(16)

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00319(13)	0.00474(14)	0.00442(13)	0.00044(14)	-0.00007(11)	0.00022(11)
Al2	0.0050(2)	0.0048(2)	0.0067(2)	-0.0001(2)	0.0	0.0
Al3	0.0039(2)	0.0042(2)	0.0064(2)	-0.0005(2)	0.0	0.0
Al4	0.0043(2)	0.0052(2)	0.0039(2)	0.00042(18)	0.0	0.0
B1	0.0047(7)	0.0058(8)	0.0050(8)	0.0003(6)	0.0	0.0
O1	0.0044(5)	0.0083(5)	0.0075(5)	-0.0037(5)	0.0	0.0
O2	0.0056(4)	0.0078(4)	0.0068(4)	0.0022(3)	0.0013(3)	0.0011(3)
O3	0.0050(5)	0.0057(5)	0.0068(5)	-0.0023(4)	0.0	0.0
O4	0.0038(5)	0.0053(5)	0.0091(6)	0.0028(5)	0.0	0.0
O5	0.0110(6)	0.0054(5)	0.0058(5)	-0.0008(4)	0.0	0.0
O6	0.0045(5)	0.0037(5)	0.0069(5)	0.0006(4)	0.0	0.0
07	0.0049(3)	0.0085(4)	0.0071(4)	0.0035(3)	-0.0007(3)	-0.0008(3)

**Table A1.3:** Anisotropic displacement parameters from single-crystal XRD data of sample a without resolution restriction and fully occupied positions (refinement strategy 1).

**Table A1.4:** Individual and mean polyhedral bond lengths of the refined structure from sample a without resolution restriction and fully occupied positions (refinement strategy 1).

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O4	1.8453(9)	Al2	O2(x2)	1.7566(9)
Al1	O6	1.8849(9)	Al2	01	1.7989(14)
Al1	O2	1.9041(11)	Al2	O5	1.8460(15)
Al1	O1	1.9054(10)	Al2	O6	2.1603(16)
Al1	O3	1.9095(10)		<al2-o></al2-o>	1.8637(13)
Al1	07	1.9538(10)			
	<al1-o></al1-o>	1.9005(10)	Al3	O4	1.7460(14)
			Al3	O5	1.7730(14)
Al4	O2 (x2)	1.7254(10)	Al3	O7 $(x2)$	1.8237(10)
Al4	O6	1.7633(14)	A13	01	2.2531(16)
Al4	O5	1.7713(15)		<al3-o></al3-o>	1.8839(13)
	<al4-o></al4-o>	1.7466(12)			
B1	O3	1.363(2)			
B1	O7 $(x2)$	1.3837(14)			
	<b1-o></b1-o>	1.3699(18)			

# Single-crystal XRD data of sample *f* refined with fully occupied positions according to refinement strategy 1

**Table A1.5:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample f without resolution restriction and fully occupied positions (refinement strategy 1).

Diffractometer: Bruker Smart APEX 2							
X-ray radiation	$MoK\alpha (0.71073 \text{ Å})$	Reflections rejected	170				
X-ray power	50  kV, 40  mA	Reflections used	9250				
Crystal size	$0.15 \cdot 0.025 \cdot 0.05 \text{ mm}^3$	Max. 20	$73.38^{\circ}$				
Measurement time	60  s/frame	Index range $h$	-9 9				
Absorption corr.	Empirical $\psi$ scans	Index range $k$	$-24 \dots 25$				
Resolution	No restriction	Index range $l$	-12 11				
Sample name	f (ABO-62, # 62-2)	Unique reflections	1618				
Temperature	298K	$ ext{Reflections} > 4\sigma(I)$	1565				
Space group	$Cmc2_1$	$R( ext{int})$	0.0258				
a-axis*	5.6618(7) Å	$R(\mathbf{\sigma})$	0.0194				
b-axis*	14.9981(12) Å	Refined parameters	82				
c-axis*	7.6806(7) Å	GoF	1.152				
$Volume^*$	652.21(12) Å <sup>3</sup>	$R1,I>4{f \sigma}(I)$	0.0228				
Ζ	4	R1, all data	0.0239				
Density $\rho$	$2.250~{ m g/cm^3}$	w $R2$ on $F^2$	0.0605				
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.50~{\rm e}^{\rm -}{\rm \cdot}{\rm \AA}^{\rm -3}$ near Al3				
Reflections collected	9420	$\Delta  ho_{ m min}$	-0.43 e <sup>-</sup> ·Å <sup>-3</sup> near Al2				

\*lattice parameters from CAD4 (15 centered reflections with  $\pm 11^\circ < \theta < \pm 30^\circ)$ 

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24812(6)	0.38351(3)	0.0	1.0	8b	0.00532(8)
Al2	0.0	0.25533(4)	0.68638(11)	1.0	4a	0.00695(10)
Al3	0.0	0.44415(4)	0.68164(11)	1.0	4a	0.00614(10)
Al4	0.0	0.29686(4)	0.33409(12)	1.0	4a	0.00626(10)
B1	0.5	0.51544(13)	0.2215(3)	1.0	4a	0.0064(3)
01	0.0	0.32905(9)	0.8706(2)	1.0	4a	0.0073(2)
O2	0.24297(16)	0.30960(7)	0.20111(17)	1.0	8b	0.00836(18)
O3	0.5	0.45349(9)	0.0927(2)	1.0	4a	0.0067(2)
O4	0.0	0.54592(9)	0.5721(2)	1.0	4a	0.0074(2)
O5	0.0	0.35200(9)	0.5380(2)	1.0	4a	0.0087(2)
O6	0.0	0.19128(9)	0.4350(2)	1.0	4a	0.0066(2)
07	0.28883(18)	0.45169(7)	0.78536(15)	1.0	8b	0.00778(17)

**Table A1.6:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample f without resolution restriction and fully occupied positions (refinement strategy 1).

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00412(14)	0.00617(14)	0.00569(15)	0.00001(15)	-0.00016(11)	0.00032(11)
Al2	0.0061(2)	0.0064(2)	0.0084(2)	0.0(2)	0.0	0.0
Al3	0.0049(2)	0.0058(2)	0.0077(2)	-0.0001(2)	0.0	0.0
Al4	0.0060(2)	0.0071(2)	0.0057(2)	0.00013(19)	0.0	0.0
B1	0.0061(7)	0.0073(8)	0.0057(8)	-0.0004(6)	0.0	0.0
01	0.0047(5)	0.0089(5)	0.0083(6)	-0.0036(5)	0.0	0.0
O2	0.0075(4)	0.0090(4)	0.0086(4)	0.0020(3)	0.0014(3)	0.0006(3)
O3	0.0058(5)	0.0062(5)	0.0080(6)	-0.0023(4)	0.0	0.0
O4	0.0052(5)	0.0060(5)	0.0109(6)	0.0027(5)	0.0	0.0
O5	0.0113(6)	0.0072(5)	0.0076(6)	-0.0009(4)	0.0	0.0
O6	0.0051(5)	0.0058(5)	0.0088(6)	0.0012(4)	0.0	0.0
07	0.0048(3)	0.0102(4)	0.0083(4)	0.0038(3)	-0.0002(3)	-0.0009(3)

**Table A1.7:** Anisotropic displacement parameters from single-crystal XRD data of sample f without resolution restriction and fully occupied positions (refinement strategy 1).

**Table A1.8:** Individual and mean polyhedral bond lengths of the refined structure from sample f without resolution restriction and fully occupied positions (refinement strategy 1).

Center	Ligand	Distance (Å)	Center	Ligand	Distance $(Å)$
Al1	O4	1.8441(9)	Al2	O2 (x2)	1.7547(10)
Al1	O6	1.8818(9)	Al2	01	1.7958(14)
Al1	O2	1.9015(12)	Al2	O5	1.8443(15)
Al1	01	1.9047(10)	Al2	O6	2.1566(16)
Al1	O3	1.9087(10)		<al2-o></al2-o>	1.8776(14)
Al1	07	1.9536(11)			
	<al1-o></al1-o>	1.8991(9)	Al3	O4	1.7428(14)
			Al3	O5	1.7685(15)
Al4	O2 (x2)	1.7240(11)	Al3	O7 (x2)	1.8225(10)
Al4	O6	1.7629(14)	Al3	01	2.2555(16)
Al4	O5	1.7709(16)		<al3-o></al3-o>	1.8824(13)
	<al4-o></al4-o>	1.7526(14)			
B1	O3	1.357(2)			
B1	O7 $(x2)$	1.3832(14)			
	<b1-o></b1-o>	1.3657(19)			

#### Single-crystal XRD data of sample *a* refined with fully occupied positions and resolution restriction

**Table A1.9:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample *a* with resolution restriction  $(0.6 \text{ Å} \leq \lambda/(2 \cdot \sin \theta) \leq 1 \text{ Å})$  and fully occupied positions.

	Diffractometer: Enraf Nonius CAD4							
X-ray radiation	$MoK\alpha (0.71073 \text{ Å})$	Reflections rejected	1434					
X-ray power	50 kV, 40 mA	Reflections used	4294					
Crystal size	$0.2 \cdot 0.13 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.93^{\circ}$					
Measurement time	Max. 120 $sec/step$	Index range $h$	-9 9					
Absorption corr.	$\psi$ scans	Index range $k$	-24 24					
Resolution	$1$ Å $\dots$ 0.6 Å	Index range $l$	-12 12					
Sample name	$a \ (ABO-01-2009, \ \# \ 01)$	Unique reflections	1166					
Temperature	298K	Reflections > $4\sigma(I)$	1045					
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0510					
<i>a</i> -axis	5.6682(2) Å	$R(\mathbf{\sigma})$	0.0375					
<i>b</i> -axis	15.0060(9) Å	Refined parameters	82					
c-axis	7.6892(4) Å	GoF	0.993					
Volume	$654.07(6) \text{ Å}^3$	$R1,I>4{f \sigma}(I)$	0.0174					
Z	4	R1, all data	0.0274					
Density $\rho$	$2.942~{ m g/cm^3}$	w $R2$ on $F^2$	0.0296					
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.24 \text{ e}^{-} \cdot \text{Å}^{-3}$ near B1					
Reflections collected	5728	$\Delta  ho_{\min}$	-0.21 e <sup>-</sup> ·Å <sup>-3</sup> near O1					

**Table A1.10:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24818(6)	0.383521(18)	0.0	1.0	8b	0.00388(6)
Al2	0.0	0.25538(3)	0.68663(7)	1.0	4a	0.00513(8)
Al3	0.0	0.44420(3)	0.68198(9)	1.0	4a	0.00450(8)
Al4	0.0	0.29699(3)	0.33427(9)	1.0	4a	0.00408(8)
B1	0.5	0.51566(10)	0.22105(19)	1.0	4a	0.0049(2)
01	0.0	0.32947(8)	0.87054(15)	1.0	4a	0.00647(16)
O2	0.24291(12)	0.30983(5)	0.20164(11)	1.0	$8\mathrm{b}$	0.00663(12)
O3	0.5	0.45338(8)	0.09228(14)	1.0	4a	0.00542(16)
04	0.0	0.54602(7)	0.57211(14)	1.0	4a	0.00587(16)
O5	0.0	0.35181(7)	0.53798(14)	1.0	4a	0.00741(17)
O6	0.0	0.19126(7)	0.43504(16)	1.0	4a	0.00507(16)
07	0.28862(14)	0.45174(5)	0.78567(11)	1.0	8b	0.00670(12)

resolution	$(0.0 \text{ M} \leq \kappa/(2 \text{ sind}) \leq 1 \text{ M})$ and tany occupied positions.						
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{1\mathcal{Q}}$ (Å <sup>2</sup> )	
Al1	0.00300(9)	0.00439(10)	0.00424(9)	0.00053(10)	-0.00006(8)	0.00021(8)	
Al2	0.00457(16)	0.00455(15)	0.00626(16)	-0.00024(14)	0.0	0.0	
Al3	0.00354(16)	0.00386(14)	0.00610(16)	-0.00050(14)	0.0	0.0	
Al4	0.00414(15)	0.00470(16)	0.00341(15)	0.00027(12)	0.0	0.0	
B1	0.0038(5)	0.0060(5)	0.0048(5)	-0.0007(4)	0.0	0.0	
01	0.0040(3)	0.0085(4)	0.0069(4)	-0.0035(3)	0.0	0.0	
O2	0.0055(3)	0.0079(3)	0.0065(3)	0.0021(2)	0.0011(2)	0.00132(19)	
O3	0.0045(4)	0.0058(4)	0.0060(4)	-0.0025(3)	0.0	0.0	
O4	0.0035(3)	0.0056(4)	0.0085(4)	0.0026(3)	0.0	0.0	
O5	0.0113(4)	0.0053(4)	0.0055(4)	-0.0007(3)	0.0	0.0	
O6	0.0040(4)	0.0037(4)	0.0074(4)	0.0013(3)	0.0	0.0	
07	0.0045(3)	0.0084(3)	0.0071(2)	0.0035(2)	-0.00067(19)	-0.0010(2)	

**Table A1.11:** Anisotropic displacement parameters from single-crystal XRD data of sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

**Table A1.12:** Individual and mean polyhedral bond lengths of the refined structure from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O4	1.8452(7)	Al2	O2(x2)	1.7593(7)
Al1	O6	1.8832(7)	Al2	01	1.7989(11)
Al1	O2	1.9046(8)	Al2	O5	1.8440(12)
Al1	01	1.9048(8)	Al2	O6	2.1606(13)
Al1	O3	1.9079(8)		<al2-o></al2-o>	1.8644(10)
Al1	07	1.9535(8)			
	<al1-o></al1-o>	1.8999(8)	Al3	O4	1.7459(12)
			Al3	O5	1.7743(11)
Al4	O2 (x2)	1.7243(8)	Al3	O7 $(x2)$	1.8235(8)
Al4	O6	1.7658(11)	Al3	01	2.2507(13)
Al4	O5	1.7692(12)		<al3-o></al3-o>	1.8836(10)
	<al4-o></al4-o>	1.7459(10)			
B1	O3	1.3615(17)			
B1	O7 $(x2)$	1.3864(10)			
	<b1-0></b1-0>	1.3968(15)			

#### Single-crystal XRD data of sample *f* refined with fully occupied positions and resolution restriction

**Table A1.13:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

	Diffractometer: Bruker Smart APEX 2							
X-ray radiation	$MoK\alpha$ (0.71073 Å)	Reflections rejected	3505					
X-ray power	50  kV, 40  mA	Reflections used	5915					
Crystal size	$0.15 \cdot 0.025 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.94^{\circ}$					
Measurement time	60  s/frame	Index range $h$	-9 9					
Absorption corr.	Empirical $\psi$ scans	Index range $k$	-24 23					
Resolution	$1 \ \mathrm{\AA} \ \ldots \ 0.62 \ \mathrm{\AA}$	Index range $l$	-12 11					
Sample name	f (ABO-62, # 62-2)	Unique reflections	1130					
Temperature	298K	Reflections $> 4\sigma(I)$	1088					
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0399					
a-axis*	5.6618(7) Å	$R(\mathbf{\sigma})$	0.0294					
b-axis*	14.9981(12) Å	Refined parameters	82					
c-axis*	7.6806(7) Å	GoF	1.043					
$Volume^*$	652.21(12) Å <sup>3</sup>	$R1,I>4{f \sigma}(I)$	0.0187					
Z	4	R1, all data	0.0201					
Density $\rho$	$2.250~{ m g/cm^3}$	wR2 on $F^2$	0.0375					
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.25~{\rm e}^{\rm -}{\rm \cdot}{\rm \AA}^{\rm -3}$ near Al2					
Reflections collected	9420	$\Delta  ho_{\min}$	-0.25 e <sup>-</sup> ·Å <sup>-3</sup> near All					

\*lattice parameters from CAD4 (15 centered reflections with  $\pm 11^{\circ} < \theta < \pm 30^{\circ}$ )

	1 9	(	- / (	/ _ /		1 1
Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24810(6)	0.38347(2)	0.0	1.0	8b	0.00474(8)
Al2	0.0	0.25544(3)	0.68639(8)	1.0	4a	0.00620(10)
Al3	0.0	0.44419(3)	0.68167(9)	1.0	4a	0.00543(10)
Al4	0.0	0.29688(3)	0.33416(9)	1.0	4a	0.00559(10)
B1	0.5	0.51556(10)	0.2212(2)	1.0	4a	0.0057(2)
O1	0.0	0.32932(8)	0.87043(16)	1.0	4a	0.00683(18)
O2	0.24280(13)	0.30975(6)	0.20137(13)	1.0	8b	0.00796(15)
O3	0.5	0.45336(8)	0.09288(16)	1.0	4a	0.00615(19)
O4	0.0	0.54600(8)	0.57220(17)	1.0	4a	0.00703(19)
O5	0.0	0.35172(8)	0.53799(16)	1.0	4a	0.0084(2)
O6	0.0	0.19123(7)	0.43511(17)	1.0	4a	0.00595(18)
07	0.28870(15)	0.45178(6)	0.78534(12)	1.0	8b	0.00727(14)
a117 1	G :			<u>,                                     </u>		

**Table A1.14:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

resolution	$(0.02 \text{ M} \le \kappa/(2 \text{ sind}) \le 1 \text{ M})$ and fully occupied positions.							
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{1\mathcal{Z}}$ (Å <sup>2</sup> )		
Al1	0.00358(12)	0.00544(12)	0.00522(12)	-0.00004(12)	-0.00014(8)	0.00037(8)		
Al2	0.00540(17)	0.00570(17)	0.00750(18)	-0.00008(15)	0.0	0.0		
Al3	0.00404(17)	0.00507(16)	0.00717(18)	-0.00011(15)	0.0	0.0		
Al4	0.00565(17)	0.00607(18)	0.00506(17)	0.00014(14)	0.0	0.0		
B1	0.0048(5)	0.0069(6)	0.0054(6)	-0.0012(4)	0.0	0.0		
01	0.0039(4)	0.0093(4)	0.0074(4)	-0.0034(3)	0.0	0.0		
O2	0.0072(3)	0.0086(3)	0.0080(3)	0.0019(2)	0.0015(2)	0.0007(2)		
O3	0.0048(4)	0.0065(4)	0.0071(4)	-0.0024(3)	0.0	0.0		
O4	0.0046(4)	0.0064(4)	0.0101(5)	0.0025(3)	0.0	0.0		
O5	0.0115(5)	0.0072(4)	0.0064(4)	-0.0015(3)	0.0	0.0		
O6	0.0039(4)	0.0056(4)	0.0083(4)	0.0015(3)	0.0	0.0		
07	0.0040(3)	0.0097(3)	0.0081(3)	0.0039(2)	0.0005(2)	-0.0012(2)		

**Table A1.15:** Anisotropic displacement parameters from single-crystal XRD data of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

**Table A1.16:** Individual and mean polyhedral bond lengths of the refined structure from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and fully occupied positions.

Center	Ligand	Distance $(\text{\AA})$	Center	Ligand	Distance $(\text{\AA})$
Al1	O4	1.8438(8)	Al2	O2 (x2)	1.7578(8)
Al1	O6	1.8808(8)	Al2	01	1.7961(12)
Al1	O2	1.9014(9)	Al2	O5	1.8396(13)
Al1	01	1.9034(8)	Al2	O6	2.1569(13)
Al1	O3	1.9084(8)		<al2-o></al2-o>	1.8616(11)
Al1	O7	1.9547(9)			
	<al1-o></al1-o>	1.8988(8)	A13	O4	1.7431(12)
			Al3	O5	1.7725(12)
Al4	O2 (x2)	1.7226(9)	Al3	O7 (x2)	1.8217(9)
Al4	O6	1.7642(12)	Al3	O1	2.2517(13)
Al4	O5	1.7684(13)		<al3-o></al3-o>	1.8821(11)
	<al4-o></al4-o>	1.7445(11)			
B1	O3	1.3570(19)			
B1	O7 $(x2)$	1.3834(11)			
	<b1-o></b1-o>	1.3746(14)			

# Single-crystal XRD data of sample *a* refined with 91% AI and 9% B at the AI4 site according to refinement strategy 2

**Table A1.17:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample a a without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

	Diffractometer: Enraf Nonius CAD4							
X-ray radiation	$MoK\alpha (0.71073 \text{ Å})$	Reflections rejected	92					
X-ray power	50  kV, 40  mA	Reflections used	5636					
Crystal size	$0.2\cdot0.13\cdot0.05~\mathrm{mm^3}$	Max. 20	$69.93^{\circ}$					
Measurement time	Max. $120 \text{ sec/step}$	Index range $h$	-9 9					
Absorption corr.	$\psi$ scans	Index range $k$	-24 24					
Resolution	No restriction	Index range $l$	-12 12					
Sample name	a  (ABO-01-2009,  # 01)	Unique reflections	1552					
Temperature	298K	$\operatorname{Reflections} > 4\sigma(I)$	1423					
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0373					
<i>a</i> -axis	5.6682(2) Å	$R(\mathbf{\sigma})$	0.0265					
b-axis	15.0060(9) Å	Refined parameters	82					
c-axis	7.6892(4) Å	GoF	1.239					
Volume	654.07(6) Å <sup>3</sup>	$R1,I>4{ m s}(I)$	0.0205					
Z	4	R1, all data	0.0268					
Density $\rho$	$2.942~{ m g/cm^3}$	w $R2$ on $F^2$	0.0533					
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.54~{\rm e}^{-}{\rm \cdot \AA^{-3}}$ near O1					
Reflections collected	5728	$\Delta  ho_{\min}$	-0.35 e <sup>-</sup> ·Å <sup>-3</sup> near Al2					

**Table A1.18:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample a without resolution restriction and fully occupied positions (refinement strategy 2).

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24813(7)	0.38355(3)	0.0	1.0	8b	0.00459(7)
Al2	0.0	0.25534(4)	0.68658(10)	1.0	4a	0.00600(10)
Al3	0.0	0.44419(4)	0.68193(11)	1.0	4a	0.00527(10)
Al4	0.0	0.29693(4)	0.33414(12)	0.91	4a	0.00254(10)
$\rm B^{b}$	0.0	0.29693(4)	0.33414(12)	0.09	4a	0.00254(10)
B1	0.5	0.51579(14)	0.2211(3)	1.0	4a	0.0056(3)
01	0.0	0.32913(10)	0.8708(2)	1.0	4a	0.0073(2)
O2	0.24309(16)	0.30963(7)	0.20124(16)	1.0	8b	0.00708(18)
O3	0.5	0.45347(10)	0.0923(2)	1.0	4a	0.0063(2)
O4	0.0	0.54587(9)	0.5717(2)	1.0	4a	0.0066(2)
O5	0.0	0.35173(9)	0.53787(19)	1.0	4a	0.0081(3)
O6	0.0	0.19132(9)	0.4350(2)	1.0	4a	0.0056(2)
07	0.28860(18)	0.45167(7)	0.78561(15)	1.0	8b	0.00734(17)

2).						
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00365(13)	0.00519(14)	0.00492(13)	0.00045(16)	-0.00002(12)	0.00008(12)
Al2	0.0053(2)	0.0055(2)	0.0072(2)	-0.0001(2)	0.0	0.0
Al3	0.0043(2)	0.0047(2)	0.0068(2)	-0.0005(2)	0.0	0.0
Al4	0.0025(2)	0.0032(2)	0.0019(2)	0.00039(19)	0.0	0.0
$\mathbf{B}^{\mathbf{a}}$	0.0025(2)	0.0032(2)	0.0019(2)	0.00039(19)	0.0	0.0
B1	0.0053(8)	0.0065(8)	0.0048(8)	0.0004(6)	0.0	0.0
01	0.0040(5)	0.0094(6)	0.0085(6)	-0.0036(5)	0.0	0.0
O2	0.0061(4)	0.0079(4)	0.0072(4)	0.0022(3)	0.0013(3)	0.0011(3)
O3	0.0058(5)	0.0063(6)	0.0069(6)	-0.0026(5)	0.0	0.0
O4	0.0043(5)	0.0058(6)	0.0098(6)	0.0025(5)	0.0	0.0
O5	0.0115(6)	0.0057(6)	0.0070(6)	-0.0006(4)	0.0	0.0
O6	0.0052(5)	0.0042(5)	0.0074(6)	0.0008(4)	0.0	0.0
07	0.0055(4)	0.0089(4)	0.0076(4)	0.0036(3)	-0.0009(3)	-0.0008(3)

**Table A1.19:** Anisotropic displacement parameters from single-crystal XRD data of sample a without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

**Table A1.20:** Individual and mean polyhedral bond lengths of the refined structure from sample a a without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O4	1.8451(9)	Al2	O2 (x2)	1.7562(10)
Al1	O6	1.8843(9)	Al2	01	1.7978(14)
Al1	O2	1.9041(11)	Al2	O5	1.8437(15)
Al1	O1	1.9058(10)	A12	O6	2.1599(17)
Al1	O3	1.9087(10)		<al2-o></al2-o>	1.8628(13)
Al1	07	1.9532(11)			
	<al1-o></al1-o>	1.9002(10)	Al3	O4	1.7453(15)
			Al3	O5	1.7753(15)
Al4	O2 (x2)	1.7261(10)	Al3	O7 $(x2)$	1.8233(10)
Al4	O6	1.7643(15)	Al3	01	2.2560(17)
Al4	O5	1.7693(16)		<al3-o></al3-o>	1.8846(13)
	<al4-o></al4-o>	1.7465(13)			
B1	O3	1.362(2)			
B1	O7 $(x2)$	1.3858(14)			
	<b1-o></b1-o>	1.3779(16)			

# Single-crystal XRD data of sample *f* refined with 91% AI and 9% B at the AI4 site according to refinement strategy 2

**Table A1.21:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample f without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

	Diffractometer: Bruker Smart APEX 2							
X-ray radiation	$MoK\alpha$ (0.71073 Å)	Reflections rejected	170					
X-ray power	50  kV, 40  mA	Reflections used	9250					
Crystal size	$0.15 \cdot 0.025 \cdot 0.05 \text{ mm}^3$	Max. 20	$73.38^{\circ}$					
Measurement time	$60 \mathrm{s/frame}$	Index range $h$	-9 9					
Absorption corr.	Empirical $\psi$ scans	Index range $k$	$-24 \dots 25$					
Resolution	No restriction	Index range $l$	-12 11					
Sample name	f (ABO-62, # 62-2)	Unique reflections	1618					
Temperature	298K	Reflections $> 4\sigma(I)$	1565					
Space group	$Cmc2_1$	R(int)	0.0258					
a-axis*	5.6618(7) Å	$R(\sigma)$	0.0194					
b-axis*	14.9981(12) Å	Refined parameters	82					
c-axis*	7.6806(7) Å	GoF	1.186					
$Volume^*$	652.21(12) Å <sup>3</sup>	$R1,I>4{f \sigma}(I)$	0.0234					
Z	4	R1, all data	0.0245					
Density $\rho$	$2.250 \text{ g/cm}^3$	w $R2$ on $F^2$	0.0629					
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.47~{\rm e}^{\rm -}{\rm \cdot}{\rm \AA}^{\rm -3}$ near Al3					
Reflections collected	9420	$\Delta  ho_{\min}$	-0.49 e <sup>-</sup> ·Å <sup>-3</sup> near Al2					

\*lattice parameters from CAD4 (15 centered reflections with  $\pm 11^{\circ} < \theta < \pm 30^{\circ}$ )

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24810(7)	0.38350(3)	0.0	1.0	8b	0.00578(8)
Al2	0.0	0.25539(4)	0.68638(11)	1.0	4a	0.00748(11)
Al3	0.0	0.44418(4)	0.68161(12)	1.0	4a	0.00659(11)
Al4	0.0	0.29687(4)	0.33412(12)	0.91	4a	0.00435(11)
$\rm B^b$	0.0	0.29687(4)	0.33412(12)	0.09	4a	0.00435(11)
B1	0.5	0.51543(13)	0.2212(3)	1.0	4a	0.0069(3)
01	0.0	0.32895(9)	0.8708(2)	1.0	4a	0.0077(2)
O2	0.24297(17)	0.30945(7)	0.20104(18)	1.0	8b	0.00867(19)
O3	0.5	0.45346(9)	0.0927(2)	1.0	4a	0.0073(2)
04	0.0	0.54593(9)	0.5718(2)	1.0	4a	0.0079(2)
O5	0.0	0.35185(9)	0.5379(2)	1.0	4a	0.0093(3)
O6	0.0	0.19127(9)	0.4350(2)	1.0	4a	0.0071(2)
07	0.28874(18)	0.45165(7)	0.78539(16)	1.0	8b	0.00828(18)

**Table A1.22:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample f without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

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Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00456(15)	0.00660(15)	0.00618(15)	0.00001(15)	-0.00011(12)	0.00030(11)
Al2	0.0064(2)	0.0071(2)	0.0089(2)	0.0(2)	0.0	0.0
Al3	0.0054(2)	0.0063(2)	0.0081(2)	-0.0001(2)	0.0	0.0
Al4	0.0043(2)	0.0049(2)	0.0038(2)	0.0002(2)	0.0	0.0
$\mathbf{B}^{\mathbf{a}}$	0.0043(2)	0.0049(2)	0.0038(2)	0.0002(2)	0.0	0.0
B1	0.0063(7)	0.0081(8)	0.0061(8)	-0.0002(6)	0.0	0.0
01	0.0046(5)	0.0095(6)	0.0090(6)	-0.0037(5)	0.0	0.0
O2	0.0081(4)	0.0091(4)	0.0088(4)	0.0018(3)	0.0015(3)	0.0006(3)
O3	0.0064(5)	0.0068(5)	0.0087(6)	-0.0022(5)	0.0	0.0
O4	0.0058(5)	0.0065(6)	0.0112(6)	0.0027(5)	0.0	0.0
O5	0.0119(6)	0.0076(6)	0.0083(6)	-0.0007(5)	0.0	0.0
O6	0.0055(5)	0.0065(5)	0.0092(6)	0.0010(5)	0.0	0.0
07	0.0053(4)	0.0107(4)	0.0088(4)	0.0040(3)	-0.0004(3)	-0.0009(3)

**Table A1.23:** Anisotropic displacement parameters from single-crystal XRD data of sample f without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

**Table A1.24:** Individual and mean polyhedral bond lengths of the refined structure from sample f without resolution restriction and a fixed ratio of 91% Al and 9% B at the Al4 site (refinement strategy 2).

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O4	1.8433(9)	Al2	O2 (x2)	1.7539(10)
Al1	O6	1.8817(9)	Al2	01	1.7954(15)
Al1	O2	1.9023(12)	Al2	O5	1.8422(16)
Al1	01	1.9045(10)	Al2	O6	2.1570(17)
Al1	O3	1.9085(10)		<al2-o></al2-o>	1.8605(14)
Al1	07	1.9531(11)			
	<al1-o></al1-o>	1.9057(10)	Al3	O4	1.7434(15)
			Al3	O5	1.7710(15)
Al4	O2 (x2)	1.7242(11)	Al3	O7 (x2)	1.8222(11)
Al4	O6	1.7631(15)	Al3	O1	2.2579(16)
Al4	O5	1.7689(17)		<al3-o></al3-o>	1.8833(14)
	<al4-o></al4-o>	1.7451(14)			
B1	O3	1.356(3)			
B1	O7 $(x2)$	1.3846(14)			
	<b1-o></b1-o>	1.375(2)			

### Single-crystal XRD data of sample *a* refined with 91% AI and 9% B at the AI4 site and resolution restriction

**Table A1.25:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

	Diffractometer: Enr	af Nonius CAD4	
X-ray radiation	Mo $K\alpha$ (0.71073 Å)	Reflections rejected	1434
X-ray power	50 kV, 40 mA	Reflections used	4294
Crystal size	$0.2 \cdot 0.13 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.93^{\circ}$
Measurement time	Max. $120 \text{ sec/step}$	Index range $h$	-9 9
Absorption corr.	$\psi$ scans	Index range $k$	-24 24
Resolution	$1$ Å $\dots$ 0.6 Å	Index range $l$	-12 12
Sample name	a  (ABO-01-2009,  # 01)	Unique reflections	1166
Temperature	298K	Reflections $> 4\sigma(I)$	1045
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0510
<i>a</i> -axis	5.6682(2) Å	$R(\sigma)$	0.0375
b-axis	15.0060(9) Å	Refined parameters	82
c-axis	7.6892(4) Å	GoF	0.822
Volume	654.07(6) Å <sup>3</sup>	$R1,I>4{ m s}(I)$	0.0810
Z	4	R1, all data	0.0279
Density $\rho$	$2.942~{ m g/cm^3}$	w $R2$ on $F^2$	0.0385
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.27 \text{ e}^{-} \cdot \text{Å}^{-3}$ near B1
Reflections collected	5728	$\Delta  ho_{\min}$	-0.19 e <sup>-</sup> ·Å <sup>-3</sup> near O6

**Table A1.26:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24815(6)	0.38353(2)	0.0	1.0	8b	0.00423(8)
Al2	0.0	0.25538(3)	0.68661(8)	1.0	4a	0.00541(9)
Al3	0.0	0.44417(3)	0.68198(10)	1.0	4a	0.00487(9)
Al4	0.0	0.29700(3)	0.33424(10)	0.91	4a	0.00254(10)
$\rm B^b$	0.0	0.29700(3)	0.33424(10)	0.09	4a	0.00254(10)
B1	0.5	0.51573(11)	0.2210(2)	1.0	4a	0.0053(2)
01	0.0	0.32949(9)	0.87047(17)	1.0	4a	0.00683(19)
O2	0.24289(14)	0.30987(6)	0.20158(12)	1.0	8b	0.00692(14)
O3	0.5	0.45335(9)	0.09233(16)	1.0	4a	0.00574(19)
O4	0.0	0.54595(8)	0.57201(17)	1.0	4a	0.00624(19)
O5	0.0	0.35171(8)	0.53810(17)	1.0	4a	0.0077(2)
O6	0.0	0.19119(8)	0.43512(19)	1.0	4a	0.00533(18)
07	0.28858(15)	0.45171(6)	0.78567(12)	1.0	8b	0.00699(14)

Site $U_{11}$ (Å2) $U_{22}$ (Å2) $U_{33}$ (Å2) $U_{23}$ (Å2) $U_{13}$ (Å2) $U_{12}$ (Å2)Al10.00342(11)0.00473(12)0.00454(11)0.00051(12) $-0.00006(9)$ 0.00019(9)Al20.00477(18)0.00491(17)0.00654(18) $-0.00022(16)$ 0.00.0Al30.00385(17)0.00416(16)0.00660(18) $-0.00057(16)$ 0.00.0Al40.00276(16)0.00310(18)0.00178(17)0.00024(14)0.00.0Ba0.00276(16)0.0061(6)0.0054(6) $-0.0003(4)$ 0.00.0B10.0043(5)0.0061(6)0.0072(4) $-0.0032(4)$ 0.00.0O10.0045(4)0.0088(4)0.0072(4) $-0.0032(4)$ 0.00.0O20.0055(3)0.0081(3)0.0071(3)0.0022(2)0.0012(2)0.0014(2)O30.0051(4)0.0060(4) $-0.0026(3)$ 0.00.0O40.0037(4)0.0060(4) $0.0078(4)$ 0.000.0O50.0118(5)0.0053(4)0.0078(4)0.0013(3)0.00.0O60.0042(4)0.0040(4)0.0078(4)0.0013(3)0.00.0O70.0047(3)0.0087(3)0.0076(3)0.0035(2) $-0.0008(2)$ $-0.0009(2)$							
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Al1	0.00342(11)	0.00473(12)	0.00454(11)	0.00051(12)	-0.00006(9)	0.00019(9)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Al2	0.00477(18)	0.00491(17)	0.00654(18)	-0.00022(16)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Al3	0.00385(17)	0.00416(16)	0.00660(18)	-0.00057(16)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Al4	0.00276(16)	0.00310(18)	0.00178(17)	0.00024(14)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\mathbf{B}^{\mathbf{a}}$	0.00276(16)	0.00310(18)	0.00178(17)	0.00024(14)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	B1	0.0043(5)	0.0061(6)	0.0054(6)	-0.0003(4)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	O1	0.0045(4)	0.0088(4)	0.0072(4)	-0.0032(4)	0.0	0.0
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	O2	0.0055(3)	0.0081(3)	0.0071(3)	0.0022(2)	0.0012(2)	0.0014(2)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	O3	0.0051(4)	0.0061(4)	0.0060(4)	-0.0026(3)	0.0	0.0
O6   0.0042(4)   0.0040(4)   0.0078(4)   0.0013(3)   0.0   0.0	O4	0.0037(4)	0.0060(4)	0.0090(5)	0.0027(4)	0.0	0.0
	O5	0.0118(5)	0.0053(4)	0.0060(4)	-0.0010(3)	0.0	0.0
O7   0.0047(3)   0.0087(3)   0.0076(3)   0.0035(2)   -0.0008(2)   -0.0009(2)	O6	0.0042(4)	0.0040(4)	0.0078(4)	0.0013(3)	0.0	0.0
	07	0.0047(3)	0.0087(3)	0.0076(3)	0.0035(2)	-0.0008(2)	-0.0009(2)

**Table A1.27:** Anisotropic displacement parameters from single-crystal XRD data of sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

**Table A1.28:** Individual and mean polyhedral bond lengths of the refined structure from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

Center	Ligand	Distance $(Å)$	Center	Ligand	Distance (Å)
Al1	O4	1.8453(8)	Al2	O2(x2)	1.7596(8)
Al1	O6	1.8827(8)	Al2	01	1.7987(12)
Al1	O2	1.9040(9)	Al2	O5	1.8422(14)
Al1	01	1.9048(9)	Al2	O6	2.1604(14)
Al1	O3	1.9079(9)		<al2-o></al2-o>	1.8641(11)
Al1	07	1.9532(9)			
	<al1-o></al1-o>	1.8997(9)	Al3	O4	1.7458(13)
			Al3	O5	1.7745(13)
Al4	O2 (x2)	1.7244(9)	Al3	O7 $(x2)$	1.8233(9)
Al4	O6	1.7671(12)	A13	01	2.2499(14)
Al4	O5	1.7695(14)		<al3-o></al3-o>	1.8834(14)
	<al4-o></al4-o>	1.7464(11)			
B1	O3	1.3621(19)			
B1	O7 (x2)	1.3864(11)			
	<b1-o></b1-o>	1.3783(14)			

### Single-crystal XRD data of sample *f* refined with 91% AI and 9% B at the AI4 site and resolution restriction

**Table A1.29:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

	Diffractometer: Bruk	er Smart APEX 2	
X-ray radiation	Mo $K\alpha$ (0.71073 Å)	Reflections rejected	3505
X-ray power	50 kV, 40 mA	Reflections used	5915
Crystal size	$0.15 \cdot 0.025 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.94^{\circ}$
Measurement time	$60 \mathrm{s/frame}$	Index range $h$	-9 9
Absorption corr.	Empirical $\psi$ scans	Index range $k$	-24 23
Resolution	$1 \ \mathrm{\AA} \ \ldots \ 0.62 \ \mathrm{\AA}$	Index range $l$	-12 11
Sample name	f (ABO-62, # 62-2)	Unique reflections	1130
Temperature	298K	Reflections $> 4\sigma(I)$	1088
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0399
a-axis*	5.6618(7) Å	$R(\mathbf{\sigma})$	0.0294
b-axis*	14.9981(12) Å	Refined parameters	82
c-axis*	7.6806(7) Å	GoF	0.815
$Volume^*$	652.21(12) Å <sup>3</sup>	$R1,I>4{f \sigma}(I)$	0.0192
Z	4	R1, all data	0.0205
Density $\rho$	$2.250~{ m g/cm^3}$	w $R2$ on $F^2$	0.0463
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.25~{ m e}^{-1} \cdot { m A}^{-3}$ near Al2
Reflections collected	9420	$\Delta  ho_{\min}$	-0.25 e <sup>-</sup> ·Å <sup>-3</sup> near Al1

\*lattice parameters from CAD4 (15 centered reflections with  $\pm 11^{\circ} < \theta < \pm 30^{\circ}$ )

**Table A1.30:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24809(6)	0.38347(2)	0.0	1.0	8b	0.00516(10)
Al2	0.0	0.25544(3)	0.68638(9)	1.0	4a	0.00654(11)
Al3	0.0	0.44419(3)	0.68166(10)	1.0	4a	0.00585(11)
Al4	0.0	0.29690(3)	0.33413(10)	0.91	4a	0.00404(11)
$\rm B^{b}$	0.0	0.29690(3)	0.33413(10)	0.09	4a	0.00404(11)
B1	0.5	0.51562(11)	0.2212(2)	1.0	4a	0.0062(3)
01	0.0	0.32925(8)	0.87053(18)	1.0	4a	0.0072(2)
O2	0.24286(14)	0.30974(6)	0.20135(14)	1.0	8b	0.00822(16)
O3	0.5	0.45338(8)	0.09282(17)	1.0	4a	0.0066(2)
O4	0.0	0.54594(8)	0.57215(19)	1.0	4a	0.0074(2)
O5	0.0	0.35170(8)	0.53801(18)	1.0	4a	0.0088(2)
O6	0.0	0.19118(8)	0.43510(19)	1.0	4a	0.00643(19)
07	0.28852(16)	0.45173(6)	0.78528(13)	1.0	8b	0.00763(15)

5100.						
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00408(13)	0.00584(13)	0.00555(14)	-0.00008(12)	-0.00011(9)	0.00038(8)
Al2	0.00567(19)	0.00616(18)	0.0078(2)	-0.00016(16)	0.0	0.0
Al3	0.00460(19)	0.00544(18)	0.0075(2)	-0.00017(16)	0.0	0.0
Al4	0.00417(18)	0.00451(19)	0.00343(18)	0.00008(15)	0.0	0.0
$\mathbf{B}^{\mathbf{a}}$	0.00417(18)	0.00451(19)	0.00343(18)	0.00008(15)	0.0	0.0
B1	0.0054(6)	0.0072(6)	0.0061(6)	-0.0012(4)	0.0	0.0
01	0.0042(4)	0.0097(5)	0.0077(5)	-0.0032(4)	0.0	0.0
O2	0.0075(3)	0.0089(3)	0.0082(3)	0.0020(2)	0.0014(2)	0.0007(2)
O3	0.0051(4)	0.0070(4)	0.0077(5)	-0.0025(4)	0.0	0.0
O4	0.0048(4)	0.0069(5)	0.0103(5)	0.0025(4)	0.0	0.0
O5	0.0118(5)	0.0078(4)	0.0068(5)	-0.0015(4)	0.0	0.0
O6	0.0044(4)	0.0060(4)	0.0089(5)	0.0014(4)	0.0	0.0
07	0.0043(3)	0.0099(3)	0.0087(3)	0.0041(3)	-0.0005(2)	-0.0012(2)

**Table A1.31:** Anisotropic displacement parameters from single-crystal XRD data of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

**Table A1.32:** Individual and mean polyhedral bond lengths of the refined structure from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and a fixed ratio of 91% Al and 9% B at the Al4 site.

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O4	1.8442(8)	Al2	O2(x2)	1.7573(9)
Al1	O6	1.8805(8)	A12	01	1.7961(13)
Al1	O2	1.9014(10)	Al2	O5	1.8394(15)
Al1	01	1.9034(9)	Al2	O6	2.1572(15)
Al1	O3	1.9083(9)		<al2-o></al2-o>	1.8615(12)
Al1	07	1.9546(10)			
	<al1-o></al1-o>	1.8987(9)	Al3	O4	1.7424(13)
			Al3	O5	1.7724(13)
Al4	O2(x2)	1.7228(10)	Al3	O7(x2)	1.8206(9)
Al4	O6	1.7650(13)	Al3	O1	2.2531(14)
Al4	O5	1.7686(14)		<al3-o></al3-o>	1.8818(12)
	<al4-o></al4-o>	1.7448(12)			
B1	O3	1.358(2)			
B1	O7(x2)	1.3841(12)			
	<b1-o></b1-o>	1.3754(15)			

# Single-crystal XRD data of sample *a* refined with vacancies at the Al2 site and resolution restriction according to refinement strategy 3

**Table A1.33:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

	Diffractometer: Enr	af Nonius CAD4	
X-ray radiation	$MoK\alpha (0.71073 \text{ Å})$	Reflections rejected	1434
X-ray power	50  kV, 40  mA	Reflections used	4294
Crystal size	$0.2 \cdot 0.13 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.93^{\circ}$
Measurement time	Max. $120 \text{ sec/step}$	Index range $h$	-9 9
Absorption corr.	$\psi$ scans	Index range $k$	-24 24
Resolution	$1 \text{ \AA} \dots 0.6 \text{ \AA}$	Index range $l$	-12 12
Sample name	$a \ (ABO-01-2009, \ \# \ 01)$	Unique reflections	1166
Temperature	298K	$\operatorname{Reflections} > 4\sigma(I)$	1045
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0510
a-axis	5.6682(2) Å	$R(\mathbf{\sigma})$	0.0375
<i>b</i> -axis	15.0060(9) Å	Refined parameters	83
c-axis	7.6892(4) Å	GoF	0.990
Volume	654.07(6) Å <sup>3</sup>	$R1,I>4{ m s}(I)$	0.0173
Ζ	4	R1, all data	0.0273
Density $\rho$	$2.942 \mathrm{~g/cm^3}$	w $R2$ on $F^2$	0.0295
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.54~{\rm e}^{\rm \cdot}{\rm \cdot}{\rm \AA}^{\rm -3}$ near Al2
Reflections collected	5728	$\Delta  ho_{\min}$	-0.29 e <sup>-</sup> ·Å <sup>-3</sup> near Al1

**Table A1.34:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24819(6)	0.383523(18)	0.0	1.0	8b	0.00397(7)
Al2	0.0	0.25538(3)	0.68660(7)	0.979(4)	4a	0.0046(2)
Al3	0.0	0.44420(3)	0.68197(9)	1.0	4a	0.00462(9)
Al4	0.0	0.29699(3)	0.33426(9)	1.0	4a	0.00419(9)
B1	0.5	0.51567(10)	0.22105(19)	1.0	4a	0.0050(2)
O1	0.0	0.32949(8)	0.87053(15)	1.0	4a	0.00659(17)
O2	0.24288(12)	0.30984(5)	0.20162(11)	1.0	8b	0.00672(13)
O3	0.5	0.45337(8)	0.09229(14)	1.0	4a	0.00551(17)
O4	0.0	0.54602(7)	0.57211(14)	1.0	4a	0.00596(16)
O5	0.0	0.35183(7)	0.53796(14)	1.0	4a	0.00750(18)
O6	0.0	0.19126(7)	0.43506(16)	1.0	4a	0.00516(17)
07	0.28862(14)	0.45173(5)	0.78568(11)	1.0	8b	0.00678(12)

resolution	$(0.0 \text{ M} \leq \kappa/(2 \text{ sind}) \leq 1 \text{ M})$ and $(1.2 \text{ vacancies (remember strategy 5)}).$						
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )	
Al1	0.00308(10)	0.00447(11)	0.00435(10)	0.00051(10)	-0.00005(8)	0.00022(8)	
Al2	0.0040(3)	0.0040(3)	0.0057(3)	-0.00022(14)	0.0	0.0	
Al3	0.00365(16)	0.00397(15)	0.00623(16)	-0.00050(14)	0.0	0.0	
Al4	0.00422(15)	0.00485(17)	0.00350(15)	0.00026(12)	0.0	0.0	
B1	0.0040(5)	0.0060(5)	0.0050(5)	-0.0007(4)	0.0	0.0	
O1	0.0042(3)	0.0087(4)	0.0069(4)	-0.0035(3)	0.0	0.0	
O2	0.0055(3)	0.0080(3)	0.0066(3)	0.0021(2)	0.0010(2)	0.00133(19)	
O3	0.0046(4)	0.0059(4)	0.0061(4)	-0.0025(3)	0.0	0.0	
O4	0.0036(3)	0.0058(4)	0.0085(4)	0.0027(3)	0.0	0.0	
O5	0.0114(4)	0.0054(4)	0.0056(4)	-0.0007(3)	0.0	0.0	
O6	0.0041(4)	0.0038(4)	0.0076(4)	0.0013(3)	0.0	0.0	
07	0.0046(3)	0.0085(3)	0.0072(2)	0.0035(2)	-0.00066(19)	-0.0010(2)	

**Table A1.35:** Anisotropic displacement parameters from single-crystal XRD data of sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

**Table A1.36:** Individual and mean polyhedral bond lengths of the refined structure from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Center	Ligand	Distance $(\text{Å})$	Center	Ligand	Distance $(\text{\AA})$
Al1	O4	1.8451(7)	Al2	O2 (x2)	1.7594(7)
Al1	O6	1.8832(7)	Al2	01	1.7991(11)
Al1	O2	1.9044(8)	Al2	O5	1.8441(12)
Al1	01	1.9047(8)	Al2	O6	2.1603(13)
Al1	O3	1.9078(8)		<al2-o></al2-o>	1.8645(10)
Al1	07	1.9534(8)			
	<al1-o></al1-o>	1.8998(8)	Al3	O4	1.7459(12)
			Al3	O5	1.7741(11)
Al4	O2 (x2)	1.7242(8)	Al3	O7 (x2)	1.8236(8)
Al4	O6	1.7658(11)	Al3	01	2.2506(13)
Al4	O5	1.7693(12)		<al3-o></al3-o>	1.8836(10)
	<al4-o></al4-o>	1.7459(10)			
B1	O3	1.3616(17)			
B1	O7 $(x2)$	1.3864(10)			
	<b1-o></b1-o>	1.3782(12)			

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
Al1-O1-Al1 (x2)	95.23(5)	Al3-O5-Al2	103.08(5)
Al1-O1-Al3	90.64(4)	Al4-O5-Al2	100.58(6)
Al2-O1-Al1 (x2)	132.38(2)	Al4-O5-Al3	156.33(7)
Al2-O1-Al3	88.07(4)	Al1-O6-Al1	98.58(5)
Al2-O2-Al1	104.86(4)	Al1-O6-Al2 $(x2)$	91.60(4)
Al4-O2-Al1	123.99(4)	Al4-O6-Al1 $(x2)$	130.68(3)
Al4-O2-Al2	129.65(4)	Al4-O6-Al2	89.59(5)
Al1-O3-Al1	96.87(5)	Al3-O7-Al1	103.37(4)
B1-O3-Al1 $(x2)$	130.36(3)	B1-O7-Al1	126.06(7)
Al1-O4-Al1	99.37(5)	B1-O7-Al3	129.83(7)
Al3-O4-Al1 (x2)	130.31(3)		

**Table A1.37:** M-O-M angles of the refined structure from sample *a* with resolution restriction (0.6  $\dot{A} \leq \lambda/(2 \cdot \sin \theta) \leq 1 \dot{A}$ ) and Al2 vacancies (refinement strategy 3).

**Table A1.38:** O-M-O angles of the refined structure from sample *a* with resolution restriction (0.6 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
O1-Al1-O3	168.80(5)	O2-Al2-O2	111.87(5)
O1-Al1-O7	82.46(4)	O2-A12-O5 (x2)	118.51(3)
O2-Al1-O1	99.59(4)	O2-Al2-O6 (x2)	79.11(3)
O2-Al1-O3	91.61(4)	O5-A12-O6	78.15(5)
O2-Al1-O7	173.12(4)	O4-Al3-O1	168.83(5)
O3-Al1-O7	86.45(4)	O4-A13-O5	112.44(5)
O4-Al1-O1	80.66(3)	O4-Al3-O7 (x2)	99.05(3)
O4-Al1-O2	94.36(4)	O5-Al3-O1	78.73(5)
O4-Al1-O3	98.28(4)	O5-Al3-O7 (x2)	108.75(3)
O4-Al1-O6	177.57(6)	O7-Al3-O1 (x2)	76.45(3)
O4-Al1-O7	92.45(4)	O7-Al3-O7	127.58(6)
O6-Al1-O1	99.65(4)	O2-Al4-O2	105.98(5)
O6-Al1-O2	83.21(4)	O2-A14-O5 (x2)	118.14(3)
O6-Al1-O3	81.89(4)	O2-A14-O6~(x2)	111.11(3)
O6-Al1-O7	89.98(5)	O6-A14-O5	91.68(5)
O1-Al2-O5	90.12(6)	O3-B1-O7 (x2)	120.19(6)
O1-Al2-O6	168.27(5)	O7-B1-O7	119.60(11)
O2-Al2-O1 $(x2)$	106.99(3)		

### Single-crystal XRD data of sample *f* refined with vacancies at the Al2 site and resolution restriction according to refinement strategy 3

**Table A1.39:** Details for single-crystal X-ray diffraction data collection and crystal structure refinement of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Diffractometer: Bruker Smart APEX 2							
X-ray radiation	Mo $K\alpha$ (0.71073 Å)	Reflections rejected	3505				
X-ray power	50 kV, 40 mA	Reflections used	5915				
Crystal size	$0.15 \cdot 0.025 \cdot 0.05 \text{ mm}^3$	Max. 20	$69.94^{\circ}$				
Measurement time	$60 \mathrm{s/frame}$	Index range $h$	-9 9				
Absorption corr.	Empirical $\psi$ scans	Index range $k$	$-24 \ldots 23$				
Resolution	$1 \text{ \AA} \dots 0.62 \text{ \AA}$	Index range $l$	-12 11				
Sample name	f (ABO-62, # 62-2)	Unique reflections	1130				
Temperature	298K	$ ext{Reflections} > 4\sigma(I)$	1088				
Space group	$Cmc2_1$	$R(\mathrm{int})$	0.0399				
a-axis*	5.6618(7) Å	$R(\sigma)$	0.0294				
b-axis*	14.9981(12) Å	Refined parameters	83				
c-axis*	7.6806(7) Å	GoF	1.046				
$Volume^*$	652.21(12) Å <sup>3</sup>	$R1,I>4{f \sigma}(I)$	0.0186				
Ζ	4	R1, all data	0.0199				
Density $\rho$	$2.250~{ m g/cm^3}$	w $R2$ on $F^2$	0.0370				
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$\Delta  ho_{ m max}$	$0.24 \text{ e}^{-} \cdot \text{Å}^{-3}$ near Al2				
Reflections collected	9420	$\Delta  ho_{\min}$	-0.24 e··Å-³ near Al1				

\*lattice parameters from CAD4 (15 centered reflections with  $\pm 11^\circ < \theta < \pm 30^\circ)$ 

**Table A1.40:** Atomic positions and isotropic displacement parameters of single-crystal XRD data from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Site	x	y	z	Occ.	W. S. <sup>a</sup>	$U_{eq.}$ (Å <sup>2</sup> )
Al1	0.24810(6)	0.38347(2)	0.0	1.0	8b	0.00489(9)
Al2	0.0	0.25544(3)	0.68636(8)	0.968(4)	4a	0.0053(3)
Al3	0.0	0.44420(3)	0.68166(9)	1.0	4a	0.00558(11)
Al4	0.0	0.29688(3)	0.33415(9)	1.0	4a	0.00576(11)
B1	0.5	0.51556(10)	0.2212(2)	1.0	4a	0.0059(2)
O1	0.0	0.32935(8)	0.87041(16)	1.0	4a	0.00700(19)
O2	0.24275(13)	0.30977(6)	0.20136(13)	1.0	8b	0.00810(15)
O3	0.5	0.45335(8)	0.09289(16)	1.0	4a	0.00630(19)
O4	0.0	0.54600(8)	0.57221(17)	1.0	4a	0.00718(19)
O5	0.0	0.35174(8)	0.53797(16)	1.0	4a	0.0085(2)
O6	0.0	0.19125(7)	0.43511(17)	1.0	4a	0.00608(18)
07	0.28868(15)	0.45178(6)	0.78535(12)	1.0	8b	0.00741(15)

resolution restriction $(0.02 \text{ A} \le k/(2.500) \le 1 \text{ A})$ and Al2 vacancies (remember strategy 5).							
Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{1\mathcal{Q}}$ (Å <sup>2</sup> )	
Al1	0.00371(12)	0.00555(12)	0.00540(13)	-0.00006(11)	-0.00013(8)	0.00037(8)	
Al2	0.0046(3)	0.0048(3)	0.0066(3)	-0.00006(15)	0.0	0.0	
Al3	0.00419(18)	0.00520(17)	0.00735(19)	-0.00011(15)	0.0	0.0	
Al4	0.00577(17)	0.00629(19)	0.00522(17)	0.00013(14)	0.0	0.0	
B1	0.0050(5)	0.0069(6)	0.0057(6)	-0.0012(4)	0.0	0.0	
01	0.0040(4)	0.0096(4)	0.0074(4)	-0.0034(3)	0.0	0.0	
O2	0.0074(3)	0.0088(3)	0.0081(3)	0.0019(2)	0.0014(2)	0.0007(2)	
O3	0.0049(4)	0.0066(4)	0.0073(4)	-0.0023(3)	0.0	0.0	
O4	0.0048(4)	0.0067(4)	0.0100(5)	0.0025(3)	0.0	0.0	
O5	0.0116(5)	0.0073(4)	0.0065(4)	-0.0015(3)	0.0	0.0	
O6	0.0041(4)	0.0057(4)	0.0085(4)	0.0015(3)	0.0	0.0	
07	0.0042(3)	0.0099(3)	0.0082(3)	0.0038(2)	-0.0005(2)	-0.0012(2)	

**Table A1.41:** Anisotropic displacement parameters from single-crystal XRD data of sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

**Table A1.42:** Individual and mean polyhedral bond lengths of the refined structure from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Center	Ligand	Distance $(\text{Å})$	Center	Ligand	Distance $(\text{\AA})$
Al1	O4	1.8439(8)	Al2	O2(x2)	1.7581(8)
Al1	O6	1.8810(8)	Al2	01	1.7964(12)
Al1	O2	1.9012(9)	Al2	O5	1.8398(13)
Al1	01	1.9034(8)	Al2	O6	2.1566(13)
Al1	O3	1.9082(8)		<al2-o></al2-o>	1.8617(11)
Al1	07	1.9546(9)			
	<al1-o></al1-o>	1.8987(8)	Al3	O4	1.7429(12)
			Al3	O5	1.7723(12)
Al4	O2 (x2)	1.7224(9)	Al3	O7 $(x2)$	1.8217(9)
Al4	O6	1.7639(12)	Al3	O1	2.2514(13)
Al4	O5	1.7685(13)		<al3-o></al3-o>	1.8820(11)
	<al4-o></al4-o>	1.7443(11)			
B1	O3	1.3570(18)			
B1	O7 $(x2)$	1.3836(11)			
	<b1-o></b1-o>	1.3747(14)			

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1 (x2)	95.13(5)	Al3-O5-Al2	103.21(6)
Al1-O1-Al3	90.60(5)	Al4-O5-Al2	100.56(6)
Al2-O1-Al1 (x2)	132.43(3)	Al4-O5-Al3	156.24(8)
Al2-O1-Al3	88.01(5)	Al1-O6-Al1	98.61(5)
Al2-O2-Al1	104.89(4)	Al1-O6-Al2 (x2)	91.65(5)
Al4-O2-Al1	124.03(5)	Al4-O6-Al1 (x2)	130.66(3)
Al4-O2-Al2	129.60(5)	Al4-O6-Al2	89.56(5)
Al1-O3-Al1	96.73(5)	Al3-O7-Al1	103.33(4)
B1-O3-Al1 (x2)	130.47(3)	B1-O7-Al1	126.00(7)
All-O4-All	99.25(5)	B1-O7-Al3	129.95(7)
Al3-O4-Al1 (x2)	130.37(3)		

**Table A1.43:** M-O-M angles of the refined structure from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

**Table A1.44:** O-M-O angles of the refined structure from sample f with resolution restriction (0.62 Å  $\leq \lambda/(2 \cdot \sin \theta) \leq 1$  Å) and Al2 vacancies (refinement strategy 3).

Sites	Angle (°)	Sites	Angle ( $^{\circ}$ )
O1-Al1-O3	168.91(5)	O2-Al2-O2	111.87(6)
O1-Al1-O7	82.49(5)	O2-Al2-O5 $(x2)$	118.51(3)
O2-Al1-O1	99.53(5)	O2-Al2-O6 (x2)	79.06(4)
O2-Al1-O3	91.55(5)	O5-Al2-O6	78.24(5)
O2-Al1-O7	173.10(4)	O4-Al3-O1	168.75(6)
O3-Al1-O7	86.54(5)	O4-Al3-O5	112.65(6)
O4-Al1-O1	80.78(4)	O4-Al3-O7 (x2)	98.98(3)
O4-Al1-O2	94.40(5)	O5-Al3-O1	78.60(5)
O4-Al1-O3	98.15(4)	O5-Al3-O7 (x2)	108.73(3)
O4-Al1-O6	177.58(7)	O7-Al3-O1 (x2)	76.48(3)
O4-Al1-O7	92.44(5)	O7-Al3-O7	127.59(6)
O6-Al1-O1	99.61(4)	O2-Al4-O2	105.87(6)
O6-Al1-O2	83.18(5)	O2-Al4-O5 (x2)	118.16(4)
O6-Al1-O3	81.93(4)	O2-Al4-O6 (x2)	111.17(4)
O6-Al1-O7	89.98(5)	O6-Al4-O5	91.64(6)
O1-Al2-O5	90.18(6)	O3-B1-O7 (x2)	120.14(6)
O1-Al2-O6	168.41(5)	O7-B1-O7	119.70(12)
O2-Al2-O1 $(x2)$	106.96(4)		

#### Bond valences of sample *a*

	0.							
Site	O1	O2	O3	O4	O5	O6	07	$\Sigma_{\rm Cation}$
Al1	$0.462^{a}$	0.464	$0.457^{\rm a}$	$0.544^{\rm a}$		$0.489^{a}$	0.406	2.822
Al2	0.617	$0.691^{\rm b}$			0.543	0.232		2.774
Al3	0.181			0.711	0.661		$0.577^{ m b}$	2.707
Al4		$0.752^{\rm b}$			0.664	0.679		2.847
B1			1.021				$0.966^{\mathrm{b}}$	2.953
$\Sigma_{\rm Anion}$	1.722	1.907	1.935	1.799	1.868	1.889	1.949	

**Table A1.45:** Bond valences calculated according to the structure refinement of sample a according to effinement strategy 1.

<sup>a</sup>bonds to two cations, <sup>b</sup>bonds to two anions

#### **FTIR correlation table**

**Table A1.46:** Correlation table of vibrational species of ideal and distorted borate groups, modified after Wilson et al. (1980).

Vibrations of the $(BO_3)^{3-}$ group	Ideal symmetry $D3h (-62m)$	Effective symmetry <sup>a</sup> $C2v \ (mm2)$	Site symmetry <sup>b</sup> Cs (m)
$v_1$ , symmetrical stretching	$A_1$ '	$A_1$	A'
$v_3$ , antisymmetrical stretching	E'	$A_1 + B_2$	A' + A"
$v_2$ , out-of-plane bending	$A_2$ "	$B_1$	A'
$v_4$ , within-plane bending	E'	$A_1 + B_2$	A' + A"

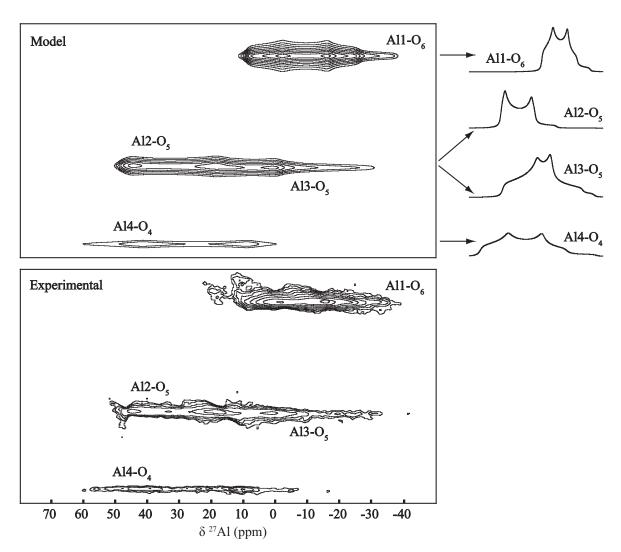
<sup>a</sup> according to the real bond lengths B-O3 (1.36 Å) and 2 x B-O7 (1.38 Å) and bond angles of  $120^{\circ}$  the borate group has the shape of an isosceles triangle, b<sup>b</sup>Wyckoff site symmetry 4a of the B and O3 atoms

#### <sup>27</sup>AI MAS NMR parameters and <sup>27</sup>AI 3Q-MAS NMR spectra

Site	Species <sup>a</sup>	$\delta^{27} \mathrm{Al}_\mathrm{iso}~(\mathrm{ppm})$	$C_q$ (MHz)	η
Al1	$Al^{(VI)}$	11.9	5.91	0.39
Al2	$\mathrm{Al^{(V)}}_{1}$	56.0	6.65	0.11
Al3	$\mathrm{Al}^{(\mathrm{V})}{}_{2}$	48.6	7.78	0.71
Al4	$\mathrm{Al}^{\mathrm{(IV)}}$	75.2	9.18	0.41

Table A1.47: Chemical shifts and quadrupolar parameters obtained by line shape analysis of the  $^{27}$ Al 3Q-MAS NMR spectrum of sample *a*.

<sup>a</sup>nomenclature in agreement with Gan et al. (2002) and Hung et al. (2006)



**Figure A1.1:** Experimental and modeled <sup>27</sup>Al 3Q-MAS NMR spectra of sample *a*. The chemical shifts and the quadrupolar parameters of the four aluminum species were evaluated by fitting the 2D NMR spectrum using the program Dmfit (Massiot et al. 2002). The simulated 1D slices of the four species are shown at the right margin of the 2D plot.

#### A.2 Supplementary Data for Chapter 5

#### Supplementary data for $AI_5BO_9$

Table A2.1: Details for	single-crystal	X-ray	$\operatorname{diffraction}$	data	$\operatorname{collection}$	and	$\operatorname{crystal}$	structure	refine-
ment of $Al_5BO_9$ .									

Diffractometer: B	Bruker Smart APEX 2, $MoK\alpha$ (0.71073 .	Å), 50 kV, 40 mA
Crystal size	$0.13 \cdot 0.07$	$\cdot 0.4 \text{ mm}^3$
Detector distance	59.5	mm
Rotation axis	φ,	ω
Collection mode	Automated	hemisphere
Frame size	$512 \cdot 51$	2 pixel
Absorption corr.	Empirical $\psi$	correction
Resolution	1 Å	0.62 Å
Sample name	ABO-01	I-CAP
Temperature	$298 K^{a}$	873K
Space group	$Cmc2_1$	$Cmc2_1$
a-axis	5.6689(6) Å	5.6825(6) Å
<i>b</i> -axis	15.0045(13) Å	15.0754(16) Å
<i>c</i> -axis	7.6897(13) Å	7.7121 Á
Volume	654.07(13) Å <sup>3</sup>	$660.67 \text{ Å}^3$
Z	4	4
Density $\rho$	$2.942~{ m g/cm^3}$	$2.913~{ m g/cm^3}$
Absorption coeff. $\mu$	$0.88 \text{ mm}^{-1}$	$0.87 \text{ mm}^{-1}$
Time per frame	10 s	30 s
Reflections collected	8857	3199
Reflections rejected	3938	1922
Reflections used	4919	1277
Max. 20	$69.91^{\circ}$	$64.64^{\circ}$
Index range $h$	-8 8	-7 7
Index range $k$	-23 23	-21 20
Index range <i>l</i>	-12 12	-10 10
Unique reflections	1110	595
Reflections $> 4\sigma(I)$	1118	581
R(int)	0.0423	0.0414
$R(\sigma)$	0.0349	0.0529
Refined parameters	83	83
GoF	1.001	1.076
$R1,I>4{ m s}(I)$	0.0195	0.0242
R1, all data	0.0211	0.0247
w $R2$ on $F^2$	0.0351	0.0557
$\Delta\rho_{\rm max}$	$0.23 e^{-} \cdot A^{-3}$ , near B1	$0.26 e^{-}$ Å <sup>-3</sup> , near Al1
$\Delta \rho_{\min}$	$-0.18 e^{-\dot{A}-3}$ , near O4	-0.26 e <sup>-</sup> ·Å <sup>-3</sup> , near Al1

<sup>a</sup>lattice parameters from CAD4

Site	T (K)	x	y	z	$U_{eq.}$ (Å <sup>2</sup> )
Al1	298	0.24807(7)	0.38356(2)	0.0	0.00437(10)
	873	0.24823(11)	0.38348(5)	0.0	0.0104(3)
Al2	$298_{\rm a}$	0.0	0.25537(3)	0.68676(8)	0.0052(3)
	$873_{ m b}$	0.0	0.25535(6)	0.68865(16)	0.0114(6)
Al3	298	0.0	0.44416(3)	0.68216(9)	0.00515(11)
	873	0.0	0.44438(6)	0.68045(19)	0.0110(3)
Al4	298	0.0	0.29692(3)	0.33421(9)	0.00471(11)
	873	0.0	0.29686(6)	0.33490(17)	0.0102(3)
B1	298	0.5	0.51545(11)	0.2209(2)	0.0057(3)
	873	0.5	0.5158(2)	0.2203(4)	0.0102(6)
01	298	0.0	0.32948(8)	0.87041(16)	0.0071(2)
	873	0.0	0.32915(18)	0.8714(3)	0.0146(5)
O2	298	0.24297(15)	0.30993(6)	0.20155(12)	0.00698(16)
	873	0.2432(3)	0.30959(15)	0.2026(3)	0.0156(4)
O3	298	0.5	0.45351(8)	0.09240(15)	0.0056(2)
	873	0.5	0.45346(18)	0.0929(3)	0.0137(5)
O4	298	0.0	0.54584(8)	0.57239(16)	0.0065(2)
	873	0.0	0.54610(17)	0.5721(4)	0.0156(5)
O5	298	0.0	0.35176(8)	0.53855(16)	0.0078(2)
	873	0.0	0.35135(16)	0.5391(3)	0.0179(5)
O6	298	0.0	0.19121(8)	0.43486(17)	0.0057(2)
	873	0.0	0.19109(15)	0.4356(3)	0.0123(4)
07	298	0.28875(17)	0.45168(6)	0.78551(12)	0.00705(15)
	873	0.2893(4)	0.45132(13)	0.7836(2)	0.0159(4)

Table A2.2: Atomic positions and isotropic displacement parameters of  $Al_5BO_9$ .

<sup>a</sup>occupancy: 0.983(10) <sup>b</sup>occupancy: 0.99(2)

**Table A2.3:** Anisotropic displacement parameters of  $\mathrm{Al}_5\mathrm{BO}$  at 298K.

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00320(14)	0.00509(13)	0.00482(13)	0.00049(12)	0.00012(9)	0.00026(9)
Al2	0.0045(3)	0.0050(3)	0.0060(3)	0.00001(15)	0.0	0.0
Al3	0.0042(2)	0.00460(17)	0.00663(18)	-0.00047(15)	0.0	0.0
Al4	0.0046(2)	0.00523(19)	0.00428(17)	0.00034(14)	0.0	0.0
B1	0.0048(6)	0.0061(6)	0.0063(6)	-0.0005(4)	0.0	0.0
01	0.0052(5)	0.0089(4)	0.0073(4)	-0.0037(4)	0.0	0.0
O2	0.0058(4)	0.0087(3)	0.0065(3)	0.0022(2)	0.0012(2)	0.0012(2)
O3	0.0040(5)	0.0066(5)	0.0062(4)	-0.0027(3)	0.0	0.0
O4	0.0048(5)	0.0054(5)	0.0094(4)	0.0023(3)	0.0	0.0
O5	0.0107(5)	0.0066(4)	0.0059(4)	-0.0009(3)	0.0	0.0
O6	0.0046(5)	0.0054(4)	0.0070(4)	0.0011(3)	0.0	0.0
07	0.0051(3)	0.0090(3)	0.0071(3)	0.0033(2)	-0.0006(2)	-0.0004(3)

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.0073(4)	0.0145(4)	0.0094(3)	0.0006(3)	-0.00054(14)	0.00090(15)
Al2	0.0090(7)	0.0130(7)	0.0120(7)	-0.0003(3)	0.0	0.0
Al3	0.0086(5)	0.0124(5)	0.0119(4)	-0.0004(3)	0.0	0.0
Al4	0.0084(4)	0.0155(5)	0.0067(4)	0.0013(3)	0.0	0.0
B1	0.0034(14)	0.0160(15)	0.0113(12)	-0.0034(9)	0.0	0.0
O1	0.0077(10)	0.0232(11)	0.0130(11)	-0.0083(8)	0.0	0.0
O2	0.0105(8)	0.0223(9)	0.0139(7)	0.0065(5)	0.0048(4)	0.0031(4)
O3	0.0088(11)	0.0197(13)	0.0125(9)	-0.0066(9)	0.0	0.0
O4	0.0107(11)	0.0161(11)	0.0199(10)	0.0089(9)	0.0	0.0
O5	0.0266(14)	0.0167(11)	0.0103(9)	-0.0015(8)	0.0	0.0
O6	0.0090(10)	0.0135(9)	0.0143(9)	0.0027(7)	0.0	0.0
07	0.0089(7)	0.0242(9)	0.0147(6)	0.0089(6)	0.0006(5)	-0.0026(5)

Table A2.4: Anisotropic displacement parameters of Al<sub>5</sub>BO at 873K.

**Table A2.5:** Polyhedral bond lengths of  $Al_5BO_9$ .

Center	Ligand	$d_{298K}$ (Å)	$d_{873K}~({\rm \AA})$	α (K <sup>-1</sup> )	$ \mathrm{esd}  \mathrel{\alpha}(\mathrm{K}^{\text{-}1})$
Al1	01	1.9050(9)	1.9091(17)	3.743E-06	1.756E-06
	O2	1.9035(9)	1.919(2)	1.416E-05	2.087 E-06
	O3	1.9095(9)	1.9165(17)	6.375E-06	1.752 E-06
	O4	1.8464(8)	1.8511(16)	4.333E-06	1.685 E-06
	O6	1.8839(8)	1.8861(16)	2.031E-06	1.651E-06
	07	1.9541(9)	1.9710(18)	1.504E-05	1.791E-06
	<al1-o></al1-o>	1.9004(9)	1.909(2)	7.672E-06	1.787 E-06
Al2	O1	1.7974(12)	1.795(2)	-2.129E-06	2.510 E-06
	O2	1.7596(9)	1.7603(19)	6.919E-07	1.467 E-06
	O2	1.7596(9)	1.7603(19)	6.919E-07	2.078E-06
	O5	1.8413(14)	1.851(3)	8.878E-06	2.957 E-06
	O6	2.1632(14)	2.179(3)	1.262E-05	2.304 E-06
	<Al2 <sup>(V)</sup> -O>	1.8642(11)	1.869(2)	4.534E-06	2.268 E-06
Al3	01	2.2486(14)	2.277(3)	2.204E-05	2.421E-06
	O4	1.7436(13)	1.746(3)	2.593E-06	2.899 E-06
	O5	1.7725(13)	1.776(3)	3.827E-06	2.852 E-06
	07	1.8231(10)	1.830(2)	6.201E-06	2.133E-06
	07	1.8231(10)	1.830(2)	6.201E-06	2.133E-06
	<Al3 <sup>(V)</sup> -O>	1.8822(12)	1.892(3)	8.870E-06	2.479E-06
Al4	O2	1.7250(9)	1.7287(17)	3.730E-06	1.939E-06
	O3	1.7250(9)	1.7287(17)	3.730E-06	1.939E-06
	O5	1.7737(13)	1.776(3)	2.255E-06	2.938E-06
	O6	1.7650(13)	1.774(2)	8.474E-06	2.689E-06
	<al4-o></al4-o>	1.7472(11)	1.752(2)	4.554E-06	2.382 E-06
B1	O3	1.3567(19)	1.360(4)	3.846E-06	5.446E-06
	07	1.3870(12)	1.385(2)	-2.884E-06	3.142E-06
	07	1.3870(12)	1.385(2)	-2.884E-06	3.142E-06
	<b1-0></b1-0>	1.3769(14)	1.377(3)	-6.736E-07	3.898E-06

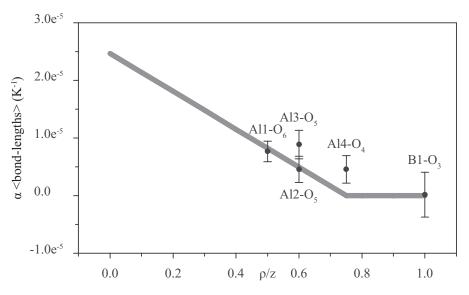


Figure A2.1: Expansion parameters  $\alpha$  calculated from averaged polyhedral bond length of Al<sub>5</sub>BO<sub>9</sub>. Expected values according to Hazen and Prewitt (1977) plot on the gray line.

Sites	Angle $(^{\circ})$	Sites	Angle (°)
Al1-O1-Al1	95.15(6)	Al3-O5-Al2	103.22(6)
Al1-O1-Al3 (x2)	90.62(5)	Al3-O5-Al4	156.18(8)
Al2-O1-Al1 (x2)	132.42(3)	Al4-O5-Al2	100.60(6)
Al2-O1-Al3	88.14(5)	Al1-O6-Al1	98.60(6)
Al2-O2-Al1	104.95(4)	Al1-O6-Al2 $(x2)$	91.54(5)
Al4-O2-Al1	124.01(5)	Al4-O6-Al1 (x2)	130.67(3)
Al4-O2-Al2	129.52(5)	Al4-O6-Al2	89.58(6)
Al1-O3-Al1	96.82(6)	Al3-O7-Al1	103.27(5)
B1-O3-Al1 (x2)	130.36(3)	B1-O7-Al1	126.16(8)
Al1-O4-Al1	99.20(6)	B1-O7-Al3	129.88(8)
Al3-O4-Al1 (x2)	130.39(3)		

Table A2.6: M-O-M angles of the refined structure of  $Al_5BO_9$  at 298K.

Table A2.7: M-O-M angles of the refined structure of  $\rm Al_5BO_9$  at 873K.

	0	5 9	
Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	95.27(11)	Al3-O5-Al2	103.58(12)
Al1-O1-Al3 (x2)	90.50(10)	Al3-O5-Al4	155.41(16)
Al2-O1-Al1 (x2)	132.35(5)	Al4-O5-Al2	101.01(13)
Al2-O1-Al3	88.01(10)	Al1-O6-Al1	98.67(11)
Al2-O2-Al1	105.12(8)	Al1-O6-Al2 $(x2)$	91.66(9)
Al4-O2-Al1	123.83(10)	Al4-O6-Al1 $(x2)$	130.63(5)
Al4-O2-Al2	129.57(12)	Al4-O6-Al2	89.57(11)
Al1-O3-Al1	96.58(11)	Al3-O7-Al1	103.42(10)
B1-O3-Al1 (x2)	130.58(6)	B1-O7-Al1	125.94(16)
Al1-O4-Al1	99.29(11)	B1-O7-Al3	130.08(14)
Al3-O4-Al1 (x2)	130.35(6)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
O1-Al1-O3	168.78(5)	O2-Al2-O2	111.80(6)
O1-Al1-O7	82.44(5)	O2-A12-O5 (x2)	118.52(3)
O2-Al1-O1	99.64(5)	O2-Al2-O6 (x2)	79.05(3)
O2-Al1-O3	91.57(5)	O5-Al2-O6	78.19(5)
O2-Al1-O7	173.04(5)	O4-Al3-O1	168.88(6)
O3-Al1-O7	86.46(5)	O4-Al3-O5	112.51(6)
O4-Al1-O1	80.79(4)	O4-Al3-O7 (x2)	99.02(4)
O4-Al1-O2	94.35(5)	O5-Al3-O1	78.61(5)
O4-Al1-O3	98.17(4)	O5-Al3-O7 (x2)	108.66(3)
O4-Al1-O6	177.60(7)	O7-Al3-O1 (x2)	76.51(3)
O4-Al1-O7	92.54(5)	O7-Al3-O7	127.76(6)
O6-Al1-O1	99.61(4)	O2-Al4-O2	105.96(6)
O6-Al1-O2	83.25(5)	O2-Al4-O5 (x2)	118.12(4)
O6-Al1-O3	81.91(5)	O2-Al4-O6 (x2)	111.16(4)
O6-Al1-O7	89.86(5)	O6-Al4-O5	91.63(6)
O1-Al2-O5	90.02(6)	O3-B1-O7 (x2)	120.29(7)
O1-Al2-O6	168.21(6)	O7-B1-O7	119.40(13)
O2-Al2-O1 $(x2)$	107.08(4)		

**Table A2.8:** O-M-O angles of the refined structure of  $Al_5BO_9$  at 298K.

Table A2.9: O-M-O angles of the refined structure of  $\rm Al_5BO_9$  at 873K.

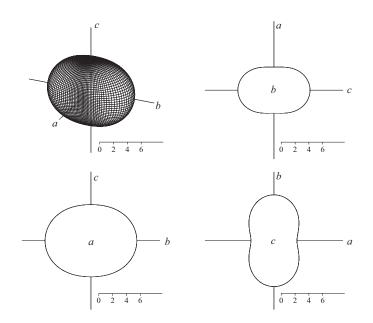
Sites	Angle (°)	Sites	Angle (°)
O1-Al1-O3	169.09(12)	O2-A12-O2	111.96(15)
O1-Al1-O7	82.55(11)	O2-Al2-O5 (x2)	118.23(8)
O2-Al1-O1	99.39(11)	O2-Al2-O6 (x2)	78.90(7)
O2-Al1-O3	91.51(10)	O5-Al2-O6	77.84(11)
O2-Al1-O7	172.84(9)	O4-Al3-O1	168.29(13)
O3-Al1-O7	86.68(10)	O4-Al3-O5	113.56(13)
O4-Al1-O1	80.75(8)	O4-Al3-O7 (x2)	99.08(7)
O4-Al1-O2	94.41(11)	O5-Al3-O1	78.15(10)
O4-Al1-O3	98.09(9)	O5-Al3-O7 (x2)	108.19(7)
O4-Al1-O6	177.50(14)	O7-Al3-O1 (x2)	76.26(6)
O4-Al1-O7	92.70(11)	O7-Al3-O7	127.96(13)
O6-Al1-O1	99.66(9)	O2-Al4-O2	106.18(14)
O6-Al1-O2	83.09(9)	O2-Al4-O5~(x2)	118.16(8)
O6-Al1-O3	81.97(8)	O2-Al4-O6 (x2)	110.99(9)
O6-Al1-O7	89.79(11)	O6-Al4-O5	91.58(12)
O1-Al2-O5	90.27(13)	O3-B1-O7 (x2)	120.16(13)
O1-Al2-O6	168.11(12)	O7-B1-O7	119.7(3)
O2-Al2-O1 $(x2)$	107.26(8)		

T (K)	a-axis (Å)	<i>b</i> -axis (Å)	<i>c</i> -axis (Å)	Volume (Å)	$R_{wp}$	GoF
	~ /	~ /	. ,			
298	5.66891(4)	15.01477(12)	7.69560(5)	655.029(9)	7.51	1.14
323	5.66936(4)	15.01715(11)	7.69656(5)	655.267(8)	7.54	1.14
373	5.67029(3)	15.02191(9)	7.69849(4)	655.747(7)	7.54	1.14
423	5.67125(3)	15.02667(7)	7.70045(3)	656.233(6)	7.59	1.15
473	5.67225(2)	15.03144(6)	7.70242(3)	656.724(5)	7.57	1.14
523	5.67327(2)	15.03621(6)	7.70442(3)	657.221(4)	7.50	1.13
573	5.67432(2)	15.04099(6)	7.70643(2)	657.724(4)	7.56	1.14
623	5.67540(2)	15.04578(6)	7.70846(2)	658.231(4)	7.63	1.15
673	5.67651(2)	15.05056(6)	7.71050(3)	658.745(4)	7.63	1.15
723	5.67766(2)	15.05536(6)	7.71257(3)	659.264(4)	7.64	1.15
773	5.67883(2)	15.06015(6)	7.71466(3)	659.788(4)	7.63	1.15
823	5.68003(2)	15.06496(6)	7.71676(3)	660.319(4)	7.71	1.16
873	5.68126(2)	15.06976(6)	7.71888(3)	660.854(4)	7.73	1.16
923	5.68253(2)	15.07457(6)	7.72102(3)	661.396(4)	7.69	1.15
973	5.68382(2)	15.07939(6)	7.72318(2)	661.942(4)	7.68	1.15
1023	5.68514(2)	15.08421(5)	7.72536(2)	662.495(4)	7.66	1.15
1073	5.68649(2)	15.08904(6)	7.72755(3)	663.053(4)	7.63	1.14
1123	5.68788(2)	15.09387(6)	7.72977(3)	663.617(5)	7.73	1.15
1173	5.68929(2)	15.09870(7)	7.73200(4)	664.186(6)	7.70	1.15
1223	5.69073(3)	15.10354(9)	7.73425(4)	664.761(7)	7.68	1.14
1273	5.69221(3)	15.10838(11)	7.73652(5)	665.342(8)	7.87	1.17

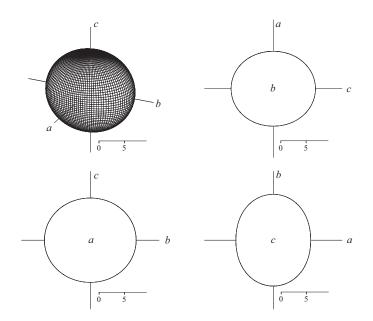
**Table A2.10:** Temperature dependent lattice parameters and agreement factors from powder XRD refinements of  $Al_5BO_9$ .

Table A2.11: Temperature dependent thermal expansion parameters  $\alpha$  of  ${\rm Al}_5{\rm BO}_9.$  AF is the anisotropy factor.

	( 1)	( 1)	( 1)	( 1)	·
T (K)	$\alpha_{a-axis}$ (K <sup>-1</sup> )	$\alpha_{b-axis}$ (K <sup>-1</sup> )	$\alpha_{c-axis}$ (K <sup>-1</sup> )	$\alpha_{Volume}$ (K <sup>-1</sup> )	AF
298	3.17949E-06	6.33479E-06	4.98378E-06	1.44906E-05	6.31
323	3.23237E-06	6.33687E-06	5.00730E-06	1.45702 E-05	6.21
373	3.33807 E-06	6.34101E-06	5.05430E-06	1.47291E-05	6.01
423	3.44372 E-06	6.34515E-06	5.10127 E-06	1.48876E-05	5.80
473	3.54932 E-06	6.34929E-06	5.14820 E-06	1.50458E-05	5.60
523	3.65486 E-06	6.35342 E-06	5.19510E-06	1.52036E-05	5.40
573	3.76035E-06	6.35755E-06	5.24195E-06	1.53611E-05	5.19
623	3.86577 E-06	6.36167 E-06	5.28878E-06	1.55182 E-05	4.99
673	3.97113E-06	6.36579E-06	5.33556E-06	1.56749E-05	4.79
723	4.07643 E-06	6.36990E-06	5.38231E-06	1.58312E-05	4.59
773	4.18167 E-06	6.37401E-06	5.42901E-06	1.59872E-05	4.38
823	4.28684 E-06	6.37811E-06	5.47568E-06	1.61427 E-05	4.18
873	4.39194 E-06	6.38221 E-06	5.52231E-06	1.62979E-05	3.98
923	4.49697 E-06	6.38631E-06	5.56891E-06	1.64527 E-05	3.78
973	4.60193 E-06	6.39040E-06	5.61546E-06	1.66071 E-05	3.58
1023	4.70682 E-06	6.39448E-06	5.66197 E-06	1.67611E-05	3.38
1073	4.81164 E-06	6.39857E-06	5.70844 E-06	1.69147 E-05	3.17
1123	4.91637 E-06	6.40264 E-06	5.75488E-06	1.70679E-05	2.97
1173	5.02103E-06	6.40671 E-06	5.80127 E-06	1.72207 E-05	2.77
1223	5.12562 E-06	6.41078 E-06	5.84762 E-06	1.73731E-05	2.57
1273	5.23012E-06	6.41484E-06	5.89393E-06	1.75251E-05	2.37



**Figure A2.2:** Representation surface of thermal expansion tensor components of  $Al_5BO_9$  at 298K. Top left: Arbitrary view, top right: Cross-section in (010), bottom left: Cross-section in (100) and bottom right: Cross-section in (001). *a*, *b* and *c* are directions parallel to unit cell axes, scale bar unit is  $10^{-6} \cdot K^{-1}$ .



**Figure A2.3:** Representation surface of thermal expansion tensor components of  $Al_5BO_9$  at 1273K. Top left: Arbitrary view, top right: Cross-section in (010), bottom left: Cross-section in (100) and bottom right: Cross-section in (001). *a*, *b* and *c* are directions parallel to unit cell axes, scale bar unit is  $10^{-6}$ ·K<sup>-1</sup>.

Plots were made with WinTensor v. 1.2 by Werner Kaminsky, Dept. of Chemistry, University of Washington.

#### Supplementary data for Al<sub>4</sub>B<sub>2</sub>O<sub>9</sub>

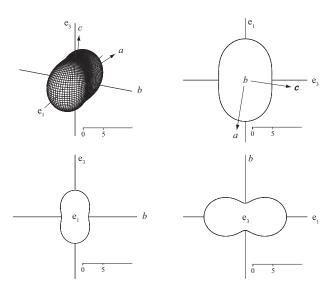
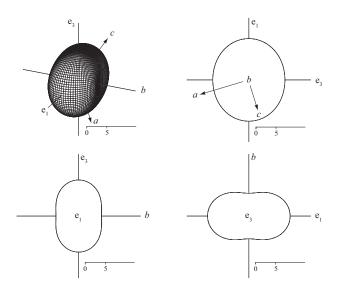


Figure A2.4: Representation surface of thermal expansion tensor components of  $Al_4B_2O_9$  at 298K after principal axis transformation.  $e_1$ ,  $e_2$  and  $e_3$  are eigenvectors of the unit cell axes, the *b*-axis is parallel to  $e_2$ . Top left: Arbitrary view, top right: Cross-section of the  $e_1$ - $e_3$ -plane, bottom left: Cross-section of the *b*- $e_3$ -plane and bottom right: Cross section of the *b*- $e_1$ -plane. *a*- and *c*-axes are plotted approximately. Scale bar unit is  $10^{-6} \cdot K^{-1}$ .



**Figure A2.5:** Representation surface of thermal expansion tensor components of  $Al_4B_2O_9$  at 1323K after principal axis transformation.  $e_1$ ,  $e_2$  and  $e_3$  are eigenvectors of the unit cell axes, the *b*-axis is parallel to  $e_2$ . Top left: Arbitrary view, top right: Cross-section of the  $e_1$ - $e_3$ -plane, bottom left: Cross-section of the *b*- $e_3$ -plane and bottom right: Cross section of the *b*- $e_1$ -plane. *a*- and *c*-axes are plotted approximately. Scale bar unit is  $10^{-6} \cdot K^{-1}$ .

Plots were made with WinTensor v. 1.2 by Werner Kaminsky, Dept. of Chemistry, University of Washington.

T (K)	<i>a</i> -axis (Å)	b-axis (Å)	c-axis (Å)	Volume (Å)	$\beta$ (°)	$R_{wp}$	GoF
298	14.8045(3)	5.54509(8)	15.07669(19)	90.8290(10)	1237.55(2)	9.50	1.24
323	14.8081(2)	5.54553(7)	15.07906(18)	90.8276(9)	1238.14(2)	9.78	1.27
373	14.8151(2)	5.54643(6)	15.08387(17)	90.8247(8)	1239.32(2)	9.83	1.27
423	14.8220(2)	5.54735(6)	15.08876(17)	90.8219(7)	1240.512(18)	9.88	1.28
473	14.82884(18)	5.54829(6)	15.09374(16)	90.8191(6)	1241.707(17)	10.02	1.30
523	14.83563(17)	5.54926(5)	15.09881(16)	90.8164(6)	1242.909(16)	9.94	1.28
573	14.84235(16)	5.55024(5)	15.10396(16)	90.8138(5)	1244.118(16)	10.07	1.30
623	14.84901(16)	5.55125(5)	15.10919(15)	90.8112(5)	1245.334(17)	9.98	1.28
673	14.85560(16)	5.55228(5)	15.11452(15)	90.8086(5)	1246.558(16)	10.03	1.29
723	14.86213(16)	5.55333(5)	15.11992(15)	90.8062(5)	1247.788(16)	10.10	1.30
773	14.86859(16)	5.5544(5)	15.12541(15)	90.8037(5)	1249.025(17)	10.05	1.29
823	14.87498(17)	5.55549(4)	15.13099(16)	90.8014(5)	1250.270(16)	10.28	1.32
873	14.8813(18)	5.55661(4)	15.13665(16)	90.7991(6)	1251.521(16)	10.02	1.28
923	14.88756(19)	5.55774(4)	15.14240(17)	90.7968(6)	1252.780(16)	10.12	1.29
973	14.8936(2)	5.55890(4)	15.14823(18)	90.7946(6)	1254.045(16)	10.19	1.30
1023	14.8999(2)	5.56008(4)	15.1542(2)	90.7925(7)	1255.318(16)	10.13	1.29
1073	14.9059(3)	5.56129(4)	15.1602(2)	90.7904(7)	1256.598(16)	10.03	1.28
1123	14.9119(3)	5.56251(5)	15.1662(3)	90.7884(8)	1257.886(17)	10.02	1.27
1173	14.9178(3)	5.56376(5)	15.1724(3)	90.7864(9)	1259.180(19)	10.27	1.30
1223	14.9237(4)	5.56502(6)	15.1787(3)	90.7845(10)	1260.48(2)	10.23	1.30
1273	14.92950(4)	5.56631(7)	15.1850(4)	90.7826(12)	1261.79(3)	10.12	1.28
1323	14.9352(5)	5.56762(8)	15.1914(4)	90.7808(13)	1263.11(3)	10.04	1.27

Table A2.12: Temperature dependent lattice parameters and agreement factors from powder XRD refinements of  $Al_4B_2O_9$ .

Table A2.13: Temperature dependent thermal expansion parameters  $\alpha$  of  ${\rm Al}_4{\rm B}_2{\rm O}_9.$  AF is the anisotropy factor.

T (K)	$\alpha_{a\text{-axis}} (\text{K}^{-1})$	$\alpha_{b\text{-axis}} \ (\mathrm{K}^{-1})$	$\alpha_{c\text{-axis}} \ (\mathrm{K}^{-1})$	$\alpha_{\mathit{Volume}}~(\mathrm{K}^{\text{-}1})$	AF
298	9.53666E-06	3.16331E-06	6.26697E-06	1.89688E-05	12.75
323	9.48925 E-06	3.20212E-06	6.32250 E-06	1.90170E-05	12.57
373	9.39455 E-06	3.27969E-06	6.43348E-06	1.91132E-05	12.23
423	9.29998 E-06	3.35724 E-06	6.54435 E-06	1.92091E-05	11.89
473	9.20554 E-06	3.43474E-06	6.65512 E-06	1.93048E-05	11.54
523	9.11123E-06	3.51220E-06	6.76577E-06	1.94001E-05	11.20
573	9.01704 E-06	3.58962 E-06	6.87631E-06	1.94951E-05	10.85
623	8.92298E-06	3.66700E-06	6.98674 E-06	1.95898E-05	10.51
673	8.82905 E-06	3.74433E-06	7.09705E-06	1.96843E-05	10.17
723	8.73523E-06	3.82163E-06	7.20724 E-06	1.97784 E-05	9.83
773	8.64154 E-06	3.89887E-06	7.31731E-06	1.98723E-05	9.49
823	8.54797E-06	3.97608E-06	7.42725E-06	1.99658E-05	9.14
873	8.45451E-06	4.05323E-06	7.53708E-06	2.00591 E-05	8.80
923	8.36118E-06	4.13034E-06	7.64678E-06	2.01520E-05	8.46
973	8.26795 E-06	4.20741E-06	7.75635E-06	2.02446E-05	8.12
1023	8.17485 E-06	4.28442 E-06	7.86580E-06	2.03370E-05	7.78
1073	8.08185 E-06	4.36138E-06	7.97511E-06	2.04290 E-05	7.44
1123	7.98897E-06	4.43829E-06	8.08429E-06	2.05208E-05	7.29
1173	7.89620E-06	4.51516E-06	8.19334E-06	2.06122 E-05	7.36
1223	7.80353E-06	4.59196E-06	8.30226E-06	2.07033E-05	7.42
1273	7.71097E-06	4.66872 E-06	8.41103E-06	2.07941 E-05	7.48
1323	7.61852E-06	4.74542E-06	8.51967E-06	2.08847E-05	7.55

### Supplementary data for grandidierite

**Table A2.14:** Details for single-crystal X-ray diffraction data collection and crystal structure refine-ment of grandidierite.

Diffractometer: Br	ruker Smart APEX 2, $MoK\alpha$ (0.71073	Å), 50 kV, 40 mA
Crystal size	$0.17\cdot 0.34$ $\cdot$	$0.27 \text{ mm}^3$
Detector distance	59.5	mm
Rotation axis	φ,	ω
Collection mode	Automated	hemisphere
Frame size	$512 \cdot 51$	2 pixel
Absorption corr.	Empirical ψ	correction
Resolution	No rest	riction
Sample name	Grandi	dierite
Temperature	298K	873K
Space group	Pnma	Pnma
a-axis	10.9967(1) Å	11.0193(1) Å
b-axis	5.7634(1) Å	5.7884(1) Å
c-axis	10.3321(1) Å	10.3434(1) Å
Volume	653.05(3)Å <sup>3</sup>	659.74(3)Å <sup>3</sup>
Z	4	4
Density $\rho$	$2.947~{ m g/cm^3}$	$2.917~{ m g/cm^3}$
Absorption coeff. $\mu$	$1.00 \text{ mm}^{-1}$	$0.99 \text{ mm}^{-1}$
Time per frame	10 s	20 s
Reflections collected	11250	15049
Reflections rejected	681	634
Reflections used	10569	14415
Max. 20	63.20°	89.18°
Index range $h$	-15 14	-16 21
Index range $k$	-8 8	-10 10
Index range $l$	-14 15	-15 18
Unique reflections	1105	2468
$\operatorname{Reflections} > 4 {f \sigma}(I)$	1118	2283
R(int)	0.0167	0.0162
$R(\mathbf{\sigma})$	0.0095	0.0110
Refined parameters	87	87
GoF	1.278	1.262
$R1,I>4{\bf \sigma}(I)$	0.0182	0.0207
R1, all data	0.0186	0.0230
wR2 on $F^2$	0.0563	0.0636
$\Delta  ho_{ m max}$	$0.46 e^{A^{-3}}$ , near O6	$0.53 e^{-1} A^{-3}$ , near O3
$\Delta  ho_{\min}$	-0.39 e <sup>-</sup> ·Å <sup>-3</sup> , near Si1	-0.37 e <sup>-</sup> ·Å <sup>-3</sup> , near Si1

Site	T (K)	x	y	z	$U_{eq.}$ (Å <sup>2</sup> )
Al1	298	0.0	0.0	0.0	0.00552(12)
	873	0.0	0.0	0.0	0.01061(5)
Al2	298	0.0	0.0	0.5	0.00582(12)
	873	0.0	0.0	0.5	0.01194(6)
Al3	298	0.44790(4)	0.25	0.22632(4)	0.00531(12)
	873	0.44836(2)	0.25	0.22639(2)	0.01087(5)
Si1	298	0.26328(4)	0.25	0.43360(4)	0.00556(11)
	873	0.26376(2)	0.25	0.43284(2)	0.01029(5)
Mg1	$298^{\mathrm{a}}$	0.21907(5)	0.25	0.09193(5)	0.00713(19)
	$873^{\rm b}$ b	0.21984(3)	0.25	0.09065(3)	0.01554(10)
01	298	0.28803(10)	0.25	0.27516(10)	0.0084(2)
	873	0.28891(6)	0.25	0.27480(6)	0.01763(12)
O2	298	0.52252(10)	0.25	0.38146(10)	0.0062(2)
	873	0.52276(6)	0.25	0.38136(6)	0.01139(9)
O3	298	0.00338(9)	0.25	-0.12097(11)	0.0072(2)
	873	0.00383(6)	0.25	-0.12137(7)	0.01455(11)
O4	298	0.11999(10)	0.25	0.47382(11)	0.0075(2)
	873	0.12076(5)	0.25	0.47120(8)	0.01589(11)
O5	298	0.10016(10)	-0.25	0.54648(10)	0.0061(2)
	873	0.09928(5)	-0.25	0.54548(6)	0.01147(9)
O6	298	0.32917(7)	0.02270(14)	0.49260(7)	0.00760(17)
	873	0.32896(4)	0.02350(8)	0.49222(5)	0.01606(9)
07	298	0.50112(7)	-0.04508(15)	0.18058(8)	0.00808(16)
	873	0.49932(5)	-0.04632(9)	0.18077(5)	0.01725(9)
B1	298	-0.00037(13)	0.25	-0.25127(16)	0.0067(3)
	873	0.00112(8)	0.25	-0.25123(9)	0.01210(14)

Table A2.15: Atomic positions and isotropic displacement parameters of grandidierite.

<sup>a</sup> occupancy: 0.973(3) Mg, 0.030(3) Fe; <sup>b</sup> occupancy: 0.9651(19) Mg, 0.0349(19) Fe

 Table A2.16:
 Anisotropic displacement parameters of grandidierite at 298K.

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.0059(2)	0.0048(2)	0.0058(2)	-0.00025(14)	0.00015(13)	-0.00004(14)
Al2	0.0074(2)	0.0052(2)	0.0049(2)	0.00019(14)	-0.00001(13)	0.00055(14)
Al3	0.0060(2)	0.0052(2)	0.0047(2)	0.0	-0.00027(14)	0.0
Si1	0.00512(18)	0.00545(18)	0.00609(19)	0.0	0.00028(13)	0.0
$Mg1^{a}$	0.0068(3)	0.0065(3)	0.0082(3)	0.0	-0.00134(17)	0.0
01	0.0069(5)	0.0121(5)	0.0064(5)	0.0	0.0008(4)	0.0
O2	0.0074(4)	0.0053(4)	0.0058(4)	0.0	-0.0004(4)	0.0
O3	0.0104(5)	0.0056(5)	0.0057(5)	0.0	-0.0001(3)	0.0
O4	0.0054(5)	0.0066(5)	0.0105(5)	0.0	0.0013(4)	0.0
O5	0.0064(4)	0.0056(4)	0.0064(4)	0.0	-0.0003(3)	0.0
O6	0.0064(4)	0.0065(3)	0.0099(4)	0.0018(3)	0.0000(2)	0.0004(3)
07	0.0128(4)	0.0058(3)	0.0057(3)	0.0003(3)	0.0002(2)	0.0009(2)
B1	0.0063(7)	0.0068(7)	0.0069(7)	0.0	0.0002(5)	0.0

<sup>a</sup>occupancy: 0.973(3) Mg, 0.030(3)

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.01149(10)	0.00909(11)	0.01127(12)	-0.00059(8)	0.00021(6)	-0.00010(6)
Al2	0.01653(11)	0.01002(12)	0.00926(11)	0.00067(8)	-0.00063(7)	0.00203(7)
Al3	0.01320(10)	0.01055(11)	0.00886(10)	0.0	-0.00076(7)	0.0
Si1	0.00890(8)	0.01062(10)	0.01134(10)	0.0	0.00026(6)	0.0
$Mg1^{a}$	0.01417(14)	0.01446(16)	0.01799(17)	0.0	-0.00357(9)	0.0
01	0.0138(2)	0.0280(4)	0.0111(2)	0.0	0.00057(18)	0.0
O2	0.0143(2)	0.0105(2)	0.0094(2)	0.0	-0.00135(16)	0.0
O3	0.0241(3)	0.0114(3)	0.0082(2)	0.0	-0.00007(17)	0.0
O4	0.00985(19)	0.0145(3)	0.0233(3)	0.0	0.00305(18)	0.0
O5	0.01085(18)	0.0110(2)	0.0126(2)	0.0	-0.00069(15)	0.0
O6	0.01218(15)	0.01259(18)	0.0234(2)	0.00506(15)	-0.00041(13)	0.00099(12)
07	0.0325(3)	0.00941(18)	0.00987(17)	0.00119(14)	0.00038(14)	0.00229(14)
B1	0.0180(3)	0.0096(3)	0.0087(3)	0.0	0.0003(2)	0.0

Table A2.17: Anisotropic displacement parameters of grandidierite at 873K.

<sup>a</sup>occupancy: 0.9651(19) Mg, 0.0349(19) Fe

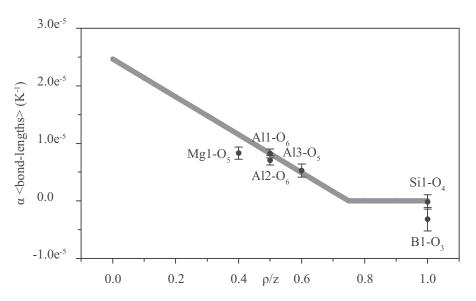


Figure A2.6: Expansion parameters  $\alpha$  calculated from averaged polyhedral bond length of grandidierite. Expected values according to Hazen and Prewitt (1977) plot on the gray line.

Center	Ligand	$d_{298K} (\text{\AA})$	d <sub>873K</sub> (Å)	$\alpha$ (K <sup>-1</sup> )	$ \mathrm{esd}  \ \alpha \ (\mathrm{K}^{\text{-}1})$
Al1	O2	1.9071(7)	1.9139(4)	6.201E-06	7.352E-07
	O2	1.9071(7)	1.9139(4)	6.201E-06	7.352E-07
	O3	1.9078(7)	1.9162(4)	7.657E-06	7.350E-07
	O3	1.9078(7)	1.9192(4)	7.657E-06	7.350E-07
	O6	1.8795(8)	1.8913(5)	1.092 E- 05	8.730E-07
	O6	1.8795(8)	1.8913(5)	1.092E-05	8.730E-07
	<al1-o></al1-o>	1.8981(7)	1.9076(4)	8.246E-06	7.804E-07
Al2	O4	1.9699(7)	1.9883(4)	1.624E-05	7.118E-07
	O4	1.9699(7)	1.9883(4)	1.624E-05	7.118E-07
	O5	1.8743(7)	1.8741(4)	-1.856E-07	7.481E-07
	O5	1.8743(7)	1.8741(4)	-1.856E-07	7.481E-07
	07	1.8838(8)	1.8889(5)	4.708 E-06	8.710E-07
	07	1.8839(9)	1.8889(5)	4.616E-06	9.504 E-07
	<al2-o></al2-o>	1.9094(8)	1.9171(4)	7.059E-06	7.890E-07
Al3	01	1.8244(12)	1.8270(7)	2.478E-06	1.324E-06
	O2	1.7998(12)	1.8004(7)	5.798E-07	1.342E-06
	O5	1.9314(11)	1.9435(7)	1.090E-05	1.174E-06
	07	1.8591(9)	1.8655(5)	$5.987 \text{E}{-}06$	9.631E-07
	07	1.8591(9)	1.8655(5)	$5.987 \text{E}{-}06$	9.631E-07
	<al3-o></al3-o>	1.8548(11)	1.8604(6)	5.270 E-06	1.151E-06
Mg1	01	$2.0386(12)^{\rm a}$	$2.0512(7)^{b}$	1.075E-05	1.185E-06
	O2	$2.1792(12)^{a}$	$2.1908(7)^{b}$	9.257E-06	1.109E-06
	O5	$2.0373(12)^{\rm a}$	$2.0472(7)^{\rm b}$	8.451E-06	1.186E-06
	O6	$1.9503(8)^{a}$	$1.9575(5)^{\rm b}$	6.420E-06	8.413E-07
	O6	$1.9503(8)^{a}$	$1.9575(5)^{\rm b}$	6.420E-06	8.413E-07
	<Mg1-O $>$	$2.0311(19)^{a}$	$2.0407(6)^{b}$	8.305E-06	1.037E-06
Si1	01	1.6593(12)	1.6580(7)	-1.363E-06	1.456E-06
	O4	1.6255(12)	1.6250(6)	-5.350E-07	1.435E-06
	O6	1.6155(8)	1.6162(5)	7.536E-07	1.016E-06
	O6	1.6158(8)	1.6162(5)	4.305E-07	1.015E-06
	<si1-o></si1-o>	1.6290(10)	1.6289(6)	-1.868E-07	1.231E-06
B1	O3	1.3469(19)	1.3436(11)	-4.261E-06	2.835E-06
	07	1.3750(11)	1.3729(7)	-2.656E-06	1.649E-06
	07	1.3750(11)	1.3729(7)	-2.656E-06	1.649E-06
	<b1-o></b1-o>	1.3656(14)	1.3631(8)	-3.184E-06	2.038E-06

 Table A2.18: Polyhedral bond lengths of grandidierite.

<sup>a</sup>occupancy: 0.973(3) Mg, 0.030(3) Fe; <sup>b</sup>occupancy: 0.9651(19) Mg, 0.0349(19) Fe

	0	0	
Sites	Angle (°)	Sites	Angle ( $^{\circ}$ )
Al3-O1-Mg1	95.72(5)	Al2-O5-Al2	100.48(5)
Si1-O1-Al3	115.47(6)	Al2-O5-Al3(x2)	94.97(4)
Si1-O1-Mg1	148.81(7)	Al2-O5-Mg1(x4)	129.00(3)
Al1-O2-Al1	98.14(5)	Al3-O5-Mg1(x2)	92.51(5)
Al1-O2-Mg1(x4)	92.71(4)	Mg1-O5-Mg1	0.0
Al3-O2-Al1(x2)	120.87(4)	Al1-O6-Mg1(x2)	101.12(4)
Al3-O2-Mg1(x2)	124.32(6)	Mg1-O6-Mg1	0.00(3)
Mg1-O2-Mg1	0.00(3)	Si1-O6-Al1	121.18(5)
Al1-O3-Al1	98.10(5)	Si1-O6-Mg1(x2)	136.94(5)
B1-O3-Al1(x2)	130.86(3)	A13-O7-A12	97.10(4)
Al2-O4-Al2	94.01(5)	B1-O7-Al2	128.72(8)
Si1-O4-Al2(x2)	132.91(2)	B1-O7-Al3	130.82(8)

Table A2.19: M-O-M angles of the refined structure of grandidierite at 298K.

 Table A2.20:
 M-O-M angles of the refined structure of grandidierite at 873K.

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
Al3-O1-Mg1	95.87(3)	Al2-O5-Al2	101.10(3)
Si1-O1-Al3	115.53(4)	Al2-O5-Al3(x2)	94.82(2)
Si1-O1-Mg1	148.60(4)	Al2-O5-Mg1(x4)	128.729(14)
Al1-O2-Al1	98.24(3)	A13-O5-Mg1(x2)	92.48(3)
Al1-O2-Mg1(x4)	92.59(2)	Mg1-O5-Mg1	0.000(19)
Al3-O2-Al1(x2)	120.74(2)	Al1-O6-Mg1(x2)	101.15(2)
Al3-O2-Mg1(x2)	124.69(3)	Mg1-O6-Mg1	0.000(19)
Mg1-O2-Mg1	0.000(18)	Si1-O6-Al1	121.16(3)
Al1-O3-Al1	98.09(3)	Si1-O6-Mg1(x2)	137.02(3)
B1-O3-Al1(x2)	130.880(16)	Al3-O7-Al2	96.95(2)
Al2-O4-Al2	93.40(3)	B1-O7-Al2	128.98(5)
Si1-O4-Al2(x2)	133.279(13)	B1-O7-Al3	131.22(5)

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
O2-Al1-O2	180.0	O1-Al3-O7(x2)	111.84(3)
O2-Al1-O3(x2)	81.23(4)	O2-Al3-O1	100.99(5)
O2-Al1-O3(x2)	98.77(4)	O2-Al3-O5	168.79(5)
O3-Al1-O3	180.0	O2-A13-O7(x2)	94.80(3)
O6-Al1-O2(x2)	92.89(4)	O7-Al3-O5(x2)	80.86(3)
O6-Al1-O2(x2)	87.11(4)	O7-Al3-O7	132.35(5)
O6-Al1-O3(x2)	89.63(4)	O3-B1-O7(x2)	120.80(7)
O6-Al1-O3(x2)	90.38(4)	O7-B1-O7	118.39(14)
O6-Al1-O6	180.0	O1-Mg1-O2	104.51(5)
O4-Al2-O4	180.00(5)	O5-Mg1-O1	81.55(5)
O5-Al2-O4(x2)	101.89(4)	O5-Mg1-O2	173.94(5)
O5-Al2-O4(x2)	78.11(4)	O6-Mg1-O1(x2)	126.11(3)
O5-Al2-O5	180.00(1)	O6-Mg1-O2(x2)	78.31(3)
O5-Al2-O7(x2)	98.28(4)	O6-Mg1-O5(x2)	98.21(3)
O5-Al2-O7(x2)	81.72(4)	O6-Mg1-O6	107.39(5)
O7-Al2-O4(x2)	87.74(4)	O4-Si1-O1	114.23(6)
O7-Al2-O4(x2)	92.26(4)	O6-Si1-O1(x2)	107.40(4)
O7-Al2-O7	180.00(5)	O6-Si1-O4(x2)	109.63(4)
O1-Al3-O5	90.22(5)	O6-Si1-O6	108.36(6)

Table A2.21: O-M-O angles of the refined structure of grandidierite at 298K.

Table A2.22: O-M-O angles of the refined structure of grandidierite at 873K.

Sites	Angle (°)	Sites	Angle ( $^{\circ}$ )
O2-Al1-O2	180.0	O1-Al3-O7(x2)	111.028(18)
O2-Al1-O3(x2)	81.15(2)	O2-Al3-O1	101.18(3)
O2-Al1-O3(x2)	98.85(2)	O2-Al3-O5	168.58(3)
O3-Al1-O3	180.0	O2-Al3-O7(x2)	95.05(2)
O6-Al1-O2(x2)	92.81(2)	O7-Al3-O5(x2)	80.663(19)
O6-Al1-O2(x2)	87.19(2)	O7-Al3-O7	133.69(4)
O6-Al1-O3(x2)	89.74(3)	O3-B1-O7(x2)	120.82(4)
O6-Al1-O3(x2)	90.26(3)	O7-B1-O7	118.36(8)
O6-Al1-O6	180.0	O1-Mg1-O2	104.18(3)
O4-Al2-O4	180.0	O5-Mg1-O1	81.41(3)
O5-A12-O4(x2)	102.06(2)	O5-Mg1-O2	174.40(3)
O5-Al2-O4(x2)	77.94(2)	O6-Mg1-O1(x2)	125.795(17)
O5-Al2-O5	180.0	O6-Mg1-O2(x2)	78.257(19)
O5-Al2-O7(x2)	98.12(2)	O6-Mg1-O5(x2)	98.55(2)
O5-Al2-O7(x2)	81.88(2)	O6-Mg1-O6	107.95(3)
O7-Al2-O4(x2)	87.57(3)	O4-Si1-O1	113.75(4)
O7-A12-O4(x2)	92.43(3)	O6-Si1-O1(x2)	107.48(2)
O7-A12-O7	180.0	O6-Si1-O4(x2)	109.77(2)
O1-Al3-O5	90.24(3)	O6-Si1-O6	108.43(4)

grandiaterite.				
T (K)	a-axis (Å)	b-axis (Å)	c-axis (Å)	Volume (Å)
298	10.9753(6)	5.7650(4)	10.3327(6)	653.78(5)
323	10.9772(6)	5.7659(3)	10.3333(6)	654.03(5)
373	10.9820(6)	5.7680(4)	10.3334(6)	654.57(5)
423	10.9854(6)	5.7697(3)	10.3341(6)	655.00(5)
473	10.9892(6)	5.7717(3)	10.3348(5)	655.50(5)
523	10.9942(7)	5.7735(4)	10.3370(6)	656.14(5)
573	10.9998(6)	5.7760(3)	10.3375(6)	656.79(5)
623	11.0029(6)	5.7780(3)	10.3387(6)	657.28(5)
673	11.0072(6)	5.7805(3)	10.3405(5)	657.94(5)
723	11.0128(7)	5.7834(4)	10.3424(6)	658.72(6)
773	11.0163(7)	5.7863(4)	10.3437(6)	659.34(6)
823	11.0225(6)	5.7874(3)	10.3455(6)	659.96(6)
873	11.0283(6)	5.7901(4)	10.3467(6)	660.69(6)
923	11.0329(7)	5.7941(4)	10.3488(6)	661.55(6)
973	11.0386(7)	5.7958(4)	10.3493(6)	662.12(6)

 Table A2.23:
 Temperature dependent lattice parameters from single-crystal unit cell refinements of grandidierite.

Table A2.24: Temperature dependent thermal expansion parameters  $\alpha$  of grandidierite. AF is the anisotropy factor.

	1400011				
T (K)	$\alpha_{a\text{-axis}}$ (K <sup>-1</sup> )	$\alpha_{b\text{-axis}}$ (K <sup>-1</sup> )	$\alpha_{c\text{-axis}}$ (K <sup>-1</sup> )	$\alpha_{Volume} ~(\mathrm{K}^{-1})$	AF
298	1.40815E-06	1.33177E-05	3.46936E-06	1.46746E-05	23.82
323	1.48210E-06	1.35266E-05	3.54291 E-06	1.49979E-05	24.09
373	1.62989E-06	1.39442 E-05	3.68996E-06	1.56432 E-05	24.63
423	1.77756E-06	1.43612E-05	3.83696E-06	1.62870E-05	25.17
473	1.92510E-06	1.47778E-05	3.98389E-06	1.69293E-05	25.71
523	2.07251 E-06	1.51939E-05	4.13077E-06	1.75699E-05	26.24
573	2.21977E-06	1.56094 E-05	4.27759E-06	1.82088E-05	26.78
623	2.36689E-06	1.60245 E-05	4.42433E-06	1.88459E-05	27.32
673	2.51385 E-06	1.64390 E-05	4.57100 E-06	1.94812 E-05	27.85
723	2.66067E-06	1.68530E-05	4.71760E-06	2.01147 E-05	28.38
773	2.80732E-06	1.72664 E-05	4.86412E-06	2.07462 E-05	28.92
823	2.95381E-06	1.76792E-05	5.01055E-06	2.13757E-05	29.45
873	3.10013E-06	1.80914 E-05	5.15691E-06	2.20032 E-05	29.98
923	3.24628E-06	1.85029E-05	5.30317 E-06	2.26287 E-05	30.51
973	3.39225 E-06	1.89139E-05	5.44934E-06	2.32519E-05	31.04

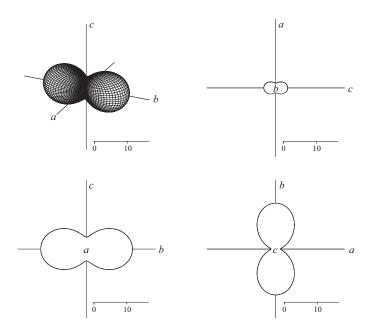


Figure A2.7: Representation surface of thermal expansion tensor components of grandidierite at 298K. Top left: Arbitrary view, top right: Cross-section in (010), bottom left: Cross-section in (100) and bottom right: Cross section in (001). a, b and c are directions parallel to unit cell axes, scale bar unit is  $10^{-6}$ ·K<sup>-1</sup>.

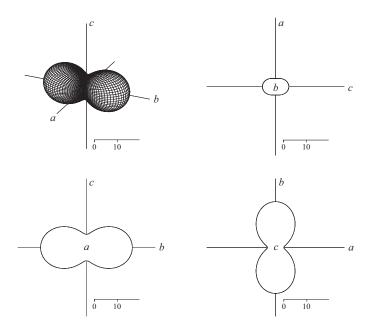


Figure A2.8: Representation surface of thermal expansion tensor components of grandidierite at 973K. Top left: Arbitrary view, top right: Cross-section in (010), bottom left: Cross-section in (100) and bottom right: Cross section in (001). a, b and c are directions parallel to unit cell axes, scale bar unit is  $10^{-6}$ ·K<sup>-1</sup>.

Plots were made with WinTensor v. 1.2 by Werner Kaminsky, Dept. of Chemistry, University of Washington.

### Supplementary data for jeremejevite

 Table A2.25: Details for single-crystal X-ray diffraction data collection and crystal structure refinement of jeremejevite.

Diffractometer: E	Bruker Smart APEX 2, $MoK\alpha$ (0.71073 A	Å), 50 kV, 40 mA
Crystal size	$0.14 \cdot 0.14$ ·	$0.07 \text{ mm}^3$
Detector distance	49.5 1	nm
Rotation axis	φ, α	υ
Collection mode	Automated h	nemisphere
Frame size	$512 \cdot 512$	2 pixel
Absorption corr.	Empirical $\psi$	correction
Resolution	No restr	iction
Sample name	Jereme	jevite
Temperature	298K	873K
Space group	$P6_3/m$	$P6_3/m$
a-axis	8.5526(3) Å	8.5840(16) Å
c-axis	8.1793(4) Å	8.2187(15) Å
Volume	518.13(4) Å <sup>3</sup>	524.46(17) Å <sup>3</sup>
Ζ	2	2
Density p	$3.288~{ m g/cm^3}$	$3.288~{ m g/cm^3}$
Absorption coeff. µ	$0.79 \text{ mm}^{-1}$	$0.78 \text{ mm}^{-1}$
Time per frame	40 s	40 s
Reflections collected	5498	4806
Reflections rejected	7	2284
Reflections used	5491	2522
Max. 20	$69.21^{\circ}$	$60.86^{\circ}$
Index range $h$	-12 11	-16 21
Index range $k$	-11 13	-11 11
Index range $l$	-12 8	-11 11
Unique reflections	781	555
$ ext{Reflections} > 4 \mathbf{\sigma}(I)$	673	465
R(int)	0.0324	0.0270
$R(\sigma)$	0.0222	0.0251
Refined parameters	50	50
GoF	1.064	1.236
$R1,I>4{\bf \sigma}(I)$	0.0258	0.0308
R1, all data	0.0326	0.0434
w $R2$ on $F^2$	0.0627	0.1145
$\Delta \rho_{\rm max}$	$0.49 e^{-}$ Å <sup>-3</sup> , near O1	$0.69 e^{-} A^{-3}$ , near O3
$\Delta \rho_{\min}$	$-0.40 e^{-} \cdot A^{-3}$ , near All	$-0.73 \text{ e}^{-3}$ , near Ol

Site	T (K)	x	y	z	$U_{eq.}$ (Å <sup>2</sup> )
Al1	298	0.33719(4)	0.35252(4)	0.57536(4)	0.00522(11)
	873	0.33696(10)	0.35223(11)	0.57509(9)	0.0120(3)
F1	298	0.38689(11)	0.49995(11)	0.75	0.00745(18)
	873	0.3868(3)	0.4994(3)	0.75	0.0178(5)
B1	298	0.2335(2)	0.0023(2)	0.75	0.0060(3)
	873	0.2320(6)	0.0031(6)	0.75	0.0121(7)
B2	298	0.6667	0.3333	0.5545(2)	0.0058(3)
	873	0.6667	0.3333	0.5539(6)	0.0120(9)
01	298	0.30281(13)	0.18634(13)	0.75	0.0063(2)
	873	0.3020(4)	0.1854(4)	0.75	0.0136(5)
O2	298	0.58909(10)	0.43927(10)	0.55453(9)	0.00632(16)
	873	0.5890(3)	0.4383(3)	0.5538(2)	0.0141(4)
O3	298	0.08819(9)	0.28503(10)	0.60141(9)	0.00630(16)
	873	0.0874(3)	0.2845(3)	0.6018(2)	0.0138(4)

Table A2.26: Atomic positions and isotropic displacement parameters of jeremejevite.

 Table A2.27: Anisotropic displacement parameters of jeremejevite at 298K.

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.00537(15)	0.00532(15)	0.00505(17)	-0.00002(10)	0.00006(10)	0.00273(11)
F1	0.0096(4)	0.0066(4)	0.0054(4)	0.0	0.0	0.0035(3)
B1	0.0054(6)	0.0061(6)	0.0065(6)	0.0	0.0	0.0029(5)
B2	0.0063(4)	0.0063(4)	0.0048(8)	0.0	0.0	0.0031(2)
01	0.0077(4)	0.0057(4)	0.0059(4)	0.0	0.0	0.0037(3)
O2	0.0058(3)	0.0057(3)	0.0074(3)	0.0006(2)	0.0004(3)	0.0029(2)
O3	0.0057(3)	0.0084(3)	0.0051(3)	0.0006(2)	0.0006(2)	0.0038(3)

Table A2.28: Anisotropic displacement parameters of jeremejevite at 873K.

Site	$U_{11}$ (Å <sup>2</sup> )	$U_{22}$ (Å <sup>2</sup> )	$U_{33}$ (Å <sup>2</sup> )	$U_{23}$ (Å <sup>2</sup> )	$U_{13}$ (Å <sup>2</sup> )	$U_{12}$ (Å <sup>2</sup> )
Al1	0.0124(4)	0.0128(4)	0.0110(4)	0.0003(3)	0.0004(3)	0.0065(3)
F1	0.0241(13)	0.0157(11)	0.0132(10)	0.0	0.0	0.0097(10)
B1	0.0133(18)	0.0130(18)	0.0106(15)	0.0	0.0	0.0071(15)
B2	0.0118(13)	0.0118(13)	0.0124(19)	0.0	0.0	0.0059(7)
01	0.0185(14)	0.0113(12)	0.0131(11)	0.0	0.0	0.0089(11)
O2	0.0134(9)	0.0132(9)	0.0171(8)	0.0023(7)	0.0010(7)	0.0076(7)
O3	0.0121(9)	0.0185(10)	0.0106(7)	0.0017(7)	0.0014(6)	0.0076(8)

Center	Ligand	$d_{298K}$ (Å)	d <sub>873K</sub> (Å)	α (Κ <sup>-1</sup> )	$ \mathrm{esd}  \; \alpha \; (\mathrm{K}^{\text{-}1})$
Al1	01	1.9310(7)	1.9438(20)	1.153E-05	1.908E-06
	O2	1.8907(8)	1.9021(20)	8.955E-06	2.053E-06
	O2	1.9033(8)	1.9131(21)	1.049E-05	1.981E-06
	O3	1.9196(8)	1.9312(21)	1.051E-05	2.036E-06
	O3-6	1.8623(8)	1.8683(20)	5.603E-06	2.012 E-06
	F1	1.8097(6)	1.8182(20)	8.169E-06	1.849E-06
	<al1-o></al1-o>	1.8861(8)	1.8961(16)	9.236E-06	1.972E-06
B1	01	1.3767(18)	1.3673(48)	-1.187E-05	6.476E-06
	O3	1.3899(10)	1.3940(26)	5.130E-06	3.486E-06
	O3	1.3899(10)	1.3940(26)	5.130E-06	3.486E-06
	<b1-0></b1-0>	1.3855(13)	1.3851(40)	-5.021E-07	4.476E-06
B2	O2	1.3646(7)	1.3622(19)	-2.931E-06	2.581E-06
	O2	1.3646(7)	1.3623(19)	-2.931E-06	2.581E-06
	O2	1.3646(7)	1.3623(19)	-2.931E-06	2.581E-06
	<b2-o></b2-o>	1.3646(7)	1.3623(19)	-2.931E-06	2.581E-06

 Table A2.29:
 Polyhedral bond lengths of jeremejevite.

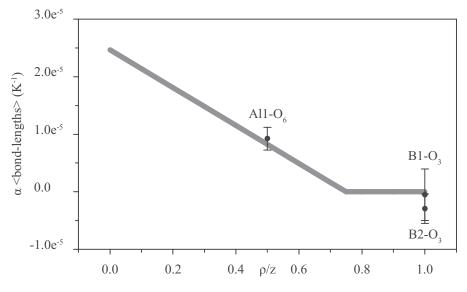


Figure A2.9: Expansion parameters  $\alpha$  calculated from averaged polyhedral bond length of jeremejevite. Expected values according to Hazen and Prewitt (1977) plot on the gray line.

Sites	Angle (°)	Sites	Angle ( $^{\circ}$ )
Al1-F1-Al1	104.24(5)	B2-O2-Al1	124.93(5)
Al1-O1-Al1	95.41(5)	Al1-O3-Al1	115.51(4)
B1-O1-Al1 $(x2)$	131.41(3)	B1-O3-Al1	120.38(7)
Al1-O2-Al1	101.35(3)	B1-O3-Al1	123.16(7)
B2-O2-Al1	125.55(8)		

Table A2.30: M-O-M angles of the refined structure of jeremejevite at 298K.

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-F1-Al1	104.49(13)	B2-O2-Al1	125.27(13)
Al1-O1-Al1	95.39(12)	Al1-O3-Al1	115.14(9)
B1-O1-Al1 (x2)	131.37(8)	B1-O3-Al1	120.9(2)
Al1-O2-Al1	101.21(9)	B1-O3-Al1	123.14(19)
B2-O2-Al1	125.65(19)		

Table A2.31: M-O-M angles of the refined structure of jeremejevite at 873K.

Table A2.32: O-M-O angles of the refined structure of jeremejevite at 298K.

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
F1-Al1-O1	79.94(3)	O2-All-O3	175.23(4)
F1-Al1-O2	86.32(3)	O3-Al1-O1	98.87(3)
F1-Al1-O2	88.47(4)	O3-Al1-O1	87.64(4)
F1-Al1-O3	178.81(4)	O3-Al1-O2	94.87(3)
F1-Al1-O3	87.74(4)	O3-Al1-O2	91.66(3)
O2-Al1-O1	164.85(4)	O3-Al1-O3	92.18(4)
O2-Al1-O1	94.55(4)	O1-B1-O3 (x2)	119.01(6)
O2-Al1-O2	78.65(4)	O3-B1-O3	121.95(12)
O2-Al1-O3	98.22(3)	O2-B2-O2 (x3)	120.0

 Table A2.33:
 O-M-O angles of the refined structure of jeremejevite at 873K.

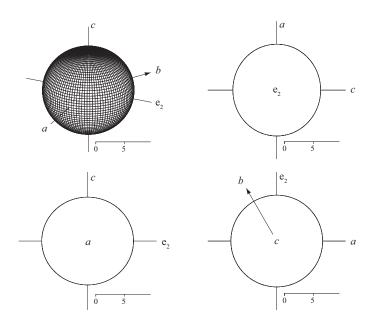
Sites	Angle (°)	Sites	Angle (°)
F1-Al1-O1	79.83(9)	O2-Al1-O3	175.42(10)
F1-Al1-O2	86.10(9)	O3-Al1-O1	99.03(9)
F1-Al1-O2	88.64(10)	O3-Al1-O1	87.40(10)
F1-Al1-O3	178.86(10)	O3-Al1-O2	95.03(9)
F1-Al1-O3	87.66(10)	O3-Al1-O2	91.55(9)
O2-Al1-O1	164.63(10)	O3-Al1-O3	92.21(11)
O2-Al1-O1	94.60(10)	O1-B1-O3 (x2)	119.13(17)
O2-Al1-O2	78.79(9)	O3-B1-O3	121.7(3)
O2-Al1-O3	98.27(9)	O2-B2-O2 (x3)	120.002(1)

	of jerennejevne.				
T (K)	a-axis (Å)	c-axis (Å)	Volume (Å)	$R_{wp}$	GoF
298	8.55670(4)	8.18027(7)	518.694(6)	13.13	1.16
323	8.55835(4)	8.18181(7)	518.991(6)	12.93	1.13
373	8.56164(4)	8.18489(6)	519.587(6)	13.05	1.14
423	8.56494(4)	8.18797(6)	520.183(6)	13.23	1.15
473	8.56825(3)	8.19105(6)	520.780(5)	13.03	1.13
523	8.57156(3)	8.19412(5)	521.377(5)	13.31	1.15
573	8.57487(3)	8.19719(5)	521.976(5)	13.36	1.15
623	8.57819(3)	8.20025(5)	522.576(5)	13.30	1.14
673	8.58152(3)	8.20331(5)	523.177(5)	13.46	1.15
723	8.58485(3)	8.20637(5)	523.778(5)	13.40	1.15
773	8.58819(3)	8.20943(5)	524.381(5)	13.27	1.13
823	8.59153(4)	8.21248(5)	524.984(5)	13.16	1.12
873	8.59488(4)	8.21553(6)	525.589(6)	13.35	1.14
923	8.59823(4)	8.21858(7)	526.194(6)	13.29	1.13
973	8.60159(5)	8.22162(8)	526.800(7)	13.38	1.14
1023	8.60495(5)	8.22467(9)	527.407(8)	13.74	1.17
1073	8.60832(6)	8.22770(10)	528.015(9)	13.55	1.15

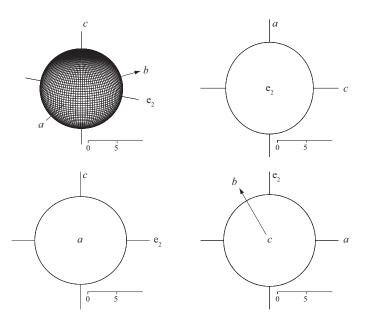
**Table A2.34:** Temperature dependent lattice parameters and agreement factors from powder XRD refinements of jeremejevite.

Table A2.35: Temperature dependent thermal expansion parameters  $\alpha$  of jeremejevite. AF is the anisotropy factor.

T (K)	$\alpha_{a-axis}$ (K <sup>-1</sup> )	$\alpha_{c-axis}$ (K <sup>-1</sup> )	$\alpha_{Volume} ~(\mathrm{K}^{-1})$	AF
298	7.68811E-06	7.54012E-06	2.29131E-05	0.30
323	7.69281E-06	7.53492E-06	2.29177 E-05	0.32
373	7.70220E-06	7.52454E-06	2.29267 E-05	0.36
423	7.71158E-06	7.51418E-06	2.29357 E-05	0.39
473	7.72095E-06	7.50382 E-06	2.29447 E-05	0.43
523	7.73030E-06	7.49347 E-06	2.29535 E-05	0.47
573	7.73965E-06	7.48313E-06	2.29623E-05	0.51
623	7.74898E-06	7.47281E-06	2.29711E-05	0.55
673	7.75830E-06	7.46249 E-06	2.29798 E-05	0.59
723	7.76761E-06	7.45219E-06	2.29884 E-05	0.63
773	7.77690E-06	7.44189E-06	2.29970 E-05	0.67
823	7.78619E-06	7.43161E-06	2.30055 E-05	0.71
873	7.79546E-06	7.42134E-06	2.30139 E-05	0.75
923	7.80472E-06	7.41107 E-06	2.30223E-05	0.79
973	7.81397E-06	7.40082 E-06	2.30306 E-05	0.83
1023	7.82321E-06	7.39058E-06	2.30389 E-05	0.87
1073	7.83243E-06	7.38034E-06	2.30471 E-05	0.90



**Figure A2.10:** Representation surface of thermal expansion tensor components of jeremejevite at 298K. Eigenvector  $e_1$  is parallel a,  $e_3$  is parallel c, the angle between the a-axis and the b-axis is 120°. Top left: Arbitrary view, top right: Cross-section of the a-c-plane, bottom left: Cross-section of the  $e_2$ -c-plane and bottom right: Cross section of the  $e_2$ -a-plane. a, b and c are directions parallel to unit cell axes, scale bar unit is  $10^{-6} \cdot K^{-1}$ .



**Figure A2.11:** Representation surface of thermal expansion tensor components of jeremejevite at 1073K. Eigenvector  $e_1$  is parallel *a*,  $e_3$  is parallel *c*, the angle between the *a*-axis and the *b*-axis is 120°. Top left: Arbitrary view, top right: Cross-section of the *a*-*c*-plane, bottom left: Cross-section of the  $e_2$ -*c*-plane and bottom right: Cross section of the  $e_2$ -*a*-plane. *a*, *b* and *c* are directions parallel to unit cell axes, scale bar unit is  $10^{-6} \cdot K^{-1}$ .

Plots were made with WinTensor v. 1.2 by Werner Kaminsky, Dept. of Chemistry, University of Washington.

## A.3 Supplementary Data for Chapter 7

	Table A3	.1: Bond distances a	t $P0$ (0.001 GF	Pa), crystal in air.	
Center	Ligand	Distance $(\text{\AA})$	Center	Ligand	Distance $(\text{\AA})$
Al1	01	1.910(3)	Al2	01	1.801(4)
Al1	O2	1.902(3)	Al2	O2 (x2)	1.753(2)
Al1	O3	1.846(2)	Al2	O5	1.851(4)
Al1	O4	1.910(2)	Al2	O6	2.171(4)
Al1	O6	1.887(2)		<al2-o></al2-o>	1.866(3)
Al1	07	1.962(3)			
	<al1-o></al1-o>	1.903(3)	Al3	01	2.267(4)
			Al3	O3	1.752(4)
Al4	O2 (x2)	1.736(3)	Al3	O5	1.770(4)
Al4	O5	1.777(4)	A13	O7 (x2)	1.825(3)
Al4	O6	1.762(3)		<al3-o></al3-o>	1.888(4)
	<al4-o></al4-o>	1.753(3)			
B1	O4	1.372(7)			
B1	O7 (x2)	1.382(4)			
	<al5-o></al5-o>	1.379(5)			

Table A3.2: Bond distances at P0 (0.001 GPa), crystal in DAC.

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	01	1.937(13)	Al2	O1	1.734(18)
Al1	O2	1.912(15)	Al2	O2 (x2)	1.778(14)
Al1	O3	1.841(15)	Al2	O5	1.76(3)
Al1	O4	1.928(14)	Al2	O6	2.194(15)
Al1	O6	1.889(17)		<al2-o></al2-o>	1.849(18)
Al1	07	1.909(12)			
	<al1-o></al1-o>	1.903(14)	Al3	01	2.24(2)
			Al3	O3	1.74(2)
Al4	O2 (x2)	1.738(6)	Al3	O5	1.82(3)
Al4	O5	1.753(16)	Al3	O7 (x2)	1.824(5)
Al4	O6	1.77(3)		<al3-o></al3-o>	1.890(14)
	<al4-o></al4-o>	1.750(15)			
B1	O4	1.38(3)			
B1	O7 (x2)	1.393(13)			
	<b1-o></b1-o>	1.389(19)			

	Table A3.3. Dolid distances at 1 1 (0.15(3) G1 a).						
Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)		
Al1	01	1.929(10)	Al2	O1	1.763(14)		
Al1	O2	1.875(11)	Al2	O2 (x2)	1.769(11)		
Al1	O3	1.846(11)	Al2	O5	1.815(19)		
Al1	O4	1.918(11)	Al2	O6	2.180(13)		
Al1	O6	1.886(13)		<al2-o></al2-o>	1.860(14)		
Al1	07	1.951(10)					
	<al1-o></al1-o>	1.900(11)	Al3	01	2.303(17)		
			Al3	O3	1.730(18)		
Al4	O2 (x2)	1.735(6)	Al3	O5	1.841(19)		
Al4	O5	1.740(12)	Al3	O7 $(x2)$	1.834(5)		
Al4	O6	1.76(2)		<al3-o></al3-o>	1.908(14)		
	<al4-o></al4-o>	1.743(11)					
B1	O4	1.38(2)					
B1	O7 $(x2)$	1.384(11)					
	<b1-o></b1-o>	1.383(14)					

Table A3.3: Bond distances at P1 (0.15(5) GPa).

Table A3.4: Bond distances at P2 (0.91(4) GPa).

Center	Ligand	Distance $(Å)$	Center	Ligand	Distance (Å)
Al1	01	1.920(9)	Al2	01	1.745(13)
Al1	O2	1.927(9)	Al2	O2 (x2)	1.740(8)
Al1	O3	1.86(1)		Al2O5	1.823(15)
Al1	O4	1.911(9)	Al2	O6	2.167(11)
Al1	O6	1.897(10)		<al2-o></al2-o>	1.843(11)
Al1	07	1.922(9)			
	<al1-o></al1-o>	1.906(9)	Al3	01	2.302(15)
			A13	O3	1.706(16)
Al4	O2 (x2)	1.727(6)	Al3	O5	1.788(16)
Al4	O5	1.791(12)	A13	O7 $(x2)$	1.835(5)
Al4	O6	1.712(16)		<al3-o></al3-o>	1.893(11)
	<al4-o></al4-o>	1.739(10)			
B1	O4	1.39(2)			
B1	O7 (x2)	1.379(10)			
	<b1-o></b1-o>	1.383(14)			

Table A3.3: Dolid distances at F3 (1.99(0) GFa).						
Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)	
Al1	01	1.921(9)	Al2	O1	1.756(12)	
Al1	O2	1.90(1)	Al2	O2 (x2)	1.754(8)	
Al1	O3	1.858(9)	Al2	O5	1.807(14)	
Al1	O4	1.900(9)	Al2	O6	2.185(11)	
Al1	O6	1.873(10)		<al2-o></al2-o>	1.851(11)	
Al1	07	1.942(10)				
	<al1-o></al1-o>	1.899(10)	Al3	01	2.244(14)	
			Al3	O3	1.698(14)	
Al4	O2 (x2)	1.728(6)	Al3	O5	1.797(15)	
Al4	O5	1.753(11)	Al3	O7 $(x2)$	1.813(6)	
Al4	O6	1.775(16)		<al3-o></al3-o>	1.873(11)	
	<al4-o></al4-o>	1.746(10)				
B1	O4	1.38(2)				
B1	O7 (x2)	1.383(10)				
	<b1-o></b1-o>	1.38(17)				

Table A3.5: Bond distances at P3 (1.99(6) GPa).

Table A3.6: Bond distances at P4 (3.32(6) GPa).

Table A5.0. Dond distances at 14 (5.52(0) Gi a).						
Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)	
Al1	01	1.914(9)	Al2	O1	1.735(14)	
Al1	O2	1.880(9)	Al2	O2 (x2)	1.767(9)	
Al1	O3	1.849(10)	Al2	O5	1.838(15)	
Al1	O4	1.894(9)	Al2	O6	2.204(13)	
Al1	O6	1.867(11)		<al2-o></al2-o>	1.862(12)	
Al1	07	1.936(9)				
	<al1-o< td=""><td>1.890(10)</td><td>Al3</td><td>01</td><td>2.249(15)</td></al1-o<>	1.890(10)	Al3	01	2.249(15)	
			Al3	O3	1.720(15)	
Al4	O2 (x2)	1.715(7)	Al3	O5	1.803(15)	
Al4	O5	1.746(12)	Al3	O7 $(x2)$	1.816(6)	
Al4	O6	1.744(17)		<al3-o></al3-o>	1.881(12)	
	<al4-o></al4-o>	1.730(11)				
B1	O4	1.38(2)				
B1	O7 $(x2)$	1.387(11)				
	<b1-o></b1-o>	1.38(14)				

	Table A3.7. Dolid distances at $T_0$ (4.76(0) GI a).						
Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)		
Al1	O1	1.888(9)	Al2	O1	1.778(14)		
Al1	O2	1.863(10)	Al2	O2 (x2)	1.733(8)		
Al1	O3	1.845(10)	Al2	O5	1.862(16)		
Al1	O4	1.912(10)	Al2	O6	2.159(13)		
Al1	O6	1.860(11)		<al2-o></al2-o>	1.853(12)		
Al1	07	1.941(10)					
	<al1-o></al1-o>	1.885(10)	Al3	01	2.241(15)		
			Al3	O3	1.712(15)		
Al4	O2 (x2)	1.733(7)	Al3	O5	1.779(16)		
Al4	O5	1.737(13)	Al3	O7 $(x2)$	1.801(6)		
Al4	O6	1.749(17)		<al3-o></al3-o>	1.867(12)		
	<al4-o></al4-o>	1.738(11)					
B1	O4	1.37(2)					
B1	O7 $(x2)$	1.372(11)					
	<b1-o></b1-o>	1.37(12)					

Table A3.7: Bond distances at P5 (4.78(6) GPa).

Table A3.8: Bond distances at P6 (5.83(5) GPa).

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	01	1.888(9)	Al2	O1	1.778(14)
Al1	O2	1.857(10)	Al2	O2 (x2)	1.730(9)
Al1	O3	1.853(10)	Al2	O5	1.840(17)
Al1	O4	1.896(9)	Al2	O6	2.164(13)
Al1	O6	1.858(11)		<al2-o></al2-o>	1.848(12)
Al1	07	1.935(10)			
	Al1-O>	1.881(10)	Al3	O1	2.188(15)
			Al3	O3	1.696(15)
Al4	O2 (x2)	1.732(7)	Al3	O5	1.743(17)
Al4	O5	1.773(13)	Al3	O7 (x2)	1.801(7)
Al4	O6	1.753(17)		<al3-o></al3-o>	1.846(12)
	<al4-o></al4-o>	1.748(11)			
B1	O4	1.36(2)			
B1	O7 $(x2)$	1.379(11)			
	<b1-o></b1-o>	1.373(14)			

Table A3.9: bold distances at <i>F1</i> (5.99(5) GFa).					
Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	O1	1.910(11)	Al2	O1	1.728(16)
Al1	O2	1.866(11)	Al2	O2 (x2)	1.757(10)
Al1	O3	1.842(12)	Al2	O5	1.834(18)
Al1	O4	1.894(11)	Al2	O6	2.178(14)
Al1	O6	1.867(12)		<al2-o></al2-o>	1.851(14)
Al1	07	1.917(11)			
	<al1-o></al1-o>	1.883(11)	Al3	01	2.263(18)
			Al3	O3	1.705(18)
Al4	O2 (x2)	1.707(8)	Al3	O5	1.826(18)
Al4	O5	1.717(14)	Al3	O7 $(x2)$	1.808(7)
Al4	O6	1.743(19)		<al3-o></al3-o>	1.882(14)
	<al4-o></al4-o>	1.719(12)			
B1	O4	1.38(3)			
B1	O7 $(x2)$	1.379(12)			
	<b1-o></b1-o>	1.379(18)			

Table A3.9: Bond distances at P7 (5.99(5) GPa).

Table A3.10: Bond distances at P8 (6.45(6) GPa).

Center	Ligand	Distance (Å)	Center	Ligand	Distance (Å)
Al1	01	1.945(12)	Al2	01	1.699(17)
Al1	O2	1.871(12)	Al2	O2 (x2)	1.743(11)
Al1	O3	1.829(13)	Al2	O5	1.84(2)
Al1	O4	1.855(12)	Al2	O6	2.160(15)
Al1	O6	1.896(14)		<al2-o></al2-o>	1.837(15)
Al1	07	1.904(12)			
	<al1-o></al1-o>	1.883(13)	Al3	01	2.263(19)
			Al3	O3	1.703(19)
Al4	O2 (x2)	1.722(8)	Al3	O5	1.79(2)
Al4	O5	1.740(16)	Al3	O7 $(x2)$	1.804(7)
Al4	O6	1.70(2)		<al3-o></al3-o>	1.873(15)
	<al4.o></al4.o>	1.721(13)			
B1	O4	1.38(3)			
B1	O7 (x2)	1.395(13)			
	<b1-o></b1-o>	1.390(19)			

Sites	Angle $(^{\circ})$	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	95.05(5)	Al3-O5-Al2	103.13(16)
Al1-O1-Al3 (x2)	90.16(3)	Al3-O5-Al4	156.69(23)
Al2-O1-Al1 (x2)	132.44(4)	Al4-O5-Al2	100.18(20)
Al2-O1-Al3	87.69(13)	Al1-O6-Al2 (x2)	91.12(4)
Al2-O2-Al1	105.17(12)	Al1-O6-Al1	98.56(5)
Al4-O2-Al1	123.75(13)	Al4-O6-Al1 (x2)	130.71(3)
Al4-O2-Al2	129.56(14)	Al4-O6-Al2	89.45(17)
Al1-O3-Al1	99.49(5)	Al3-O7-Al1	103.15(12)
Al3-O3-Al1 (x2)	130.23(3)	B1-O7-Al1	125.85(16)
Al1-O4-Al1	97.01(5)	B1-O7-Al3	130.26(19)
B1-O4-Al1 (x2)	130.27(3)		

Table A3.11: M-O-M angles at P0 (0.001 GPa), crystal in air.

Table A3.12: O-M-O angles at P0 (0.001 GPa), crystal in air.

Sites	Angle (°)	Sites	Angle (°)
O1-Al1-O4	169.00(7)	O2-Al2-O5 (x2)	118.53(7)
O1-Al1-O7	82.78(10)	O2-Al2-O6 (x2)	79.10(8)
O2-Al1-O1	99.34(10)	O2-Al2-O2	111.86(10)
O2-Al1-O4	91.66(9)	O5-A12-O6	78.21(16)
O2-Al1-O7	172.76(12)	O3-Al3-O1	169.21(16)
O3-Al1-O1	80.73(4)	O3-A13-O5	111.98(16)
O3-Al1-O2	94.85(9)	O3-Al3-O7 (x2)	99.07(8)
O3-Al1-O4	98.41(5)	O5-Al3-O1	78.81(14)
O3-Al1-O6	178.21(7)	O5-Al3-O7 (x2)	108.64(8)
O3-Al1-O7	92.33(10)	O7-Al3-O1 (x2)	76.63(7)
O4-Al1-O7	86.31(10)	O7-Al3-O7	128.09(13)
O6-Al1-O1	99.31(5)	O2-Al4-O5 (x2)	118.00(8)
O6-Al1-O2	83.37(9)	O2-Al4-O6 (x2)	110.98(8)
O6-Al1-O4	81.90(4)	O2-Al4-O2	106.04(11)
O6-Al1-O7	89.45(9)	O6-Al4-O5	92.17(17)
O1-Al2-O5	90.37(16)	O4-B1-O7 (x2)	119.89(10)
O1-Al2-O6	168.58(17)	O7-B1-O7	120.22(17)
O2-Al2-O1 (x2)	106.84(8)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	93.65(11)	Al2-O5-Al3	102.18(153)
Al1-O1-Al3 (x2)	90.21(14)	Al4-O5-Al2	104.45(42)
Al2-O1-Al1 (x2)	133.15(10)	Al4-O5-Al3	153.37(59)
Al2-O1-Al3	87.91(101)	Al1-O6-Al1	98.19(11)
Al2-O2-Al1	106.61(33)	Al1-O6-Al2 (x2)	92.59(11)
Al4-O2-Al1	123.53(31)	Al4-O6-Al1 (x2)	130.87(15)
Al4-O2-Al2	128.64(25)	Al4-O6-Al2	88.27(40)
Al1-O3-Al1	100.20(11)	Al3-O7-Al1	105.41(25)
Al3-O3-Al1 (x2)	129.90(14)	B1-O7-Al1	125.80(35)
Al1-O4-Al1	95.54(11)	B1-O7-Al3	128.70(34)
B1-O4-Al1 $(x2)$	130.95(12)		

Table A3.13: M-O-M angles at P0 (0.001 GPa), crystal in DAC.

Table A3.14: O-M-O angles at P0 (0.001 GPa), crystal in DAC.

Sites	Angle (°)	Sites	Angle ( $^{\circ}$ )
O2-Al1-O1	99.11(19)	O2-Al2-O6 (x2)	77.3(2)
O2-Al1-O4	92.87(23)	O2-Al2-O2	110.00(19)
O3-Al1-O1	81.18(18)	O5-Al2-O2 $(x2)$	117.92(25)
O3-Al1-O2	94.42(24)	O5-Al2-O6	77.25(45)
O3-Al1-O4	97.95(10)	O3-Al3-O1	169.05(63)
O3-Al1-O6	176.73(18)	O3-Al3-O5	114.34(113)
O3-Al1-O7	94.92(21)	O3-Al3-O7 (x2)	101.33(25)
O4-Al1-O1	168.02(16)	O5-Al3-O1	76.61(60)
O6-Al1-O1	98.75(10)	O5-Al3-O7 (x2)	105.97(21)
O6-Al1-O2	82.35(19)	O7-Al3-O1 (x2)	74.5(2)
O6-Al1-O4	82.79(23)	O7-Al3-O7	128.11(23)
O6-Al1-O7	88.29(25)	O2-Al4-O5 (x2)	119.32(18)
07-Al1-O1	80.49(31)	O2-Al4-O6 (x2)	110.82(26)
O7-Al1-O2	170.47(33)	O2-Al4-O2	105.52(19)
O7-Al1-O4	87.70(22)	O5-Al4-O6	90.03(66)
O1-Al2-O2 (x2)	107.69(22)	O4-B1-O7 (x2)	120.49(25)
O1-Al2-O5	93.30(48)	O7-B1-O7	118.97(30)
O1-Al2-O6	170.55(46)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	93.70(9)	Al2-O5-Al3	103.36(108)
Al1-O1-Al3 (x2)	89.47(14)	Al4-O5-Al2	103.74(38)
Al2-O1-Al1 (x2)	133.11(10)	Al4-O5-Al3	152.90(53)
Al2-O1-Al3	88.69(77)	Al1-O6-Al2 $(x2)$	90.6(1)
Al2-O2-Al1	105.33(34)	Al1-O6-Al1	98.73(9)
Al4-O2-Al1	124.79(32)	Al4-O6-Al1 (x2)	130.63(12)
Al4-O2-Al2	127.94(28)	Al4-O6-Al2	89.82(41)
Al1-O3-Al1	99.38(9)	Al3-O7-Al1	104.33(24)
Al3-O3-Al1 (x2)	130.30(12)	B1-O7-Al1	126.36(34)
Al1-O4-Al1	96.50(9)	B1-O7-Al3	128.8(3)
B1-O4-Al1 (x2)	130.36(10)		

**Table A3.15:** M-O-M angles at *P*1 (0.15(5) GPa).

**Table A3.16:** O-M-O angles at *P*1 (0.15(5) GPa).

		$(0.10(0) \ 01 \ a).$	
Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
O1-Al1-O7	82.52(20)	O2-Al2-O5 (x2)	118.01(25)
O2-Al1-O1	99.88(20)	O2-Al2-O6 (x2)	78.66(20)
O2-Al1-O4	91.80(22)	O2-Al2-O2	110.86(22)
O2-Al1-O6	84.28(21)	O5-A12-O6	76.13(42)
O2-Al1-O7	172.97(33)	O3-Al3-O1	169.63(57)
O3-Al1-O1	81.49(18)	O3-A13-O5	114.46(87)
O3-Al1-O2	94.06(23)	O3-Al3-O7 (x2)	100.41(24)
O3-Al1-O4	98.05(8)	O5-Al3-O1	75.91(52)
O3-Al1-O6	178.34(13)	O7-Al3-O1 (x2)	75.49(19)
O3-Al1-O7	92.83(19)	O7-Al3-O5 (x2)	107.08(20)
O4-Al1-O1	168.32(13)	O7-Al3-O7	127.37(20)
O4-Al1-O7	85.85(19)	O2-A14-O5 (x2)	118.17(18)
O6-Al1-O1	98.65(8)	O2-A14-O6 (x2)	112.02(26)
O6-Al1-O4	82.15(22)	O2-Al4-O2	105.62(22)
O6-Al1-O7	88.83(21)	O5-Al4-O6	90.31(60)
O1-Al2-O2 (x2)	107.55(23)	O4-B1-O7 (x2)	120.30(24)
O1-Al2-O5	92.03(46)	O7-B1-O7	119.40(27)
O1-Al2-O6	168.17(48)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
Al1-O1-Al1	94.08(11)	Al3-O5-Al2	104.59(88)
Al1-O1-Al3 (x2)	88.67(15)	Al3-O5-Al4	154.22(63)
Al2-O1-Al1 (x2)	132.87(12)	Al4-O5-Al2	101.19(45)
Al2-O1-Al3	88.67(65)	Al1-O6-Al2 (x2)	91.49(12)
Al2-O2-Al1	105.31(34)	Al1-O6-Al1	97.72(11)
Al4-O2-Al1	123.71(33)	Al4-O6-Al1 (x2)	131.04(13)
Al4-O2-Al2	129.74(31)	Al4-O6-Al2	91.31(44)
Al1-O3-Al1	98.15(11)	Al3-O7-Al1	104.11(27)
Al3-O3-Al1 (x2)	130.91(14)	B1-O7-Al1	126.73(37)
Al1-O4-Al1	96.80(11)	B1-O7-Al3	128.07(30)
B1-O4-Al1 (x2)	131.05(12)		

**Table A3.17:** M-O-M angles at *P*2 (0.91(4) GPa).

**Table A3.18:** O-M-O angles at *P*2 (0.91(4) GPa).

Sites	Angle ( $^{\circ}$ )	Sites	Angle (°)
O1-Al1-O2	98.52(22)	O2-Al2-O5 (x2)	117.46(24)
O1-Al1-O7	83.76(22)	O2-Al2-O6 (x2)	79.35(20)
O3-Al1-O1	82.39(19)	O2-Al2-O2	113.42(24)
O3-Al1-O2	93.49(26)	O5-A12-O6	76.81(44)
O3-Al1-O4	97.61(10)	O3-Al3-O1	169.26(63)
O3-Al1-O6	175.75(17)	O3-A13-O5	114.96(81)
O3-Al1-O7	93.65(21)	O3-Al3-O7 (x2)	99.99(25)
O4-Al1-O1	170.66(17)	O5-Al3-O1	75.78(54)
O4-Al1-O2	90.82(27)	O5-Al3-O7 (x2)	107.58(22)
O4-Al1-O7	86.92(21)	O7-Al3-O1 (x2)	75.67(21)
O6-Al1-O1	98.54(10)	O7-A13-O7	126.68(20)
O6-Al1-O2	82.28(23)	O2-Al4-O5 $(x2)$	118.28(19)
O6-Al1-O4	82.16(21)	O2-Al4-O2	106.07(24)
O6-Al1-O7	90.58(24)	O6-Al4-O2 (x2)	111.37(24)
O7-Al1-O2	172.74(36)	O6-Al4-O5	90.69(60)
O1-Al2-O5	90.96(51)	O7-B1-O4 (x2)	119.89(27)
O1-Al2-O6	167.77(50)	O7-B1-O7	119.88(27)
O2-Al2-O1 (x2)	106.86(22)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	94.26(12)	Al3-O5-Al2	103.61(80)
Al1-O1-Al3 (x2)	90.28(16)	Al4-O5-Al2	102.25(45)
Al2-O1-Al1 (x2)	132.87(13)	Al4-O5-Al3	154.13(63)
Al2-O1-Al3	89.32(60)	Al1-O6-Al1	98.85(11)
Al2-O2-Al1	105.28(36)	Al1-O6-Al2 (x2)	91.12(15)
Al4-O2-Al1	123.54(36)	Al4-O6-Al1 (x2)	130.57(14)
Al4-O2-Al2	129.99(32)	Al4-O6-Al2	88.23(42)
Al1-O3-Al1	98.56(11)	Al3-O7-Al1	104.05(31)
Al3-O3-Al1 (x2)	130.70(14)	B1-O7-Al1	125.63(40)
Al1-O4-Al1	96.99(11)	B1-O7-Al3	128.59(30)
B1-O4-Al1 $(x2)$	130.83(13)		

Table A3.19: M-O-M angles at P3 (1.99(6) GPa).

Table A3.20: O-M-O angles at P3 (1.99(6) GPa).

		11  gres at  1.5 (1.55(0)  Gr a).	
Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
O1-Al1-O7	81.81(32)	O2-Al2-O5 (x2)	118.11(26)
O2-Al1-O1	99.84(24)	O2-Al2-O6 (x2)	78.65(22)
O2-Al1-O4	91.92(25)	O2-Al2-O2	111.95(24)
O2-Al1-O7	173.34(42)	O5-Al2-O6	77.96(42)
O3-Al1-O1	81.44(21)	O3-Al3-O1	170.48(61)
O3-Al1-O2	94.54(26)	O3-Al3-O5	112.90(73)
O3-Al1-O4	97.92(10)	O3-Al3-O7 (x2)	99.83(27)
O3-Al1-O6	178.22(17)	O5-Al3-O1	76.62(52)
O3-Al1-O7	92.09(21)	O5-Al3-O7 (x2)	108.08(25)
O4-Al1-O1	168.24(17)	O7-Al3-O1 (x2)	76.35(25)
O4-Al1-O7	86.48(22)	O7-Al3-O7	127.39(20)
O6-Al1-O1	99.22(11)	O2-Al4-O5 (x2)	119.32(21)
O6-Al1-O2	83.71(23)	O2-Al4-O6~(x2)	109.83(27)
O6-Al1-O4	81.78(22)	O2-Al4-O2	105.76(25)
O6-Al1-O7	89.65(24)	O5-Al4-O6	91.56(60)
O1-Al2-O5	90.45(51)	O7-B1-O4 (x2)	118.83(30)
O1-Al2-O6	168.41(50)	O7-B1-O7	121.34(27)
O2-Al2-O1 (x2)	107.34(24)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )	
Al1-O1-Al1	93.81(14)	Al3-O5-Al2	101.48(85)	
Al1-O1-Al3 (x2)	89.09(15)	Al4-O5-Al2	101.13(49)	
Al2-O1-Al1 (x2)	133.05(12)	Al4-O5-Al3	157.39(69)	
Al2-O1-Al3	89.13(67)	Al1-O6-Al1	99.91(14)	
Al2-O2-Al1	104.66(39)	Al1-O6-Al2 (x2)	89.94(13)	
Al4-O2-Al1	124.94(38)	Al4-O6-Al1 (x2)	130.02(15)	
Al4-O2-Al2	128.78(38)	Al4-O6-Al2	88.27(48)	
Al1-O3-Al1	98.25(14)	Al3-O7-Al1	102.63(31)	
Al3-O3-Al1 (x2)	130.74(15)	B1-O7-Al1	127.21(41)	
Al1-O4-Al1	98.03(14)	B1-O7-Al3	128.80(34)	
B1-O4-Al1 (x2)	129.40(13)			

Table A3.21: M-O-M angles at P4 (3.32(6) GPa).

Table A3.22: O-M-O angles at P4 (3.32(6) GPa).

Sites	Angle (°)	Sites	Angle (°)
01-Al1-07	83.37(25)	O2-Al2-O2	112.01(29)
O2-Al1-O1	99.82(26)	O2-A12-O5 (x2)	117.95(26)
O2-Al1-O4	92.37(29)	O2-Al2-O6 (x2)	78.82(22)
O2-Al1-O7	173.80(39)	O5-A12-O6	77.24(48)
O3-Al1-O1	81.77(19)	O3-Al3-O1	172.65(66)
O3-Al1-O2	94.51(30)	O3-Al3-O5	109.66(76)
O3-Al1-O4	98.26(13)	O3-Al3-O7 (x2)	99.78(26)
O3-Al1-O6	179.16(21)	O5-Al3-O1	77.69(58)
O3-Al1-O7	91.20(23)	O5-Al3-O7 (x2)	108.81(23)
O4-Al1-O1	167.77(20)	O7-Al3-O1 (x2)	77.29(24)
O4-Al1-O7	84.40(23)	O7-Al3-O7	128.12(23)
O6-Al1-O1	99.05(13)	O2-Al4-O5 (x2)	117.95(21)
O6-Al1-O2	85.48(27)	O2-Al4-O6 (x2)	110.98(27)
O6-Al1-O4	80.90(21)	O2-Al4-O2	105.21(29)
O6-Al1-O7	88.78(26)	O6-Al4-O5	93.36(67)
O1-Al2-O2 (x2)	106.91(24)	O4-B1-O7 (x2)	120.49(29)
O1-Al2-O5	91.70(57)	O7-B1-O7	118.94(30)
O1-Al2-O6	168.94(58)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	95.78(14)	Al3-O5-Al2	100.93(92)
Al1-O1-Al3 (x2)	90.40(16)	Al4-O5-Al2	101.22(53)
Al2-O1-Al1 (x2)	132.09(14)	Al4-O5-Al3	157.85(76)
Al2-O1-Al3	87.87(65)	Al1-O6-Al2 (x2)	90.83(15)
Al2-O2-Al1	105.87(40)	Al1-O6-Al1	99.83(14)
Al4-O2-Al1	124.11(40)	Al4-O6-Al1 (x2)	130.07(15)
Al4-O2-Al2	129.00(35)	Al4-O6-Al2	90.15(50)
Al1-O3-Al1	98.79(14)	Al3-O7-Al1	103.54(35)
Al3-O3-Al1 (x2)	130.59(15)	B1-O7-Al1	124.68(45)
Al1-O4-Al1	96.21(14)	B1-O7-Al3	131.20(35)
B1-O4-Al1 (x2)	130.18(14)		

Table A3.23: M-O-M angles at P5 (4.78(6) GPa).

Table A3.24: O-M-O angles at P5 (4.78(6) GPa).

Table A5.24. 0-M-0 angles at 1.5 (4.16(0) G1 a).				
Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )	
O1-Al1-O4	168.12(20)	O2-Al2-O5 (x2)	116.73(28)	
01-Al1-O7	82.28(31)	O2-Al2-O6 (x2)	78.34(25)	
O2-Al1-O1	99.32(27)	O2-A12-O2	113.47(27)	
O2-Al1-O4	92.52(28)	O5-A12-O6	76.69(49)	
O2-Al1-O7	172.73(45)	O3-Al3-O1	168.25(68)	
O3-Al1-O1	81.02(21)	O3-Al3-O5	111.88(81)	
O3-Al1-O2	93.43(30)	O3-Al3-O7 $(x2)$	99.16(29)	
O3-Al1-O4	97.54(13)	O5-Al3-O1	79.86(60)	
O3-Al1-O6	176.87(21)	O5-Al3-O7 $(x2)$	108.24(26)	
O3-Al1-O7	93.82(24)	O7-Al3-O1 (x2)	76.26(27)	
O4-Al1-O7	86.07(24)	O7-Al3-O7	128.86(23)	
O6-Al1-O1	100.32(13)	O2-Al4-O5 (x2)	118.59(23)	
O6-Al1-O2	83.57(26)	O2-Al4-O6 (x2)	111.12(28)	
O6-Al1-O4	81.73(22)	O2-Al4-O2	104.97(27)	
O6-Al1-O7	89.17(26)	O5-Al4-O6	91.94(69)	
O1-Al2-O5	91.34(58)	O4-B1-O7 (x2)	118.97(33)	
O1-Al2-O6	168.03(58)	O7-B1-O7	122.0(3)	
O2-Al2-O1 (x2)	107.67(26)			

Sites     Angle (°)     Sites				
	- · · /		Angle $(^{\circ})$	
Al1-O1-Al1	95.69(14)	Al3-O5-Al2	102.20(94)	
Al1-O1-Al3 $(x2)$	91.80(16)	Al3-O5-Al4	157.52(76)	
Al2-O1-Al1 (x2)	132.14(14)	Al4-O5-Al2	100.28(53)	
Al2-O1-Al3	88.78(67)	Al1-O6-Al2 (x2)	90.52(15)	
Al2-O2-Al1	105.96(40)	Al1-O6-Al1	99.67(14)	
Al2-O2-Al4	129.29(35)	Al4-O6-Al1 (x2)	130.16(15)	
Al4-O2-Al1	123.49(41)	Al4-O6-Al2	89.55(50)	
Al1-O3-Al1	98.07(14)	Al3-O7-Al1	103.46(36)	
Al3-O3-Al1 (x2)	130.96(15)	B1-O7-Al1	125.65(47)	
Al1-O4-Al1	97.00(14)	B1-O7-Al3	130.04(38)	
B1-O4-Al1 (x2)	130.33(15)			

Table A3.25: M-O-M angles at P6 (5.83(5) GPa).

**Table A3.26:** O-M-O angles at *P*6 (5.83(5) GPa).

Sites	Angle (°)	Sites	Angle (°)
O1-Al1-O4	167.39(21)	O2-Al2-O5 (x2)	117.81(28)
01-Al1-07	81.12(42)	O2-Al2-O6 (x2)	78.38(25)
O2-Al1-O1	100.94(27)	O2-Al2-O2	112.35(27)
O2-Al1-O4	91.65(30)	O5-Al2-O6	78.15(50)
O2-Al1-O6	83.91(26)	O3-Al3-O1	168.73(70)
O2-Al1-O7	172.75(46)	O3-A13-O5	111.55(83)
O3-Al1-O1	80.96(21)	O3-Al3-O7 (x2)	99.14(29)
O3-Al1-O2	94.0(3)	O5-Al3-O1	79.73(62)
O3-Al1-O4	97.63(13)	O5-Al3-O7 (x2)	108.57(27)
O3-Al1-O6	177.66(21)	O7-Al3-O1 (x2)	76.41(27)
O3-Al1-O7	93.20(26)	O7-Al3-O7	128.42(25)
O4-Al1-O7	86.47(26)	O2-Al4-O5 $(x2)$	118.40(23)
O6-Al1-O1	100.45(13)	O2-Al4-O6 (x2)	110.71(28)
O6-Al1-O4	81.41(23)	O2-Al4-O2	105.84(27)
O6-Al1-O7	88.88(27)	O6-Al4-O5	92.01(69)
O1-Al2-O5	89.30(59)	O4-B1-O7 (x2)	119.69(34)
O1-Al2-O6	167.45(59)	O7-B1-O7	120.61(33)
O2-Al2-O1 (x2)	108.03(27)		

Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )
Al1-O1-Al1	93.95(17)	Al3-O5-Al2	100.35(98)
Al1-O1-Al3 (x2)	88.44(19)	Al4-O5-Al2	102.99(59)
Al2-O1-Al1 (x2)	132.90(16)	Al4-O5-Al3	156.65(81)
Al2-O1-Al3	88.48(78)	Al1-O6-Al2 $(x2)$	90.19(17)
Al2-O2-Al1	104.88(44)	Al1-O6-Al1	99.39(16)
Al4-O2-Al1	124.38(43)	Al4-O6-Al1 (x2)	130.31(18)
Al4-O2-Al2	129.10(41)	Al4-O6-Al2	89.59(56)
Al1-O3-Al1	98.63(17)	Al3-O7-Al1	103.27(34)
Al3-O3-Al1 (x2)	130.40(18)	B1-O7-Al1	126.50(46)
Al1-O4-Al1	97.48(17)	B1-O7-Al3	128.84(38)
B1-O4-Al1 $(x2)$	129.76(17)		

Table A3.27: M-O-M angles at P7 (5.99(5) GPa).

Table A3.28: O-M-O angles at P7 (5.99(5) GPa).

Table A5.20. 0-10-0 angles at 1 1 (5.55(5) G1 a).				
Sites	Angle ( $^{\circ}$ )	Sites	Angle ( $^{\circ}$ )	
O1-Al1-O7	83.66(27)	O2-Al2-O5 (x2)	117.50(29)	
O2-Al1-O1	99.28(31)	O2-Al2-O6 (x2)	78.84(25)	
O2-Al1-O4	92.21(33)	O2-Al2-O2	111.95(32)	
O2-Al1-O6	84.91(30)	O5-A12-O6	75.84(53)	
O2-Al1-O7	173.77(48)	O3-Al3-O1	173.63(77)	
O3-Al1-O1	81.53(26)	O3-A13-O5	108.52(89)	
O3-Al1-O2	95.06(34)	O3-Al3-O7 (x2)	100.72(28)	
O3-Al1-O4	98.23(16)	O5-Al3-O1	77.85(67)	
O3-Al1-O6	179.61(25)	O7-Al3-O1 (x2)	76.77(27)	
O3-Al1-O7	90.79(27)	O7-Al3-O5 (x2)	108.47(26)	
O4-Al1-O1	168.48(25)	O7-Al3-O7	128.05(25)	
O4-Al1-O7	84.83(28)	O2-A14-O5 (x2)	118.25(24)	
O6-Al1-O1	98.85(16)	O2-Al4-O6 $(x2)$	110.96(30)	
O6-Al1-O4	81.39(27)	O2-Al4-O2	106.08(32)	
O6-Al1-O7	89.22(30)	O5-Al4-O6	91.58(78)	
O1-Al2-O2 $(x2)$	106.79(27)	O7-B1-O4 (x2)	119.79(32)	
O1-Al2-O5	93.32(66)	O7-B1-O7	120.22(33)	
O1-Al2-O6	169.16(67)			

Sites	Angle (°)	Sites	Angle (°)
51005			,
Al1-O1-Al1	92.10(19)	Al3-O5-Al2	101.83(121)
Al1-O1-Al3 $(x2)$	87.18(21)	Al4-O5-Al2	100.62(59)
Al2-O1-Al1 (x2)	133.78(17)	Al4-O5-Al3	157.55(84)
Al2-O1-Al3	89.56(87)	Al1-O6-Al2 (x2)	89.93(17)
Al2-O2-Al1	105.17(46)	Al1-O6-Al1	96.96(19)
Al4-O2-Al1	124.18(44)	Al4-O6-Al1 (x2)	131.52(21)
Al4-O2-Al2	129.58(42)	A14-O6-A12	90.26(61)
Al1-O3-Al1	99.94(19)	Al3-O7-Al1	103.38(35)
Al3-O3-Al1 (x2)	129.84(21)	B1-07-Al1	127.76(47)
Al1-O4-Al1	99.84(19)	B1-O7-Al3	127.66(38)
B1-O4-Al1 (x2)	128.89(19)		

Table A3.29: M-O-M angles at P8 (6.45(6) GPa).

**Table A3.30:** O-M-O angles at *P*8 (6.45(6) GPa).

Sites	Angle (°)	Sites	Angle (°)
O2-Al1-O1	98.38(32)	O2-Al2-O5 (x2)	117.61(31)
O2-Al1-O6	83.76(32)	O2-Al2-O6 (x2)	79.45(25)
O2-Al1-O7	173.22(49)	O2-A12-O2	113.01(32)
O3-Al1-O1	82.08(27)	O5-A12-O6	76.52(57)
O3-Al1-O2	94.92(37)	O3-Al3-O1	174.88(82)
O3-Al1-O4	100.03(19)	O3-A13-O5	108.36(98)
O3-Al1-O6	178.13(29)	O3-Al3-O7 (x2)	100.85(31)
O3-Al1-O7	91.79(29)	O5-Al3-O1	76.76(74)
O4-Al1-O1	168.68(28)	O5-Al3-O7 (x2)	108.15(27)
O4-Al1-O2	92.53(36)	O7-Al3-O1 (x2)	77.16(28)
O4-Al1-O6	81.36(30)	O7-Al3-O7	128.61(25)
O4-Al1-O7	85.39(29)	O2-Al4-O5 $(x2)$	118.88(24)
O6-Al1-O1	96.80(18)	O2-Al4-O2	105.00(32)
O6-Al1-O7	89.55(34)	O6-Al4-O2 (x2)	110.38(32)
07-Al1-O1	83.43(31)	O6-Al4-O5	92.61(87)
O1-Al2-O2 (x2)	106.52(28)	O4-B1-O7 (x2)	121.02(33)
O1-Al2-O5	91.85(69)	O7-B1-O7	117.62(33)
O1-Al2-O6	168.36(69)		

# **B** Powder X-ray Diffraction Tips & Tricks

Appendix B is a reprint of a Manual for using the PANalytical X'Pert Pro MPD Powder Diffractometer at the Institute of Geological Sciences, University of Bern

## B.1 Configuring a Measurement Program for the X'Pert Pro Diffractometer

Most used programs are Absolute Scan, which contains information for a simple Theta-2-Theta scan, General Batch, which can contain several batched absolute scans combined with diffractometer settings and Non-ambient Program (a special batch for the HTK 1200 furnace).

#### Absolute Scan (Theta-2-Theta Scan Program)

To create a new absolute scan, choose 'New' from the 'File' menu in X'Pert DataCollector software. When asked for the measurement type, choose 'Absolute Scan'.

🚯 Prepare Absolute Scan [Program1]	]		_ 🗆 ×
Configuration Spinner Juni 2008 Scan Axis Gonio	Scan mode C Step C Continuous C Pre-set counts		Comment Settings
Cother gonio angle Cother gonio	Start angle (*): End angle (*): Step size (*): Time per step (s):	5.000 75.000 0.020 0.50	
Repeat       Number of repeats:       1       Total time (h:m:s):	Scan speed (*/s): Pre-set counts (counts): Number of steps:	0.040000	
	Total time (h:m:s):	00:29:13	

Fig. B1.1: Absolute scan program window.

In the scan program window (Figure B1.1), select a configuration: 'Spinner' for normal usage, 'HTK' for high-temperature measurements or 'Bracket'.

The scan axis is always 'Gonio' and scan mode is 'Continuous'. At this time, no diffractometer setup is defined and the measurement program will not work correctly!

Click 'Settings...'. A new window opens (Figure B1.2) where the diffractometer setup can be entered. At the moment, all settings are 'actual', which means that they are undefined.

Click through all items in the tree list and change all values highlighted in red from 'actual' to the setting as in Figure B1.2. Values can be changed in the field highlighted green. This field is different for each item in the tree list. Items highlighted in red must be defined exactly as shown in Figure B1.2, items highlighted in yellow depend on the setup. Choosing the right setup is described in the part 'Configuring the X'Pert Pro MPD Diffractometer'. Do not enter any offsets!

Blue values are sent to the diffractometer by the DataCollector software and the hardware is automatically adjusted. Black values are for information purposes only. They are written

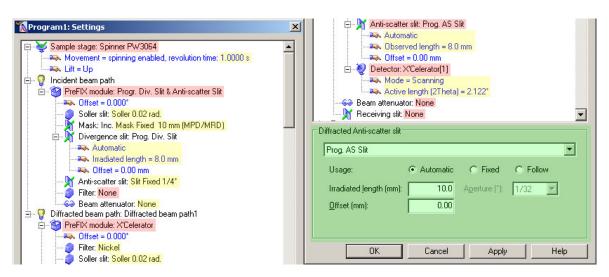


Fig. B1.2: Measurement program settings.

to the .xrdml file but the hardware configuration must be manually changed and verified by the user. When finished, click 'Apply' and then 'OK'. The step size has now changed in the 'Absolute Scan Program' window.

**2-Theta start and end angle.** The required scan range depends on the purpose of the measurement. For phase identification,  $10^{\circ}$  to  $55^{\circ}$  2-Theta is usually enough because most PDF entries do not contain high-angle data. Make sure the scan range covers all strong peaks. For lattice parameter refinements, make sure that peaks with low hkl values are covered by the scan range. For crystal structure analysis, the more the better but high-angle data is essential.

**Step size.** The X'Celerator detector allows step sizes of 0.033xx, 0.016xx, 0.008xx, 0.004xx and 0.002xx °/step. For phase identification and phase quantification by the Rietveld method, a step size of 0.016xx °/step is usually small enough. For crystal structure refinements and determination of precise lattice parameters, a smaller step size is recommended.

**Time per step.** Longer step times result in higher intensities and better counting statistics. For phase identification, there is no need to acquire data with very high intensities. As long as the peaks separate well from the noise, they can be identified. For phase quantification, a maximum intensity of 1'500 to 2'000 counts is usually good enough. For structural analyses, ca. 10'000 counts is a good start because there will still be peaks at high angles.

**Repeat.** A scan program can be repeated by entering a value in the 'Number of repeats' field. The .xrdml measurement file then contains more than one measurement which can be

summed or averaged in HighScore Plus. This is useful to reduce noise.

Users			
- Availabl	e to all users		
L Grandbi	e to all users		
	e tu all users		
Name	e to all users		

Fig. B1.3: Do not check 'Available to all users'!

When saving the program, enter a program name and a description of the program (optional). Be sure to uncheck the option 'Available to all users' (Figure B1.3).

#### General Batch (to set the generator to stand-by after measurement)

Use this program type to set up a batch of different measurement files, ambient settings or generator settings. The general batch is especially useful to set the generator to stand-by after the measurement.

Select 'New' from the 'File' menu and chose 'General Batch' from the program type pull-down menu.

In the 'General Batch' window (Figure B1.4), add one or more previously defined scan programs by clicking 'Insert Measurement Program...', if needed, add timer settings or batch settings. Batch settings can be different temperature (when using the HTK furnace) or generator settings. In a general batch the shutter is not closed between individual measurements. This can result in radiation damage on the sample surface. The following picture shows a batch that sets the generator to stand-by after completion of a single measurement:

- Item number 1 is the measurement program (added by clicking 'Insert Measurement Program...'),
- Item number 2 is a general batch setting to set the generator to 30 kV, 10 mA.

#### Non-Ambient Program (for high-temperature measurements)

The 'Non-ambient program' (Figure B1.5) is tailored for high-temperature measurements. It is set up the same way as the general batch but together with the diffraction patterns, an .xrdml file containing measured temperature and temperature set-points is saved.

When performing high-temperature measurements, it is recommended to insert a timer setting before a measurement is started to allow temperature equilibration in the furnace (see the example in Figure B1.5). It is not recommended to use the maximum heating rate. Try ca

NPrepare Batch [Program3]			TA Pr	epare No	on-ambient Program [Program	m5]	
Configuration: Spinner Juni 2008	Total time (h:m:s): 02:55:36		Con	figuration:	HTK Juni 2008		
Number Command				Number	Command		File Name
1 Fisch_FixDiv;				start	Temperature = 25; Time = 00:00;	:00; Slope = 60;	
2 kV = 30; mA = 10;				2	Fisch_AutDiv;		C:\Documents and Setti
*				3	Temperature = 100; Time = 00:01	1:15; Slope = 60;	
General Batch Settings				4	Wait 00:10:00;		
	usl		_	5	Fisch_AutDiv;		C:\Documents and Setti
Position Changer/Non-ambient X-ray/P	HU	OK		6	Temperature = 200; Time = 00:0*	1:40; Slope = 60;	
X-Ray generator		Cancel		/	Wait 00:10:00;		010 10 10 10
Switch: C On	C Off	Cancer	-	8	Fisch_AutDiv;	FF. Cl	C:\Documents and Setti
		Help		reset finish	Temperature = 25; Time = 00:02; Temperature = 25; Time = 00:00;		
I Set: kV:	30 mA: 10			nnisri	Temperature = 25, Time = 00.00.	.00, Slupe = 60,	1
				Non-an	nbient Setting		×
					Set	Increment	ОК
				Femperatu		o [	
Insert program item				remperatu	ie ( c). (0   200	<u>ا ا</u>	Cancel
Insert Measurement Program	Define file names with measured	arements					
Invest Dedah Celling	Cie Marca Callinar	10		Use Maxin	num Slope 🔽		Help
Insert Batch Settings	File Name Settings			r: a	s)· 00:00:00		
Insert Timer Settings	List of File Names	1	-	[ime (h:m::	s): 00.00.00		
				olope (*C/i	min ): 60		
1	1		-				
Comment Modify Dele	ete Cycles C	Check	-				

Fig. B1.4: General batch program.

Fig. B1.5: Non-ambient program.

### $10^{\circ}\mathrm{C/min}.$

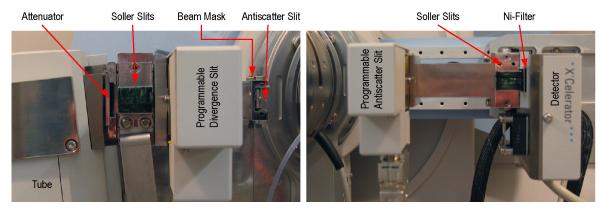
In a non-ambient program the shutter is not closed between individual measurements. This can result in radiation damage on the sample surface.

# **Running a Measurement Program**

After connecting to the diffractometer ('Instrument' menu  $\rightarrow$  'Connect'), choosing the required configuration and setting the generator to 40kV, 40mA, select 'Program...' from the 'Measure' menu and choose a previously defined program. Before the program starts, program name and comments are displayed and the file name and file path can be entered (Figure B1.6).

🗙 Start		×
Program Name:	Scan4-60 rot	
Type:	Absolute scan	
Description:		
File		
<u>N</u> ame:	tutorial1,xrdml	r El
<u>F</u> older:	E:\Documents and Settings\All Users\Documents\N	
Comment:		<b>A</b>

Fig. B1.6: Starting the measurement.



# B.2 Configuring the X'Pert Pro MPD Diffractometer

Fig. B1.7: Primary and secondary beam path PreFix modules.

# **Primary Beam Path Options**

**Attenuation foils.** Cu attenuation foils of 0.1mm and 0.2mm thickness are available (Figure B1.8). They are inserted directly after the X-ray tube in the primary beam path (Figure B1.7) and used for alignment purposes only.



Fig. B1.8: Cu attenuator foil.

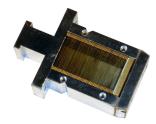


Fig. B1.9: Soller slits.

**Soller Slits.** Soller slits of 0.02 rad. and 0.04 rad. opening (Figure B1.9) can be inserted in the primary and secondary beam path (Figure B1.7). Use the same soller slit size in the primary and diffracted beam path. Smaller soller slits reduce axial divergence resulting in lower peak asymmetry at low 2-Theta angles (Figure B1.10), however, intensity is vastly reduced. Whenever possible, use the 0.02 rad. soller slits.

# Handle them with care. Do not touch the lamellae!

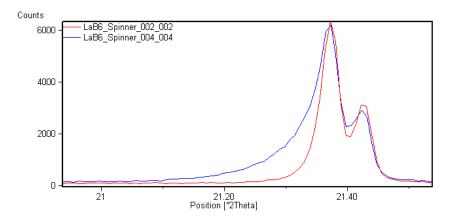


Fig. B1.10: Scans with 0.04 rad. (blue) and 0.02 rad. soller slits (red). Intensities are normalized.

**Programmable divergence slit.** The divergence slit in the primary beam path (Figure B1.7) reduces the equatorial divergence of the primary beam (in the plane perpendicular to the sample, parallel to the primary beam).

🕅 ADS: Settings			×
Divergence sli Automatic Irradiated	length = 10.0 mn		
Filter: None Mask: Inc. Ma Soller slit: Solle Diffracted beam pa Diffracted beam pa Automatic Mode = 0 Detector: X'Ce Mode = Sol	ath: Diffracted be Prog. AS Slit length = 10.0 mr 00 mm lerator[1]	am path1	
Divergence slit	gth (2Theta) = 2.	122*	
Prog. Div. Slit			•
Usage: Irradiated length (mm):	Automatic     10.0	C Fixe	d T
Offset (mm):	0.00		
OK	Cancel	Apply	Help

Fig. B1.11: Options for fixed or automatic divergence slits.

In the measurement program settings, check the options marked yellow in Figure B1.11 and be sure that the same setting is applied to both the programmable divergence slit in the incident beam path and the programmable antiscatter slit in the diffracted beam path. The programmable antiscatter slit in the diffracted beam path can also be set to 'follow'.

**Fixed slit mode:** The divergence of the incident beam is constant over the 2-Theta range (Figure B1.12). This agrees with the assumption of a constant sample volume: At low angles, a large sample area but low depth is irradiated. At high angles, a small area but large sample depth is irradiated. Large divergence slits can result in primary beam overspill at low 2-Theta angles!

**Constant irradiated length mode:** The slit size varies with 2-Theta and the irradiated sample surface is kept constant. This results in higher intensities, better s/n ratio at high 2-Theta and is a favorable option for fast data collection, e.g. phase identification. Usually, a length of 10 mm to 20 mm is appropriate. However, the assumption of constant volume is no longer valid, relative intensities are not correct and resolution is a bit lower (Figure B1.13). For matching intensities with PDF files, the pattern must be intensity-corrected to a fixed divergence of 0.5°. This is done in HighScore Plus in the 'Treatment' menu  $\rightarrow$  'Corrections'  $\rightarrow$  'Convert Divergence Slit'.

**Samples prepared on silicon plates:** As the sample is usually very thin, constant volume is not guaranteed. Samples prepared on silicon plates must be measured with automatic slits. Do not set the irradiated length larger than the physical length of the sample.

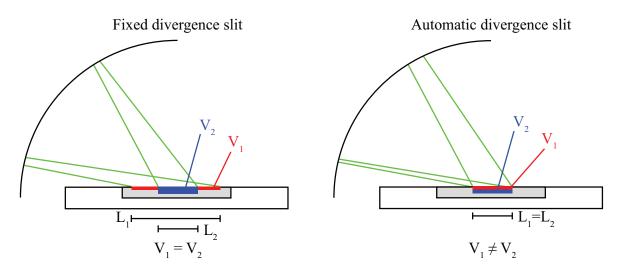


Fig. B1.12: Difference in sample volume when measuring with fixed or automatic divergence slits.

#### Choosing the right slits...

When using automatic divergence slits with constant irradiated area, the illuminated length should not be set larger than the sample length.

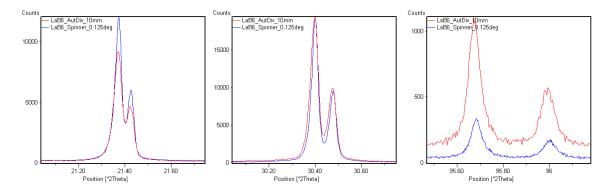


Fig. B1.13: The same pattern was measured with a fixed divergence slit of  $1/8^{\circ}$  (blue) and a constant irradiated area of 10 mm (red). Intensities were normalized to the peak at 30.4° 2-Theta. Note different relative intensities and the lower resolution at low angles.

When using fixed divergence slits, use the table below to determine correct slit size. Irradiated sample length with fixed slits can be calculated according to Equation B.1 and most often used slit sizes and irradiated lengths at the sample position are listed in Table B1.1. Selecting the correct slit size depends on the starting 2-Theta angle of the measurement and the sample size. For a 25 mm sample, the lowest starting angle with a  $1/4^{\circ}$  divergence slit is 5° 2-Theta. Select the slit as large as possible, larger slits result in higher intensities.

Irradiated length (mm) = R<sub>Goniometer</sub> (mm) 
$$\cdot \frac{\text{DivSlit}(^{\circ}) \cdot \frac{\pi}{180}}{\sin\theta(^{\circ})}$$
 (B.1)

2-Theta (°)	$1/16^{\circ}$	$1/8^{\circ}$	$1/4^{\circ}$	$1/2^{\circ}$	1°	$2^{\circ}$
1	30.0	60.0	120.0	240.0	480.0	960.0
2	15.0	30.0	60.0	120.0	240.0	480.0
3	10.0	20.0	40.0	80.0	160.0	320.0
4	7.5	15.0	30.0	60.0	120.0	240.0
5	6.0	12.0	24.0	48.0	96.0	192.1
6	5.0	10.0	20.0	40.0	80.0	160.1
7	4.3	8.6	17.2	34.3	68.6	137.2
8	3.8	7.5	15.0	30.0	60.0	120.1
9	3.3	6.7	13.3	26.7	53.4	106.8
10	3.0	6.0	12.0	24.0	48.1	96.1
11	2.7	5.5	10.9	21.9	43.7	87.4
12	2.5	5.0	10.0	20.0	40.1	80.1

Table B1.1: 2-Theta dependent irradiated length in mm at sample position with slit size.

**Beam mask.** The beam mask (Figure B1.14) limits the axial beam width (parallel to the sample surface on the goniometer axis) on the sample. They are inserted in the primary beam path after the divergence slit (Figure B1.7). Masks with a beam width of ca. 5 mm, 10 mm, 15 mm and 20 mm at the sample position are available. The 10 mm mask suits most requirements.





Fig. B1.14: Beam masks.

Fig. B1.15: Fixed antiscatter slits.

Antiscatter slit. The antiscatter slit (Figure B1.15) reduces scattering of the incident X-ray beam (Figure B1.7). When measuring with a fixed divergence slit, the size of the antiscatter should be double the size of the divergence slit. If automatic divergence slits are used, remove it or make sure that the antiscatter slit is not smaller than the maximum opening of the automatic divergence slit: With an irradiated length of 10 mm, the  $2^{\circ}$  fixed antiscatter slit can be used for measurements up to  $60^{\circ}$  2-Theta, with 8 mm irradiated length it can be used up to  $80^{\circ}$  2-Theta (Table B1.2). Select the smallest slit possible. For measurements with automatic divergence slits starting at angles below  $4^{\circ}$  2-Theta, a fixed antiscatter slit must be inserted!

Table B1.2: Max. 2-Theta angle with constant irradiated length and fixed antiscatter slits.

Irradiated Length	$5 \mathrm{mm}$	$7.5 \mathrm{~mm}$	$10 \mathrm{~mm}$	$12.5~\mathrm{mm}$	$15 \mathrm{~mm}$	$17.5~\mathrm{mm}$	$20 \mathrm{~mm}$
4° Slit (6 mm)	180°	$180^{\circ}$	$180^{\circ}$	$110^{\circ}$	$85^{\circ}$	$70^{\circ}$	$60^{\circ}$
2° Slit (3 mm)	180°	$85^{\circ}$	$60^{\circ}$	$48^{\circ}$	$40^{\circ}$	$35^{\circ}$	$30^{\circ}$

#### **Diffracted Beam Path Options**

**Programmable antiscatter slit.** The programmable antiscatter slit in the diffracted beam path (Figure B1.7) reduces air and sample scattering that does not contribute to Bragg reflections. In order to make sure that the detector observes all of the irradiated sample, this slit has to be set to the same setting as the programmable divergence slit. Set it to the same irradiated length as the divergence slit, to the same fixed divergence or set it to 'follow' in the measurement program settings.

**Soller slits.** Diffracted beam path soller slits (Figure B1.7) serve the same purpose as in the incident beam path and should have the same size.

Filter. The nickel filter is inserted before the detector (Figure B1.7) to cut off K $\beta$  radiation. Do not remove it.

**X'Celerator detector.** The X'Celerator detector (Figure B1.7) is a linear position sensitive multiple strip (LPSD) solid state detector. It consists of 127 channels, which work like an array of 127 individual point detectors. For normal measurements it should be set to 'Scanning' mode with an active area of 2.122°. Reducing the active area may decrease tangential aberration (because the detector is not only sensitive at a single point on the goniometer radius but at an area that is tangential to the goniometer circle). Since each strip of the detector works similar to a single point detector, the height of each strip (0.07 mm) is comparable to the detector slit in conventional point detectors.

# **B.3 Theory on Rietveld Refinements**

#### General Remarks on Data Quality

- Data quality is crucial! Read McCusker et al. (1999) and Toby (2006),
- Use the 0.02 rad. soller slits in any case,
- Acquire data with good s/n ratio,
- Acquire data with a max. count rate of the strongest peak of ca. 10'000 counts,
- Careful sample preparation is crucial (backloaded sample holders),
- Samples prepared on silicon plates must be measured with automatic divergence slits,
- For quantification, the 2-Theta range must cover all strong peaks,
- For precise lattice parameter determination, measure all peaks with low *hkl* indices,
- For structure refinements, measure at least up to 80° 2-Theta and use single-phase samples.

#### How a Rietveld Refinement Works

A calculated pattern of one or more phases is fitted to the measured one. Each phase is linear scaled by the scale factor. Parameters influencing peak positions (lattice parameters and space group) and relative peak intensities (atomic content of the unit cell and preferred orientation correction) are refined until the fit is satisfying. Thus, all phases present in the sample and their crystal structure must be known.

Crystal structure information is usually available from .cif files (Crystallographical Information File). Most important sections of a LaB6.cif file are described below (Table B1.3).

#### Calculated reflection positions depend on:

- Wavelength  $\lambda$  (for Cu X-radiation  $\lambda = 1.540598$  Å),
- Specimen displacement (mm) and zero shift (° 2-Theta),
- Lattice parameters a, b, c (Å) and  $\alpha, \beta, \gamma$  (°),
- Crystal system and space group.

The position of a reflection is calculated from Bragg's law (Equation B.2) and the line-spacing equation of the corresponding crystal-system (Equation B.3 for the cubic crystal system). Resulting values for LaB<sub>6</sub> with  $\lambda = 1.54$  Å are shown in Table B1.4.

$$\sin \theta = \frac{\lambda}{2 \cdot d} \tag{B.2}$$

_cell_length_a _cell_length_b	4.157 4.157	This part of the .cif file describes the unit cell			
_cell_length_c 4.157 _cell_angle_alpha 90.0		lengths and angles, formula units per unit cell			
		(Z = 1  means one La and six B atoms per  (Z = 1)			
_cell_angle_alpha	90.0	unit cell) and the space group setting.			
_cell_angle_alpha	90.0				
_symmetry_space_group_name_H-M	Pm-3m				
_symmetry_Int_Tables_number	221	Peak positions are calculated from this			
_cell_formula_units_Z	1	information.			
_atom_site_label		Atom label (name), atom type, site multiplic-			
_atom_site_type_symbol					
_atom_symmetry_multiplicity		ity and Wyckoff site (the latter two are also			
_atom_site_Wyckoff_symbol		defined by the space group symbol above),			
_atom_site_fract_x		fractional coordinates, $x, y$ and $z$ , isotropic			
_atom_site_fract_y		displacement parameters $B$ or $U$ and the site			
_atom_site_fract_z		occupacy (from $0$ to $1$ ) are defined.			
_atom_site_B_iso_or_equiv					
_atom_site_occupancy		Deele interesition and colored from this			
Lal La0+ 1 a 0.0 0.0 0.0 0.0	052 1.0	Peak intensities are calculated from this			
B1 B0+ 6 f 0.1975 0.5 0.5 0.0	041 1.0	information.			

Table B1.3: LaB6.cif file information relevant for Rietveld refinements.

$$\frac{1}{d^2} = \frac{h^2 + k^2 + l^2}{a^2} \tag{B.3}$$

			perior i cam posicio	$110$ of $202_6$ .
h	k	l	d-spacing (Å)	2-Theta (°)
1	0	0	4.157	21.380
1	1	0	2.937	30.407
1	1	1	2.399	37.463
2	1	0	1.858	48.978
2	1	1	1.696	54.009
2	2	0	1.469	63.238

**Table B1.4:** Peak positions of  $LaB_6$ .

Specimen displacement and zero shift cannot be refined together. Refining specimen displacement also accounts to some extent to sample transparency (for materials with low mass absorption coefficient) and should therefore be preferred.

#### Calculated reflection intensities depend on:

• Atom types,

- Atom coordinates (x, y, z between 0 and 1),
- Atomic displacement parameters B or U (usually a positive number < 3),
- Preferred orientation correction (value around 1, *hkl* direction must be known!).

Intensity of a reflection is calculated from the atomic content of the unit cell (structure factor F from atomic types and positions) and thermal vibration of atoms (displacement parameter B or U, with  $B = 8 \cdot \pi \cdot U^2$ ). If displacement parameters are missing in the .cif file, use 0.5 as start value. Equation B.4 describes intensity of a reflection in a powder diffractometer. Sample dependent parameters are the scale factor S, the structure factor F (atomic content of the unit cell), the atomic displacement parameter B and intensity corrections p (e.g. preferred orientation).

$$I \sim |F|^2 \cdot S \cdot p \cdot \left(\frac{1 + \cos^2 2\theta}{\sin^2 \theta \cdot \cos \theta}\right) \cdot e^{-2 \cdot \left(\frac{\sin \theta}{\lambda}\right)^2 \cdot B}$$
(B.4)

**Peak shape.** The shape of measured diffraction peaks is a convolution of peak shapes resulting from the diffractometer setup and microstructural properties of the phase like size of coherently scattering domains ('crystallite size' and microstrain). Calculated peaks are usually modeled with a pseudo Voigt function (linear combination of Gauss and Lorentz functions) with additional corrections for asymmetry and microstructural effects. The broader a peak, the smaller the 'crystallites' in the corresponding sample. In conventional powder diffractometers, peak width at half maximum (FWHM) is a function of Theta and is described by the Caglioti function (Equation B.5).

$$FWHM_k^2 = U \cdot \tan^2 \theta_k + V \cdot \tan \theta_k + W \tag{B.5}$$

The setup dependent peak shape and peak width of the diffractometer is determined by measuring a standard with (theoretically) infinite size of coherently scattering domains and no microstrain, such as NIST SRM 660a  $LaB_6$ .

#### Profile Shape Function Parameters in PANalytical HighScore Plus

The following parameters are valid for the X'Pert Pro MPD diffractometer when using a 10 mm beam mask, 0.02 rad. soller slits and the detector set to a sensitive area of 2.122°.

Intensities of data collected with constant irradiated sample length must be converted. In HighScore Plus, go to 'Treatment'  $\rightarrow$  'Corrections'  $\rightarrow$  'Convert Divergence Slit'.

After adding phases, double click 'Global Parameters' in the 'Lists Pane' (Figure B1.16) to open the 'Object Inspector' (Figure B1.17). In the 'Refinement' submenu (Figure B1.17), choose 'Pseudo Voigt 3(JFC Asymmetry)' and enter 20 in the 'Profile Base Width' field. In the 'Lists Pane', enter the values shown in Figure B1.16 in the 'Profile Parameters' submenu for each phase as start values. Calculated peaks will be too narrow because microstructural effects are not considered.

Lists Pane		
Pattern List   Scan List   Peak List   Anchor Scan	Data 0	Quantification
Name	Refine	Value
± Global Parameters		
三直 LaB6	~	
- Rofile Parameters		
U		0.001363
V		0.001565
W		0.000238
Anisotropic Broadening		0.000000
Asymmetry (Finger, Cox, Jephcoat) S/L		0.015000
Asymmetry (Finger, Cox, Jephcoat) D/L		0.015000
Peak Shape 1		0.699679
Peak Shape 2		-0.000891
Peak Shape 3		0.000029

Dbiect Ins Selected object: Global Parameters 🗄 Background Agreement Indices General Properties E Refinement Against lobs Weighting Scheme Use summed March/Dollase Function ~ Recycle HKL's (Le-Bail Fit only) **Profile Function** Pseudo Voigt 3(FJC Asymm Profile Base Width 20.00000 Intensity Limit 0.00000 Max. No. of Least-Squares Cycles 20 EPS (stop if all shifts < EPS\*esd) 0.50000 Job Type X-rays Lonat Cour . Ture Dankas

Fig. B1.16: Profile parameters in Lists Pane.

Fig. B1.17: Object Inspector options.

#### **Manual Refinement Steps**

Use steps 1, 2, 3 and 4 for quantitative phase analysis and steps 1, 2, 4, 5 and 6 for structural analysis. Check all values for their physical meaning (e.g. a specimen displacement of 0.5 mm or negative displacement parameters are physically not feasible!). A sample refinement of  $LaB_6$  is shown in Figure B1.18.

- 1. Refine specimen displacement (global parameters) and the scale factor for each phase,
- 2. Refine lattice parameters of each phase. Take care when refining lattice parameters with low intensities,
- 3. Refine B overall for each phase,
- 4. Refine profile parameters in the following sequence: U (if data > 90° 2-Theta, otherwise leave fixed!), V, W, Peak Shape 1, Peak Shape 2 for each phase. If the refinement is unstable, refine only V and W. Verify the values, especially for phases with low intensities. Instead of refining the values manually, the size-strain batch in HighScore Plus can be used. Set the refinement mode to 'Automatic' and choose 'size-strain' from the menu next to the refinement button. This batch also refines other parameters in the following order: Specimen displacement, scale, U, V, W, lattice parameters, shape 1 and shape 2,
- 5. Refine atomic positions. Never refine special positions (e.g.  $0, \frac{1}{4}, \frac{1}{3}, \frac{1}{2}, ...)$ ,

6. Refine individual B's of the atoms. Constraining B's of atoms (e.g. one parameter for each atom type) can help if the refinement is unstable.

Pattern List Scan List Peak List Anche	or Scan I	Data 🛛 Quan	tification R			X		0.000000	0.000000
Name	Refine	Value	Deviation			Y		0.000000	0.00000
🖃 🔁 Global Parameters						Z		0.000000	0.00000
Zero Shift [*2Theta]		0.000000	0.000000			B isotropic 6	3 🗹	0.018844	0.00486
Specimen Displacement [mm] 1	<b>v</b>	-0.045239	0.000083			Occupancy		1.000000	0.00000
Wavelength [Å]		1.540598	0.000000			+ B anisotropic			
🖃 🛱 LaB6	<b>&gt;</b>				닏	B1			
Scale factor 1	✓	0.008972	0.000012			X (	5	0.199256	0.00035
Preferred Orientation		1.000000	0.000000			Y _		0.500000	0.00000
B overall		0.000000				Z		0.500000	0.00000
Extinction		0.000000	0.000000			B isotropic 6		0.000000	0.02998
Flat Plate Absorption Correction		0.000000	0.000000			Occupancy		1.000000	0.00000
Porosity		0.000000	0.000000			+ B anisotropic			
Roughness		0.000000	0.000000	1-	A	Profile Parameters			
🔄 📑 🔤 Unit Cell					H	U 4.1		0.001547	0.00005
a [Å] 2	✓	4.156921	0.000003			V 4.2		0.001206	0.00005
ь [Å]	<b>v</b>	4.156921	0.000003		L	W 4.0	3	0.000350	0.00001
c [Å]	<b>v</b>	4.156921	0.000003			Anisotropic Broadening	<u>    </u>	0.000000	0.00000
alpha (*)		90.000000	0.000000			Asymmetry (Finger,Cox,Jeph.,		0.015000	0.00000
beta (*)		90.000000	0.000000			Asymmetry (Finger,Cox,Jeph.		0.015000	0.00000
gamma (*)		90.000000	0.000000			Peak Shape 1 4.4		0.737512	0.00703
- Atomic coordinates						Peak Shape 2 4.	5 🗹	-0.000815	0.00010
🖃 La1						Peak Shape 3		0.000029	0.00000

Fig. B1.18: Refined parameters are numbered according to the refinement sequence above. Step 3 was not performed because displacement parameters were refined individually for each atom.

# Agreement Factors (R-Values)

Agreement factors or *R*-values provide information about the quality of the fit of the calculated pattern to the measured one. They are only meaningful if all phases are included in the refinement. However, the most important way to determine the quality of a Rietveld fit is by viewing the observed and calculated patterns graphically (difference plot) and to ensure that the model is chemically plausible (McCusker et al. 1999 and Toby 2006).

- $R_{\rm wp}$  The weighted-profile *R*-factor (Equation B.6) is an indicator for the fit quality. Misfits at low intensities are weighted higher than misfits at peak tops.  $R_{\rm wp}$  is influenced by the background, therefore, a refinement with a high background and poor counting statistics can result in good  $R_{\rm wp}$  values even if the overall refinement is of low quality. Good values for  $R_{\rm wp}$  are around 10%.
- $R_{\text{exp}}$  The expected *R*-factor (Equation B.7) reflects data quality (counting statistics). It indicates the lowest possible  $R_{\text{wp}}$  when using available data with the current refinement model.
  - $\chi^2$  For a perfect refinement, the goodness of fit parameter  $\chi^2$  (Equation B.8) approaches

1. If data are over collected (i.e. high-intensity data where errors are dominated by the structural model and not the data),  $R_{\rm exp}$  is very small, resulting in a large  $\chi^2$ . If data have been collected too quickly,  $R_{\rm exp}$  will be large and  $\chi^2$  can be less than 1.

 $R_{\rm B}$  The Bragg *R*-factor (Equation B.9) is not used in the refinement process, however, it can be used to monitor the improvement of the model. The better the model, the lower *R*-Bragg.

$$R_{\rm wp} = \sqrt{\frac{\sum_{i} w_i (y_{i,obs} - y_{i,calc})^2}{\sum_{i} w_i (y_{i,obs})^2}}$$
(B.6)

$$R_{\exp} = \sqrt{\frac{N - P}{\sum_{i}^{N} w_i(y_{i,obs})^2}}$$
(B.7)

$$\chi^2 = \frac{R_{wp}}{R_{exp}} \tag{B.8}$$

$$R_{\rm B} = \sum_{hkl} \frac{|I_{hkl,obs} - I_{hkl,calc}|}{\sum_{hkl} |I_{hkl,obs}|} \tag{B.9}$$

Where  $y_{i,obs}$  is observed intensity at step *i*;  $y_{i,calc}$  is calculated intensity at step *i*;  $w_i$  is a weighting factor; *N* is number of data points; *P* is number of refined parameters and  $I_{hkl}$  is intensity of a Bragg reflection.

# B.4 High-Temperature Measurements Using the Anton-Paar HTK 1200 Furnace

## Mounting the Anton-Paar HTK 1200 Furnace





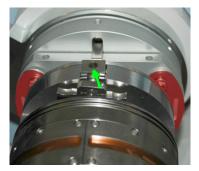


Fig. B1.19: Spinner unmount.

Fig. B1.20: Insert the screws.

Fig. B1.21: Attach the HTK.

Unmount the sample spinner by loosening the four black screws (red arrows, Figure B1.19). The spinner is still held by the latch at the top (green arrow, B1.19). Hold the spinner with one hand and lift the latch. Pull the spinner away towards you.

Insert the four hex-head bolts with washers and let them stick out ca. 1.5 cm (marked red, Figure B1.20).

With one hand, attach the HTK onto the goniometer while simultaneously lifting the latch. As soon as the latch has snapped in place, the HTK is fixed. Add the metal plates (marked red, Figure B1.21) between the HTK stage and the hex-head screws and tighten them carefully.



Fig. B1.22: Micrometer.



Fig. B1.23: Sample cup.



Fig. B1.24: Alignment slit.

Mount the micrometer on top of the motorized stage of the HTK (green arrow, Figure B1.22) and tighten it using the butterfly nut.

Screw a sample cup onto the HTK spinner (Figure B1.23). Tighten gently!

Attach the alignment slit onto the sample cup. The two slit openings must be perpendicular to the goniometer axis and face the X-ray tube and the detector (arrows, Figure B1.24).

Slide the spinner into the HTK by making sure the guidance is aligned (Figure B1.25). Tighten all four knurled screws gently by hand. **Do not switch on the sample spinner!** 

Slide in the 0.1 mm Cu attenuator foil in the primary beam path (Figure B1.26). **Do not continue without the attenuator foil!** Open the cooling-water tap (Figure B1.27) and check for leaking water. Switch on the heater (gray switch on HTK controller at the lower left of diffractometer).



Fig. B1.27: Water tap.

Fig. B1.25: Insert the spinner. Fig. B1.26: Insert attenuator.

Start the X'Pert DataCollector software. Before connecting to the instrument, go to the 'Tools' menu and select 'Exchange Sample Stage'. Change the sample stage to 'Anton Paar HTK-1200 Oven' and click 'Apply'. Confirm all messages and connect to the diffractometer.

# Adjusting the Height and Rotation of the High-Temperature Furnace

**Adjusting HTK rotation.** This is an essential step. With the Cu attenuator foil inserted, set the generator to 40 kV, 30 mA. In the same panel where generator settings are set (Instrument Settings Tab), apply the settings (and change/check them if necessary) as listed in Table B1.5.

	Tuble D1.0. Institutient settings for anguing the HTH fotation.					
Instrument settings	Incident beam optics	Diffracted beam optics				
Generator: 40 kV, 30 mA Temperature: 25° C Positions: All 0°	Fixed DivSlit $= 1/16^{\circ}$ Antiscatter Slit $=$ none Soller Slits $= 0.02$ rad. Mask $= 10$ mm	Fixed Antiscatter $= 4^{\circ}$ Soller Slits $= 0.02$ rad. Detector $=$ receiving slit				

Table B1.5: Instrument settings for aligning the HTK rotation.

Choose 'Manual Scan' from the 'Measure' menu and perform a manual scan of type 'Omega Scan' with a step size of  $0.02^{\circ}$ /step, a scan range of  $2^{\circ}$  Omega and a counting time of 0.25

sec/step. This moves the X-ray tube from  $-1^{\circ}$  to  $+1^{\circ}$  Omega (angle between the horizontal and the X-ray source) by keeping the detector at Omega + 180°.

When the scan has finished, immediately close the shutter by clicking on the arrow (four icons to the left of the stop icon) in the DataCollector toolbar. The resulting curve shows an intensity distribution of the primary beam passing through the alignment slit. The curve looks similar to Figure B1.28 with the maximum intensity most likely not at 0° Omega because the HTK is rotated.

Right-click into the measured peak and select 'Peak Parameters' from the popup menu. The appearing window shows peak information and the peak position (Figure B1.29).

In the 'Peak Parameters' window, press the 'Move To' button. The goniometer will move to that position, keeping the detector at Omega + 180°. Note the new position and the offset in the Instrument Settings Tab on the left side of DataCollector. This position will be the new zero position for the high-temperature measurement.

In the 'User Settings' menu, select 'Fine Calibration Offsets'. Press 'Set New = 0'. In the Instrument Settings Tab, new Omega and 2-Theta positions and the offset will now be zero.

Repeat the manual Omega scan. The peak position should now be close to  $0^{\circ}$  Omega. If this is not the case, reset the fine calibration offsets and start again.

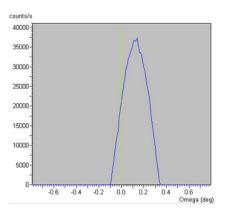


Fig. B1.28: Omega scan.

Parameters		Close
<sup>p</sup> osition (Omega (deg)):	0,124	Background>>
Net height (Counts/s):	36214,02	Print
Background height at peak position (Counts/s):	0,00	Move To
Net area (Omega (deg) * Counts/s):	9702	Help
3ackground area (Omega (deg) * Counts/s):	0	
WHM (Omega (deg)):	0,278	
ntegral breadth (Omega (deg)):	0,268	
FWHM / Integral breadth:	1,036	
Asymmetry:	0.98	

Fig. B1.29: Peak parameters window.

**Ajdusting HTK height.** Switch the micrometer on. Apply settings (and change/check them if necessary) as listed in Table B1.6.

Start a manual scan of '2-Theta' type with a step size of  $0.016xx^{\circ}/\text{step}$ , a scan range of 1° 2-Theta and a counting time of 2 s/step. This keeps the X-ray tube at 0° Omega and moves the detector from  $-0.5^{\circ}$  to  $+0.5^{\circ}$  2-Theta (angle between horizontal and the detector). Close the shutter immediately! If the resulting peak maximum is close to 0° 2-Theta, height

	8 8 8	0
Instrument settings	Incident beam optics	Diffracted beam optics
Generator: 40 kV, 30 mA Temperature: 25° C Positions: All 0°	Fixed DivSlit = $1/8^{\circ}$ Antiscatter Slit = none Soller Slits = 0.02 rad. Mask = 10 mm	Fixed Antiscatter $= 1/8^{\circ}$ Soller Slits $= 0.02$ rad. Detector $=$ scanning mode

Table B1.6: Instrument settings for aligning the HTK height.

is almost properly aligned. Right-click the peak and select 'Peak Parameters'. Write down intensity and position. If the peak maximum is at negative 2-Theta angles, the height is too low and vice versa.

Set the motor speed of the alignment stage to a low setting (black remote control on the right side of the diffractometer) and move the HTK 0.02 mm up. Perform the same scan again, close the shutter. Write down the intensity of the peak. If intensity is higher, go further upwards to find the maximum intensity. If intensity is lower, go back to 0 mm, then to -0.02 mm and repeat this step. When you have found the maximum intensity, the HTK is correctly aligned.

Always insert the beam attenuator when aligning to avoid detector damage!

# Never start a measurement with the micrometer attached. The X-ray tube and detector will hit the micrometer!

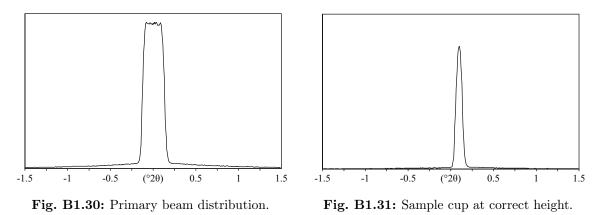
Remove the 0.1 mm Cu Attenuator before you start measuring.

When finished, reset all fine calibration offsets and change the sample stage back to spinner in DataCollector.

If you are lost...Coarse alignment when the HTK is completely misaligned. Be sure what you do and clear all fine calibration offsets before starting this procedure!

With the 0.1 mm Cu attenuator foil inserted and no sample spinner in the furnace, move Omega and 2-Theta to 0°. Set the same settings as in Table B1.5 and start a manual '2-Theta' scan with a scan range of 2° 2-Theta and 1 s/step to measure the intensity distribution of the primary beam (Figure B1.30). Select 'Peak Mode'. The primary beam center must be located at 0°. Print out this measurement. Mount an empty sample holder and perform the same manual scan again. Adjust the height of the HTK until the negative 2-Theta part of the primary beam is obscured by the sample cup (i.e. from  $-0.25^{\circ}$  to 0° in Figure B1.31). The maximum of the measured peak should be located in the middle between 0° 2-Theta and the right flank of the primary beam measured before (compare with the printout). The left flank of the peak should be close to 0° (Figure B1.31). Lowering the furnace will move the peak to lower angles, raising it will move the peak to higher angles. When an intensity distribution

as in Figure B1.31 is reached, continue with adjusting HTK rotation.



#### After the Measurement: Correct Diffraction Patterns for Height-Displacement

Due to thermal expansion of the HTK sample spinner, the sample moves out of the goniometer focus at high temperature. Peak positions will be shifted. To correct this, open the .xrdml file(s) in HighScore Plus Plus and select 'Correct Sample Displacement' from the 'Corrections' submenu in the 'Treatments' menu. Enter the negative value from Table B1.7. Note: HighScore Plus Plus appends data points with zero intensity at the end of the pattern, which must be clipped. For Rietveld refinements, it is also possible to use uncorrected patterns, however, height displacement starting values should be used according to Table B1.7.

Temp.	Displ.	Temp.	Displ.	Temp.	Displ.	Temp.	Displ.
25	0.000	325	0.098	625	0.220	925	0.360
50	0.010	350	0.105	650	0.230	950	0.370
75	0.020	375	0.113	675	0.240	975	0.380
100	0.030	400	0.120	700	0.250	1000	0.390
125	0.038	425	0.133	725	0.263	1025	0.405
150	0.045	450	0.145	750	0.275	1050	0.420
175	0.053	475	0.158	775	0.288	1075	0.435
200	0.060	500	0.170	800	0.300	1100	0.450
225	0.068	525	0.180	825	0.313	1125	0.465
250	0.075	550	0.190	850	0.325	1150	0.480
275	0.083	575	0.200	875	0.338	1175	0.495
300	0.090	600	0.210	900	0.350	1200	0.510

Table B1.7: Height displacement due to expansion of the HTK 1200 sample spinner (mm, °C).

# B.5 Rietveld Refinements with Topas Academic v. 4.1

# Using Topas Academic with jEdit

Topas Academic (Coelho 2007a) is operated with input text files containing specific keywords. The files are updated after each refinement. In order to easily access all keywords, a macro creating color-coded Topas input files for the java based text editor jEdit (www.jedit.org) is available (Figure B1.32). The macro was written by John S. O. Evans and can be obtained at the Topas Academic webpage of the University of Durham, UK at www.dur.ac.uk/john.evans/topas\_academic/topas\_main.htm. When both jEdit and Topas Academic programs are running, an input file edited in jEdit can be sent to Topas Academic by pressing the button marked red in Figure B1.32. All keywords can be accessed from the tree in the left side of the jEdit window (highlighted green in Figure B1.32).

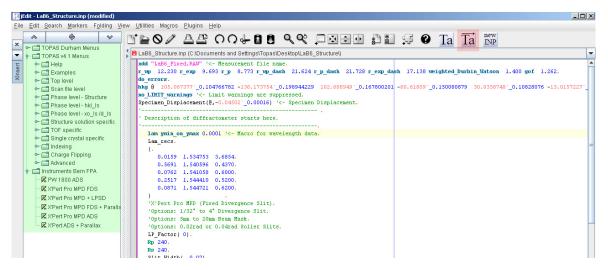


Fig. B1.32: Using jEdit with Topas Academic. The red button sends the input file to Topas. From the area highlighted in green, keywords can be inserted to the input file by simply clicking them.

# Peak Shape Modeling in Topas Academic

Topas allows peak shape modeling by using the Fundamental Parameter Approach (Cheary and Coelho 2004) in which peak profiles are simulated depending on the diffractometer setup instead of being determined from a standard material (usually NIST SRM 660a  $LaB_6$ ) and approximated using profile functions such as pseudo Voigt, split pseudo Voigt etc. Setting up a description of a laboratory diffractometer (PANalytical X'pert Pro MPD) for the fundamental parameter approach is shown in the commented example below:

```
lam ymin_on_ymax 0.0001
Lam recs
 { 0.0159 1.534753 3.6854
  0.5691 1.540596 0.4370
  0.0762 1.541058 0.6000
  0.2517 1.544410 0.5200
  0.0871 1.544721 0.6200 } 'the first seven lines describe the wavelengths
LP_Factor(0) 'monochromator angle (°)
Rp 240 'primary beam radius (mm)
Rs 240 'diffracted beam radius (mm)
Slit_Width(0.07) 'receiving slit width (0.07 mm for X'Celerator detector)
Divergence(0.125) 'divergence slit (°)
axial conv
filament_length 12 'length of X-ray tube filament (mm)
sample_length 10 'irradiated sample length parallel to the goniometer axis in (mm)
receiving_slit_length 15 'length of receiving slit parallel to the goniometer axis (mm)
primary_soller_angle 2.55 'opening of primary soller slits (°)
secondary_soller_angle 2.55 'opening of diffracted beam soller slits (°)
```

If the sample is measured with constant irradiated length instead of a fixed divergence slit, replace 'Divergence(0.125)' with the two keywords below:

```
Variable_Divergence_Shape(10) ' irradiated sample length (mm)
Variable_Divergence_Intensity
```

#### Creating an Input File for Topas Academic in jEdit

When creating an input file in jEdit for Topas Academic some important points have to be considered:

- Wavelength profiles are entered by clicking 'Emission profile(s)' from the 'Scan file level' branch. For laboratory diffractometers without primary beam monochromator, select 'CuKa5\_Berger.lam'. Use 'CuKa4\_Holzer.lam' if the instrument is equipped with a diffracted beam monochromator.
- The diffractometer description for the fundamental parameter approach can be inserted by selecting the corresponding diffractometer from the 'Instruments FPA Bern' menu. Check if you are using the correct configuration!
- The simplest way to add a structure is by clicking 'Insert CIFs in INP format' from the 'Phase level Structure' menu. Do not forget to add the scale factor ('scale') and 'CS\_L', 'CS\_G', 'Strain\_L' and 'Strain\_G' keywords to allow for peak broadening due to microstructural effects.
- Refining parameters in Topas is done by adding an @ sign in front of them. Basically, any character combination can be used as parameter name and this is the only way to refine more than one parameter simultaneously. With: x @ 0.1234 y @

0.1234 z @ 0.1234, each coordinate is refined individually. By replacing the @ symbol with any character combination (in this example 'refinetogether'), coordinates are refined together: x refinetogether 0.1234 y refinetogether 0.1234 z refinetogether 0.1234. If you do not want a parameter to be refined, simply remove the @ sign or write an ! before the parameter name: !refinetogether is not refined. Special positions must be input as fraction followed by a semicolon: do not enter x 0.33 or y 0.66 but x = 1/3; and y = 2/3; respectively. Note the = sign and the semicolon. Ignoring this can lead to very wrong results!

- Lattice parameters can be defined in two ways: For each crystal system, a macro is available (in the 'Lattice parameters' menu in jEdit). For example, the 'Hexagonal (@ 5.123, @ 6.321)' macro implies that a = b = 5.123 Å, c = 6.321 Å,  $\alpha = \beta = 90^{\circ}$  and  $\gamma = 120^{\circ}$ . If you do not want to use the macro, you can add lattice parameters by writing a, b, c, al, be, ga followed by @ and the value on a new line each. The hexagonal macro above is then replaced with: a @ 5.123, b = Get(a);, c @ 6.321, al 90, be 90, ga 120. For equivalent lattice parameters it is mandatory to use the =Get(); function or at least the same parameter name.
- When the input file is updated, values and standard deviations are displayed as follows:
  4.15 \\_0.01. The number before the \sign is the refined value. The number after the \_ sign is the estimated standard deviation (esd) of the value.

#### **Example Topas Input Files**

The following input files and corresponding patterns are from real refinements. However, they are printed here for demonstration purposes only and do not claim to be accurate as some parts have been deleted for reasons of clarity, some numbers have been shortened and esd values deleted.

**LeBail Refinement of LaB**<sub>6</sub>. An input file of a LeBail refinement of LaB<sub>6</sub> is listed below and a final refinement plot is shown in (Figure B1.33). In contrast to the Rietveld method, structure factors are not calculated from the structure but are set to an equal value at the beginning of the refinement. During the refinement, those values are used as start values to calculate 'observed intensities', which are then again used as new 'calculated intensities' and so on. The LeBail refinement can be used to determine structurally independent parameters (e.g. lattice parameters, peak shape, microstructural effects), and to extract intensities which can then be used for structure solution. In a LeBail refinement, intensities do not contribute to the correlation matrix.

```
xdd "LaB6_Fixed.XY"
x_calculation_step 0.01
r_wp 11.083 r_exp 9.693 r_p 7.801 r_wp_dash 19.685 r_p_dash 19.821
r_exp_dash 17.216 weighted_Durbin_Watson 1.724 gof 1.143
do_errors
bkg @ 109.216959 ... ... 9.0985598
no_LIMIT_warnings
Specimen_Displacement(@, -0.04025)
```

This section contains the measurement file name (xdd LaB6\_Fixed.XY), agreement factors, refined Chebyshev background parameters (two shown), the command to calculate errors (do\_errors), a command to not display limit warnings (no\_LIMIT\_warnings) and refined specimen displacement in mm. Because the diffraction pattern is input as an .xy file, the step size has to be defined using the keyword x\_calculation\_step.

```
lam ymin_on_ymax 0.0001
Lam_recs
 { 0.0159 1.534753 3.6854
  0.5691 1.540596 0.4370
  0.0762 1.541058 0.6000
  0.2517 1.544410 0.5200
  0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence(0.125)
axial_conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
```

This part contains information to describe diffractometer settings and the wavelength for the fundamental parameter approach. The sample was measured with a fixed divergence slit with an opening of  $0.125^{\circ}$ .

```
hkl_Is
lebail 1
phase_name LaB6
Cubic(@ 4.15712`_0.00000)
LVol_FWHM_CS_G_L(1, 2388.65106, 0.89, 3171.56454, csgc, 9765.63203, cslc, 4143.99750)
e0_from_Strain(0.00097, sgc1, 0.00010, slc1, 0.00390)
volume 71.8416653
space_group "Pm-3m"
load hkl_m_d_th2 I
{ 0      0      1      6      4.15711      21.35683      11.2092731
      ... not all shown }
```

The 'hkl\_Is' and 'lebail 1' keywords indicate that structural information for a LeBail refinement follows. Required information are the lattice parameters ('Cubic' macro where a = b = c and angles = 90°) and space group. The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\varepsilon_0 = \Delta d/d$ . As LaB<sub>6</sub> should have infinite crystallite size and no microstrain, the

pattern does not significantly change when no microstructural model is applied. The 'volume' keyword calculates the unit cell volume, 'space\_group' indicates the space group. The 'load hkl\_m\_d\_th2 I' keyword is added automatically in the first refinement step and includes all reflections according to the space group and lattice parameters with their *hkl*, multiplicity, d-spacing, 2-Theta values and intensities. In a LeBail refinement, no intensity errors are calculated.

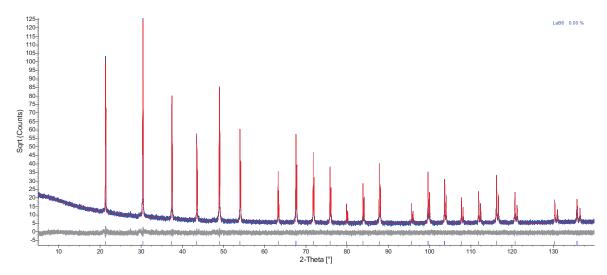


Fig. B1.33: Final refinement plot of the LeBail refinement of  $LaB_6$ . The blue curve is the measured data, the red curve is the calculated pattern and the gray plot is the difference between observed and calculated data. hkl ticks of  $LaB_6$  are shown in blue. The quantitative result is zero because no structural information is provided.

**Pawley Refinement of LaB\_6.** In a Pawley refinement, structure factors are not calculated from the structure but are refined as independent parameters for each reflection. They contribute to the correlation matrix. As a LeBail refinement, it can be used to determine structurally independent parameters and to extract intensities which can then be used for structure solution. In contrast to a LeBail refinement, correlation of the individual intensities can be used as additional information for structure solution. The input file is shown below and the final refinement plots are shown in Figure B1.34.

```
xdd "LaB6_Fixed.raw"
r_wp 10.966 r_exp 9.690 r_p 7.591 r_wp_dash 19.476 r_p_dash 19.284
r_exp_dash 17.208 weighted_Durbin_Watson 1.741 gof 1.132
do_errors
bkg @ 109.178088 ... ... 0.928684605
no_LIMIT_warnings
Specimen_Displacement(@, -0.04115)
```

This section contains the measurement file name (xdd LaB6\_Fixed.raw), agreement factors, refined Chebyshev background parameters (two shown), the command to calculate errors (do\_errors), a command to not display limit warnings (no\_LIMIT\_warnings) and refined specimen displacement in mm.

```
lam ymin_on_ymax 0.0001
Lam recs
 { 0.0159 1.534753 3.6854
   0.5691 1.540596 0.4370
   0.0762 1.541058 0.6000
   0.2517 1.544410 0.5200
   0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence(0.125)
axial conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
```

This part contains information to describe diffractometer settings and the wavelength for the fundamental parameter approach. The sample was measured with a fixed divergence slit with an opening of  $0.125^{\circ}$ .

```
hkl_Is
phase_name LaB6
Cubic(@ 4.15713`_0.00000)
LVol_FWHM_CS_G_L(1, 642.87869, 0.89, 696.16852, csgc, 962.01337, cslc, 2435.82068)
e0_from_Strain(0.00022, sgc1, 0.00010, slc1, 0.00086)
volume 71.8421992`_0.000166559246
space_group "Pm-3m"
load hkl_m_d_th2 I
{ 0      0      1      6      4.15714      21.35670 @ 11.2568643`_0.0345796072
      0      1      12      2.93954      30.38313 @ 37.2642101`_0.0896171236
      ... not all shown }
```

The 'hkl\_Is' keyword indicates that structural information for a Pawley refinement follows. Required information are the lattice parameters ('Cubic' macro where a = b = c and angles = 90°) and space group. The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\varepsilon_0 = \Delta d/d$ . As LaB<sub>6</sub> should have infinite crystallite size and no microstrain, the pattern does not significantly change when no microstructural model is applied. The 'volume' keyword calculates the unit cell volume, 'space\_group' indicates the space group. The 'load hkl\_m\_d\_th2 I' keyword is added automatically in the first refinement step and includes all reflections according to the space group and lattice parameters with their *hkl*, multiplicity, d-spacing, 2-Theta values and intensities with esd.

**Rietveld Refinement of LaB\_6.** In a Rietveld refinement, peak intensities (structure factors) are calculated from the structural content (atoms) of the unit cell. With this method, the unit cell weight is known and quantitative phase analysis of crystalline powder mixtures is possible. The input file is shown below and the final refinement plot is shown in Figure B1.35.

```
xdd "LaB6_Fixed.RAW"
r_wp 12.230 r_exp 9.693 r_p 8.773 r_wp_dash 21.624 r_p_dash 21.728
```

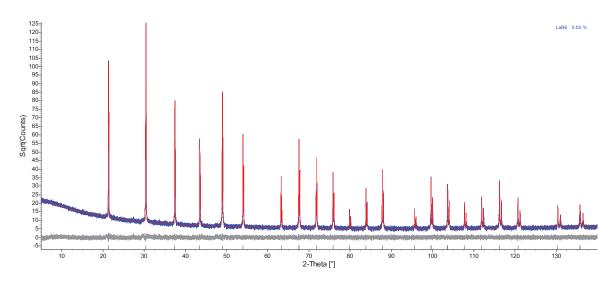


Fig. B1.34: Final refinement plot of the Pawley refinement of  $LaB_6$ . The blue curve is the measured data, the red curve is the calculated pattern and the gray plot is the difference between observed and calculated data. hkl ticks of  $LaB_6$  are shown in blue. The quantitative result is zero because no structural information is provided.

```
r_exp_dash 17.138 weighted_Durbin_Watson 1.400 gof 1.262
do_errors
bkg @ 105.067377 ... -13.0157227
no_LIMIT_warnings
Specimen_Displacement(@, -0.04502)
```

This section contains the measurement file name (xdd LaB6\_Fixed.raw), agreement factors, refined Chebyshev background parameters (two shown), the command to calculate errors (do\_errors), a command to not display limit warnings (no\_LIMIT\_warnings) and refined specimen displacement in mm.

```
lam ymin_on_ymax 0.0001
Lam_recs
 { 0.0159 1.534753 3.6854
  0.5691 1.540596
                    0.4370
  0.0762 1.541058
                    0.6000
  0.2517 1.544410
                    0.5200
  0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence(0.125)
axial_conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
```

str

This part contains information to describe diffractometer settings and the wavelength for the fundamental parameter approach. The sample was measured with a fixed divergence slit with an opening of  $0.125^{\circ}$ .

```
phase_name LaB6

Cubic(@ 4.15700`_0.00001)

LVol_FWHM_CS_G_L(1, 892.44090, 0.89, 1048.48902, csgc, 1768.25873, cslc, 2201.61914)

e0_from_Strain(0.00005, sgc1, 0.00010, slc1, 0.00012)

MVW(203.778, 71.8356248`_0.000299689962, 100.000`_0.000)

volume 71.8356248`_0.000299689962

space_group "Pm-3m"

scale @ 0.00133340363

site La1 x 0.0 y 0.0 z 0.0 occ La 1.0 beq @ 0.1444`_0.0068

site B1 x @ 0.19968`_0.00048 y 0.5 z 0.5 occ B 1.0 beq @ 0.0145`_0.0413
```

The 'str' keyword indicates that a structural description of a crystallographic unit cell follows. The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\varepsilon_0 = \Delta d/d$ . As LaB<sub>6</sub> is the NIST standard reference material 660a used for line profile calibration of a diffractometer. The pattern does not significantly change when no microstructural model is applied! Lattice parameters are defined by using a macro ('Cubic' macro where a = b = c and angles = 90°). The 'volume' keyword outputs the unit cell volume, 'space\_group' indicates the space group. The 'MVW' macro outputs unit cell mass, volume and weight percentage of the phase. The scale factor is used for linear scaling of the calculated pattern. It is also used for the calculation of the phase amount in quantitative analyses. The next two lines contain atomic site information. The site name follows the keyword 'site'. Next, x, y and z coordinates are given. The 'occ' keyword is followed by the atom type and its occupancy fraction. 'beq' is the keyword for isotropic displacement parameters.

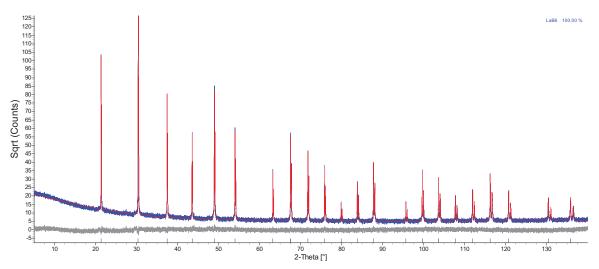


Fig. B1.35: Final refinement plot of the Rietveld refinement of  $LaB_6$ . The blue curve is the measured data, the red curve is the calculated pattern and the gray plot is the difference between observed and calculated data. *hkl* ticks of  $LaB_6$  are shown in blue. The quantitative result top right is Rietveld wt.-%, which is 100% because only one phase is present.

**Indexing of LaB\_6.** The following example demonstrates how to index a powder pattern. For obtaining good results, the pattern must have a reasonnable s/n ratio and peaks must be well resolved. It is best to use single phase samples. Lattice parameters found by indexing

can then be used with Pawley and LeBail refinements, as well as for phase identification and structure solution.

- 1. Open Topas Academic and load the powder pattern (previously converted to Bruker .raw format). Go to 'File' menu and select 'Load Scan Files'.
- Once the pattern is loaded, select 'Search Peaks' from the 'View' menu. Change peak width and noise threshold values until all peaks are found. You should see a screen as in Figure B1.36.
- 3. Open the tree view in the 'Parameter Window', select the list with the peak positions. Right-click and select copy all/selection (Figure B1.37).

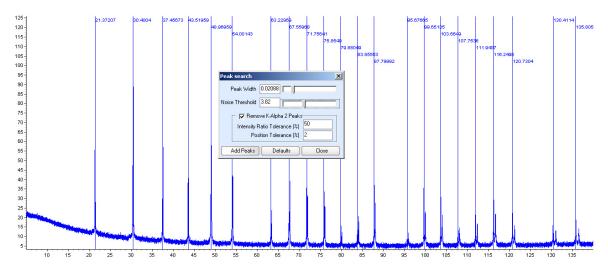


Fig. B1.36: Topas screen after peak search. The blue lines indicate identified peaks, 2-Theta angles are written next to the lines.

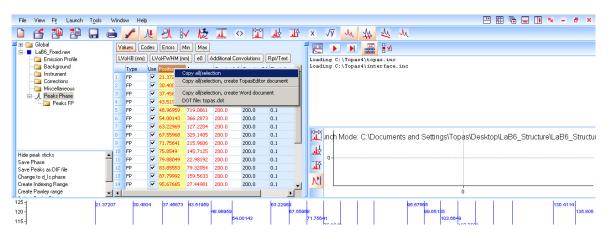


Fig. B1.37: Copy all peak positions from the 'Parameter Window'.

4. Open jEdit and create a new input file for Topas containing the measurement file name, diffractometer setup, wavelength information and background (see below). Add the keyword 'xo Is', followed by the 'CS L', 'CS G',

'Strain\_L' and 'Strain\_G' keywords in order to allow peak broadening. Next, paste the peak list copied in step 3. Refine all positions by adding an @ sign in front of each peak position. Add @ 1 after each peak. This value is the intensity, which will also be refined. Send the file to Topas and run. After refining peak positions and intensities, you should get a screen as in Figure B1.38.

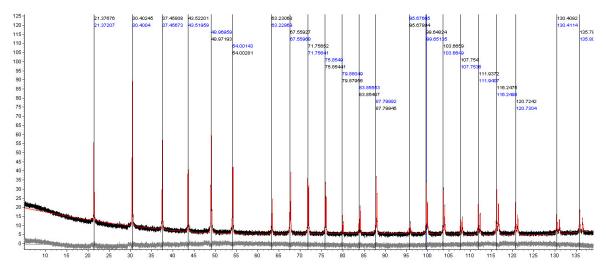


Fig. B1.38: The red curve is the pattern calculated from refined background, peak positions and peak intensities. Peak positions written in blue are positions from peak search, those in black after the refinement.

```
xdd LaB6_Fixed.raw
lam ymin_on_ymax 0.0001
Lam_recs
 { 0.0159
          1.534753
                     3.6854
   0.5691
           1.540596
                     0.4370
   0.0762
          1.541058
                     0.6000
   0.2517
          1.544410
                     0.5200
   0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence (0.125)
axial_conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
bkg @ 0 0 0 0 0
xo_Is
CS_L(0, 300)
CS_G(@, 300)
Strain_L(0, 0.05)
Strain_G(@, 0.05)
load xo I {
0 21.3767632 0 1
0 30.4024544 0 1
```

```
... not all shown }
```

5. Go back to jEdit and copy the list of refined peak positions and peak intensities from the updated file to the clipboard. Create a new input file with the keywords 'seed', 'index\_lam 1.540596', 'index\_zero\_error' and keywords for each crystal system that should be included in the indexing process. As we know that LaB<sub>6</sub> is a cubic phase, we only include the cubic crystal systems with the 'Cubic\_F', 'Cubic\_I', 'Cubic\_P' keywords. Add 'load index\_th2 index\_I' and paste the reflection positions and intensities copied from the old file. Make sure they are no longer refined by removing the @ sign. Compare with the input shown below, send to Topas and run.

6. When indexing has finished, go back to jEdit and open the new .ndx file (see below, only the first few lines are shown). Each row of the file contains a solution from the indexing process with line number, space group, status, number of unidentified reflections, cell volume, figure of merit, zero shift and lattice parameters (Coelho 2003). The first found solution with the highest figure of merit is already the correct one despite the obviously incorrect space group. This is because in powder diffraction, several space groups can not be distinguished and space group P23 has the same hkl conditions as space group Pm3m, which is the correct space group of LaB<sub>6</sub> (consult section 12.6 of the Topas Technical Reference). Select the contents of the .ndx file and copy all to the clipboard.

```
' Indexing method - Alan Coelho (2003), J. Appl. Cryst. 36, 86-95
' Time: 0.578 seconds
' Sg Status UNI Vol Gof Zero Lps...
Indexing_Solutions_With_Zero_Error_2 {
0) P23 2 0 71.854 2074.62 0.0203 4.1573 4.1573 4.1573 90.000 90.000 90.000 '== 22
1) F-43c 2 0 574.830 1675.65 0.0203 8.3147 8.3147 8.3147 90.000 90.000 90.000 '== 22
... not all shown
```

7. Go back to Topas and select 'Load INP File' from the 'File' menu. Load the input file used for indexing and select it in the 'Parameter Window'. At the botton of the 'Parameter Window', select 'Paste INP to Node/Selections' to paste all indexing solutions from the .ndx file. On the right side of the parameter window, click the 'Solutions' button to get a graphical display of cell volume vs. figure of merit (Figure B1.39).

**Quantification of known phases in a crystalline sample.** The following input file shows a quantification of two known phases CaF and Na<sub>2</sub>Al<sub>2</sub>Ca<sub>3</sub>F<sub>14</sub> in a crystalline sample. The keyword 'weight\_percent' outputs weight-% of each phase, normalized to 100%. Due to absorption of the Na<sub>2</sub>Al<sub>2</sub>Ca<sub>3</sub>F<sub>14</sub> phase, peaks are slightly asymmetric. This is corrected using a special keyword. Quantitative results are 99.468(5) wt.-% Na<sub>2</sub>Al<sub>2</sub>Ca<sub>3</sub>F<sub>14</sub> and 0.532(5) wt.-% CaF. The final refinement plot is shown in Figure B1.40.

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Miscellaneous	2 🔲 F-43	3c 2 0 574.830 1675.65		Time 0.2 Gof 2074.62 Vol 72 Best 2074.62	7
Display	3 🗌 F23	2 0 574.830 1281.3		Time 0.2 Gof 2074.62 Vol 72 Best 2074.62 Time 0.2 Gof 2074.62 Vol 72 Best 2074.62	9
imp2.inp	4 🔲 Ia-3	2 0 574.830 1013.19		Time 0.2 Gof 2074.62 Vol 72 Best 2074.62	10
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	6 🗌 Pn-3	3n 2 0 574.830 680.73	0.0203 8.3147	irying to determine space group	
4 Þ				File C:\Documents and Settings\Topas\Desktop\LaB6_St	ructure\inp2.ndx written to
Gof vs Volume Details for selected					•
	2000	Gof vs Volume		2000 Launch Mode: C:\Documents and Settings	pas\Desktop\LaB6 Structure\inp2.in
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Select next	1'000-2	<b>}──┤ġ───┼</b> ───┼─		1'000	
Sort by Gof-decreasing	500			500	
Sort by UNI/Gof-decreasing Sort by Volume		3]\ 🕴			
		and a series and a series of the series of t	mo apro A .		800 1'000 1'200
H	<b>V</b>	500 1'000 1'500 2'000	2'500 3'000		
600 001 011	111 00:	2 201 02	2 003	004 401	422 40:
580		211	301	401	
560 -			311	331	
260				222 40	
540-				302 231	421
520 -	1				

Fig. B1.39: Display of indexing solutions. Red lines are peaks belonging to the selected space group.

```
xdd "NAC.RAW"
r_wp 6.071 r_exp 3.920 r_p 4.577 r_wp_dash 10.995 r_p_dash 10.971
r_exp_dash 7.099 weighted_Durbin_Watson 0.895 gof 1.549
do_errors
bkg @ 426.075391 ... ... -2.45058623
no_LIMIT_warnings
Specimen_Displacement(@, -0.07846)
```

This section contains the measurement file name (xdd NAC.raw), agreement factors, refined Chebyshev background parameters (two shown), the command to calculate errors (do\_errors), a command to not display limit warnings (no\_LIMIT\_warnings) and refined specimen displacement in mm.

```
lam ymin_on_ymax 0.0001
Lam_recs
 { 0.0159 1.534753 3.6854
  0.5691 1.540596 0.4370
  0.0762 1.541058 0.6000
  0.2517 1.544410 0.5200
  0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Variable_Divergence_Shape(10)
Variable_Divergence_Intensity
axial_conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
```

This part contains information to describe diffractometer settings and the wavelength for the fundamental parameter approach. The sample was measured with an automatic divergence slit irradiating 10 mm of the sample surface. The

keyword 'Variable\_Divergence\_Intensity' corrects measured intensities.

```
str
phase_name NAC
Cubic(@ 10.25089'_0.00000)
LVol_FWHM_CS_G_L(1, 445.03952, 0.89, 491.92040, csgc, 710.14683, cslc, 1474.39087)
e0_from_Strain(0.00004, sgc1, 0.00010, slc1, 0.00010)
MVW(1944.625, 1077.17204`_0.00133777644, 99.468`_0.005)
weight_percent 99.468 \_0.005
space_group "I213"
scale @ 0.000721118589
site Cal x 0.46618 y 0.00000 z =1/4;
                                      occ Ca 1.0 beq 0.5771
site Al1 x 0.24754 y 0.24754 z 0.24754 occ Al 1.0 beq 0.5170
site Na1 x 0.08492 y 0.08492 z 0.08492 occ Na 1.0 beg 1.9929
site F1 x 0.13817 y 0.30571 z 0.12093 occ F 1.0 beq 0.7395
site F2 x 0.36357 y 0.36250 z 0.18886 occ F 1.0 beq 1.0556
site F3 x 0.46098 y 0.46098 z 0.46098 occ F 1.0 beq 0.4369
Absorption_With_Sample_Thickness_mm_Shape_Intensity(0, 69.10407, 0, 3.11216)
```

This part contains structural information of the Na<sub>2</sub>Al<sub>2</sub>Ca<sub>3</sub>F<sub>14</sub> phase. The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\varepsilon_0 = \Delta d/d$ . The 'MVW' macro outputs unit cell mass, volume and wt.-% (also output by the 'weight\_percent' keyword). Then follow space group, scale factor and atomic coordinates, occupancies and isotropic  $B_{eq}$ . The last macro corrects peak asymmetry due to sample absorption.

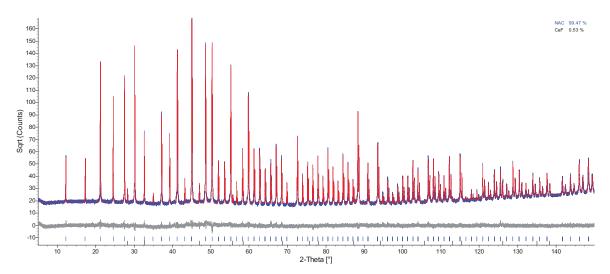


Fig. B1.40: Final refinement plot of the phase quantification. The blue curve is the measured data, the red curve is the calculated pattern and the gray plot is the difference between observed and calculated data. hkl ticks of Na<sub>2</sub>Al<sub>2</sub>Ca<sub>3</sub>F<sub>14</sub> (blue) and CaF (black) are shown. The quantitative results top right are Rietveld wt.-%.

str
phase\_name CaF
Cubic(@ 5.46555`\_0.00003)
LVol\_FWHM\_CS\_G\_L(1, 445.03952, 0.89, 491.92040, csgc, 710.14683, cslc, 1474.39087)
e0\_from\_Strain(0.00004, sgcl, 0.00010, slcl, 0.00010)

```
MVW(312.300, 163.268453`_0.00229605502, 0.532`_0.005)
weight_percent 0.532`_0.005
space_group "Fm-3m"
scale @ 0.000158590228
site Cal x 0.00 y 0.00 z 0.00 occ Ca 1.0 beq 0.5
site F1 x 0.25 y 0.25 z 0.25 occ F 1.0 beq 0.5
```

This part contains structural information of the CaF phase. The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\varepsilon_0 = \Delta d/d$ . The 'MVW' macro outputs unit cell mass, volume and wt.-% (also output by the 'weight\_percent' keyword). Then follow space group, scale factor and atomic coordinates, occupancies and isotropic  $B_{eq}$ . Due to the low amount of the CaF phase,  $B_{eq}$  of CaF can not be refined.

**Quantification of amorphous content.** The following input file describes how to quantify amorphous content. The sample was prepared from a mixture of 1g quartz (supposed as unknown), 0.5g silicon (supposed as admixed internal standard) and 0.5g glass (supposed as unknown amorphous phase). It is very important that the spiked phase has similar density and absorption properties as the other phases in the sample. The example shown below does not work when e.g. LiF is used as spike phase!

A simple quantification of quartz and silicon is performed and the Rietveld wt.-% are returned. However, they represent 100% of the crystalline phases of the mixture but not 100% of the total mixture. Ideally, Rietveld results should be 66.6% wt.-% quartz and 33.3 wt.-% silicon. It is known that 25 wt.-% silicon were added, therefore, 33.3 wt.-% silicon (from Rietveld refinements) represent 25 wt.-% of the total mixture. Normalization of the 33.3 wt.-% silicon to the 25 wt.-% silicon as admixed internal standard results in a corrected wt.-percentage of 50 wt.-% quartz for the measured 66.6 wt.-% quartz (from Rietveld). The difference to 100% represents the amorphous content of the sample: 100 wt.-% total – 25 wt.-% silicon – 50 wt.-% quartz = 25 wt.-% amorphous content (glass). This is done using the keyword 'spiked\_phase\_measured\_weight\_percent 25' in the structure description of the phase used as internal standard and 'corrected\_weight\_percent' in all other structures. The amorphous phase fraction can be output with 'weight\_percent\_amorphous'. The final refinement plot is shown in Figure B1.41.

```
xdd "Mixture.raw"
r_wp 10.747 r_exp 7.651 r_p 8.215 r_wp_dash 18.996 r_p_dash 19.772 r_exp_dash 13.523
weighted_Durbin_Watson 1.150 gof 1.405
do_errors
bkg @ 134.411527 ... ... -7.08022216
no_LIMIT_warnings
Specimen_Displacement(@, -0.03148)
```

This section contains the measurement file name (xdd Mixture.raw), agreement factors, ten refined Chebyshev background parameters (two shown), the command to calculate errors (do\_errors), a command to not display limit warnings (no\_LIMIT\_warnings) and refined specimen displacement in mm.

```
lam ymin_on_ymax 0.0001
Lam_recs
 { 0.0159 1.534753 3.6854
  0.5691 1.540596 0.4370
  0.0762 1.541058 0.6000
  0.2517 1.544410 0.5200
  0.0871 1.544721 0.6200 }
LP_Factor( 0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence (0.125)
axial conv
filament_length 12
sample_length 10
receiving_slit_length 15
primary_soller_angle 2.55
secondary_soller_angle 2.55
```

This part contains information to describe diffractometer settings and the wavelength for the fundamental parameter approach. This sample was measured with a fixed divergence slit with  $0.125^{\circ}$  opening.

```
str
phase_name "Silicon"
Cubic(@ 5.43258)
MVW (224.683, 160.331254`_0.0021742123, 31.532`_0.082)
scale @ 0.00111033118
space_group "Fd-3m:2"
LVol_FWHM_CS_G_L(1, 383.35715, 0.89, 534.46669, csgc1, 10000.00000, cslc1, 601.91460)
e0_from_Strain(0.02212, sgc1, 0.00010, slc1, 0.08849)
site Si1 x 0.125 y 0.125 z 0.125 occ Si 1.0 beq 0.85
phase_MAC 63.9439604
Phase_LAC_1_on_cm(148.79935)
Phase_Density_g_on_cm3(2.32703)
weight_percent 31.532`_0.082
spiked_phase_measured_weight_percent 25
corrected_weight_percent 25`_0
```

Structural information of the spiked silicon phase. The 'MVW' macro calculates unit cell mass, volume and weight.-%, followed by the scale factor and space group  $(2^{nd}$  setting of space group  $Fd\bar{3}m!$ ). The 'LVol\_FWHM\_CS\_G\_L' macro calculates the size of coherently scattering domains ('crystallite size'), whereas the macro 'e0\_from\_Strain' calculates  $\epsilon_0 = \Delta d/d$ . Next follow atomic coordinates, occupancies and isotropic  $B_{eq}$ . The 'phase\_MAC' keyword calculates the phase mass absorption coefficient, 'Phase\_LAC\_1\_on\_cm' calculates the linear absorption coefficient, 'Phase\_Density\_g\_on\_cm3' calculates the phase density. 'Weight\_percent' calculates Rietveld wt.-%. The 'spiked\_phase\_measured\_weight\_percent' corrects this value. The 'corrected\_weight\_percent' corrects Rietveld wt.-%

str
phase\_name "Quartz"
a @ 4.91496`\_0.00003
b = Get(a);
c @ 5.40662`\_0.00004
al 90

```
be 90
ga 120
MVW(180.252, 113.108808'_0.00145825088, 68.468'_0.082)
scale @ 0.00425994962
space_group "P3221"
LVol_FWHM_CS_G_L(1, 398.51387, 0.89, 401.71678, csgc2, 486.52138, cslc2, 3352.67231)
e0_from_Strain(0.01835, sgc2, 0.04814, slc2, 0.04032)
site Sil x 0.46832 y 0.0
                            z =2/3;
                                      occ Si 1.0 beq 0.82
site 01 x 0.41353 y 0.27029 z 0.78974 occ 0 1.0 beq 1.35
phase_MAC 36.0110617
Phase_LAC_1_on_cm(95.29491)
Phase_Density_g_on_cm3(2.64627)
weight_percent 68.468'_0.082
corrected_weight_percent 54.2851069'_0.0647018326
```

This part is the same as above but with structural information for quartz. 'Weight\_percent' calculates Rietveld wt.-%, which are corrected by the spike above. Corrected wt.-% are output with the 'corrected\_weight\_percent' keyword.

weight\_percent\_amorphous 20.7148931 \\_0.205195603

Weight\_percent\_amorphous is the difference of 100% - sum of corrected wt.-%.

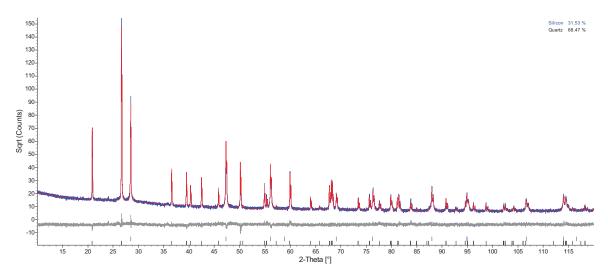


Fig. B1.41: Final refinement plot of the amorphous content quantification. hkl ticks of silicon (blue) and quartz (black) are shown. The blue curve is the measured data, the red curve is the calculated pattern and the gray plot is the difference between observed and calculated data. Note the very slight background bump between 20° and 35° 2-Theta which may result from the added glass phase. The quantitative results top right are Rietveld wt.-% without the spike correction.

#### Refinement of temperature-dependent diffraction patterns

Multipattern refinement in Topas with uncorrected diffraction patterns and no height correction. Several measurements of  $Al_2O_3$  were performed at different temperatures in order to refine temperature-dependent lattice parameters and crystal structure of corundum. Each pattern is refined independently of the others. Measured patterns were not corrected for height displacement, thus specimen displacement must follow the values from the table describing sample holder expansion of the Anton-Paar HTK 1200 furnace (Table B1.8). This method is the easiest to use when refining temperature-dependent diffraction patterns. However, no physical model expressing temperature-dependence of lattice parameters is considered. Further, when using diffraction patterns that are not corrected for height displacements, the refinement can be unstable because of unappropriate start values of specimen displacement towards higher temperatures.

Table B1.8: Sample height displacement of the HTK 1200 (°C, mm).

Temp.	Displ.	Temp.	Displ.	Temp.	Displ.	Temp.	Displ.
100	0.030	400	0.120	700	0.250	1000	0.390
200	0.060	500	0.170	800	0.300	1100	0.450
300	0.090	600	0.210	900	0.350	1200	0.510

r\_exp 11.810 r\_exp\_dash 20.004 r\_wp 15.268
r\_wp\_dash 25.862 r\_p 11.450 r\_p\_dash 23.840
weighted\_Durbin\_Watson 17.399 gof 1.293
do\_errors
no\_LIMIT\_warnings

This section contains agreement factors, the keyword to calculate errors (do\_errors) and a keyword to not display limit warnings (no\_LIMIT\_warnings).

```
macro diffractometer
{ lam ymin_on_ymax 0.0001
 Lam_recs
  { 0.0159 1.534753 3.6854
   0.5691 1.540596 0.4370
    0.0762 1.541058 0.6000
    0.2517 1.544410 0.5200
    0.0871 1.544721 0.6200 }
LP Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
Divergence (0.25)
axial_conv
filament_length 12
sample_length 10
receiving_slit_length 15
```

```
primary_soller_angle 2.55
secondary_soller_angle 2.55
start_X 10
finish_X 140 }
```

The diffractometer setup and wavelengths for the fundamental parameter approach are defined. Instead of repeating this part in each temperature-dependent refinement below, it is defined as a macro namend 'diffractometer' which is inserted below.

```
xdd Al2O3_0025.raw
diffractometer
bkg @ 46.7328227 ... ... 0.863553211
```

This part contains the datafile name (xdd Al2O3\_0025.raw). The diffractometer setup is input using the macro 'diffractometer' defined above. Next, the Chebyshev background parameters are refined (two shown).

This part contains all structural information of the  $Al_2O_3$  phase for the 25°C refinement. Note that parameters of interest are given a parameter name which is different for each temperature. Both displacement parameters are refined together using the parameter name biso\_T25. The last two lines are macros for peak broadening due to microstructural effects.

An output file is generated (Al2O3\_LPs\_Output.TXT) and values of some of the parameters above are written to the file. Out(a\_al2o3\_T25, "%11.5f", "%11.5f") means that the value of the parameter a\_al2o3\_T25 is written to the file as floating number with five digits, as is the value of the esd. The first two lines are a file header.

Similar inputs for the  $100^{\circ}$ C to  $1200^{\circ}$ C patterns follow here. Temperature dependent parameters have to be changed to the corresponding temperature in each refinement section (e.g. a\_al2o3\_T25 becomes a\_al2o3\_T100 in the  $100^{\circ}$ C refinement. Below, the last refinement ( $1200^{\circ}$ C) is shown.

```
xdd Al203_1200.raw
diffractometer
bkg @ 45.1037186 ... ... ... ... 0.563962219
str
phase_name Al2O3_T1200
r_bragg 3.9436839
Trigonal(a_al2o3_T1200 4.80746, c_al2o3_T1200 13.13868)
cell_volume v_al2o3_T1200 262.974101
Specimen_Displacement (displ_T1200, -0.54138)
scale @ 0.000642774818
space_group "R-3c"
                   y 0 z @ 0.35223 occ Al 1.0 beq biso_T1200 0.9284
site All x O
site O1 x @ 0.69378 y 0 z = 1/4; occ O 1.0 beq biso_T1200 0.9284
LVol_FWHM_CS_G_L(1, 77.91640, 0.89, 91.39145, csqc_T1200, 153.4010, cslc_T1200, 193.4558)
e0_from_Strain( 0.02323, sqc_T1200, 0.07527, slc_T1200, 0.03029)
out "Al2O3_LPs_Output.TXT" append
Out(=Get(phase_name), " \n ?%s?")
Out(a_al2o3_T1200, " %11.5f", " %11.5f")
Out(c_al2o3_T1200, " %11.5f", " %11.5f")
Out(v_al2o3_T1200, " %11.3f", " %11.3f")
Out(Get(r_wp), " %11.2f")
Out(Get(gof), " %11.2f")
Out(displ_T1200, "%11.3f", " %11.3f")
```

Multipattern refinement in Topas with uncorrected input files and known height correction. Instead of using a specimen displacement of zero as start value for each temperaturedependent refinement, patterns can be corrected with the known height displacement as listed in the Anton-Paar table (Table B1.8). The same input file as above is shown with the known displacements implemented. At each temperature, specimen displacement is calculated as: displacement<sub>total</sub> = displacement<sub>refined</sub> – known value from the table. With this method, appropriate starting values for specimen displacement are already provided at the beginning of the refinement.

```
r_wp 15.268 r_exp 11.810 r_p 11.450
r_wp_dash 25.862 r_p_dash 23.840 r_exp_dash 20.004
weighted_Durbin_Watson 17.399 gof 1.293
do_errors
no_LIMIT_warnings
```

This section contains agreement factors, the keyword to calculate errors (do\_errors) and a keyword to not display limit warnings (no\_LIMIT\_warnings).

```
macro diffractometer
{ lam ymin_on_ymax 0.0001
  Lam_recs
   { 0.0159 1.534753 3.6854
     0.5691 1.540596 0.4370
```

0.0762 1.541058 0.6000 0.2517 1.544410 0.5200 0.0871 1.544721 0.6200 } LP\_Factor(0) Rp 240 Rs 240 Slit\_Width(0.07) Divergence (0.25) axial\_conv filament\_length 12 sample\_length 10 receiving\_slit\_length 15 primary\_soller\_angle 2.55 secondary\_soller\_angle 2.55 start\_X 10 finish\_X 140 }

The diffractometer setup and wavelengths for the fundamental parameter approach are defined. Instead of repeating this part for each temperature-dependent refinement below, it is defined as a macro namend 'diffractometer' which is then inserted below.

```
xdd Al2O3_0025.raw
diffractometer
bkg @ 46.7328227 ... ... ... ... 0.863553211
```

This part contains the datafile name (xdd Al2O3\_0025.raw), then the diffractometer setup is input using the macro 'diffractometer' defined above. Next, the Chebyshev background parameters are refined (two shown).

This part contains all structural information of the  $Al_2O_3$  phase for the 25°C refinement. Note that parameters of interest are given a parameter name which is different for each temperature. Both displacement parameters are refined together using the parameter name biso\_T25. The last two lines are macros for peak broadening due to microstructural effects. A new parameter 'height\_error\_T25' is defined. The actual specimen displacement is then displ\_T25 = height\_error\_T25 - 0.0. Nothing is subtracted because the sample holder expansion at 25°C is expected to be zero. If the Anton-Paar table (Table B1.8) is exact, the value of height\_error\_Txy should be the same at each temperature.

```
out "Al2O3_LPs_Output.TXT" append
Out_String("\n Name a-axis esd c-axis esd volume esd Rwp GoF Ht_uncorr esd Ht_corr esd")
```

An output file is generated (Al2O3\_LPs\_Output.TXT) and values of some of the parameters above are written to the file. Out(a\_al2o3\_T25, "%11.5f", "%11.5f") means that the value of the parameter a\_al2o3\_T25 is written to the file as floating number with five digits, as is the value of the esd. The first two lines are a file header.

Similar inputs for the  $100^{\circ}$ C to  $1200^{\circ}$ C patterns follow here. Temperature dependent parameters have to be changed to the corresponding temperature in each refinement section (e.g. a\_al2o3\_T25 becomes a\_al2o3\_T100 in the  $100^{\circ}$ C refinement. Below, the last refinement ( $1200^{\circ}$ C) is shown.

```
xdd Al203_1200.raw
diffractometer
bkg @ 46.7552041 ... ... ... ... 1.06807956
str
phase name Al2O3 T1200
r_bragg 3.9436839
Trigonal(a_al2o3_T1200 4.80746, c_al2o3_T1200 13.13868)
cell_volume v_al2o3_T1200 262.974101
prm height_error_T1200 -0.03132
Specimen_Displacement (displ_T1200, =height_error_T1200 - 0.51;: -0.54132)
scale @ 0.000642774818
space_group "R-3c"
site All x O
                   y 0 z @ 0.35223 occ Al 1.0 beq biso_T1200 0.9284
site O1 x @ 0.69378 y 0 z = 1/4; occ 0 1.0 beg biso_T1200 0.9284
LVol_FWHM_CS_G_L(1, 77.91640, 0.89, 91.39145, csqc_T1200, 153.4010, cslc_T1200, 193.4558)
e0_from_Strain( 0.02323', sqc_T1200, 0.07527, slc_T1200, 0.03029)
out "Al2O3_LPs_Output.TXT" append
Out(=Get(phase_name), " \n ?%s?")
Out(a_al2o3_T1200, " %11.5f", " %11.5f")
Out(c_al2o3_T1200, " %11.5f", " %11.5f")
Out(v_al2o3_T1200, " %11.3f", " %11.3f")
Out(Get(r_wp), " %11.2f")
Out(Get(gof), " %11.2f")
Out(displ_T1200, "%11.3f", " %11.3f")
```

Out (height\_error\_T1200, "%11.3f", " %11.3f")

Specimen displacement used to correct the pattern here is displ\_T1200 = height\_error\_T1200 - 0.51. The value of 0.51 mm is the height displacement as listed in the Anton-Paar table (Table B1.8). The value of the 'height\_error\_T1200' parameter is the difference between actual specimen displacement and the value listed in the table. If the Anton-Paar table is exact, the value of height\_error\_Txy should be the same at each temperature. Note that height\_error\_T25 is 0.02462 whereas height\_error\_T1200 is -0.03132.

**Parameterized refinement of temperature-dependent diffraction patterns.** The idea of the parameterization shown here is based on Stinton and Evans (2007).

Traditionally, refining temperature-dependent parameters is done by refining all parameters at each temperature individually, as shown with the two examples above. However, with this method, no physical model is considered, i.e. it is not taken in account that temperaturedependency of e.g. lattice parameters is in fact a function of temperature. This can be expressed with a mathematical function.

Instead of refining each pattern individually, it is possible to refine them simultaneously and express parameters of interest (lattice parameters, specimen displacement) as a temperature-dependent function (Figure B1.42). This eliminates outliers (which are common when treating high-temperature data individually) and it drastically reduces the number of refined parameters: Refining lattice parameters of an orthorhombic phase (a, b, c), corundum as internal standard (a, c) and specimen displacement for 21 diffraction patterns (25°C, 50°C and then each 50°C up to 1000°C) results in a total of 126 refined parameters. If lattice parameters and specimen displacement are expressed as functions of temperature, the only variable in a pattern is temperature, which is supposed to be known. Polynomial coefficients are the same for each temperature and can therefore be refined simultaneously over all measurements. This method results in a total of 18 parameters for 21 diffraction patterns. Further, the physical property implicating that lattice parameter evolution is temperature-dependent is now considered in the refinement. Prior to setting up the refinement, verify if the chosen function is accurate!

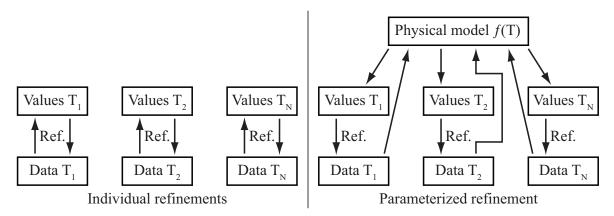


Fig. B1.42: Parameterized refinements allow consideration of a physical model (e.g. lattice parameter evolution as a function of temperature, expressed by a second order polynomial).

Results can be monitored from refined lattice parameters of an internal standard in the sample (in this case corundum) and by comparing specimen displacement values to those in the Anton-Paar table (Table B1.8).

A drawback of this method is error propagation (i.e. errors become larger at the beginning and

the end of the temperature range) when using the traditional method of error determination from a correlation matrix. This can be overcome by using the bootstrap method implemented in Topas: A series of refinements is performed, each with a fraction of observed data modified to obtain a new bootstrap sample (equals one cycle). After a given number of cycles, standard deviations of the refined values become the bootstrap errors (Coelho 2007b).

In the sample file shown below, the physical model expressing lattice parameters a and c of corundum and a, b and c of  $Al_5BO_9$  and specimen displacement (h) as a function of temperature have been defined as shown in Equations B.10 to B.15.

$$a_al2o3(T) = a0_al2o3 + ax1_al2o3 \cdot T + ax2_al2o3 \cdot T^2$$
 (B.10)

$$c_{al2o3}(T) = c_{al2o3} + c$$

$$a_al5bo9(T) = a0_al5bo9 + ax1_al5bo9 \cdot T + ax2_al5bo9 \cdot T^2$$
(B.12)

$$b_al5bo9(T) = b0_al5bo9 + bx1_al5bo9 \cdot T + bx2_al5bo9 \cdot T^2$$
 (B.13)

$$al5bo9(T) = c0 al5bo9 + cx1 al5bo9 \cdot T + cx2 al5bo9 \cdot T^2$$
(B.14)

$$h(T) = height_0 + height_x 1 \cdot T + height_x 2 \cdot T^2$$
(B.15)

~

At each temperature, lattice parameters and sample displacement are expressed with Equations B.10 to B.15. The only value that is changed each pattern is temperature T as a fixed parameter. The final refinement plot is shown in Figure B1.43.

```
r_exp 6.714 r_exp_dash 10.117 r_wp 7.912
r_wp_dash 11.923 r_p 6.013 r_p_dash 10.736
weighted_Durbin_Watson 37.137 gof 1.178
bootstrap_errors 200
no_LIMIT_warnings
```

 $\mathbf{c}$ 

This section contains agreement factors, the command to calculate bootstrap errors with 200 cycles (bootstrap\_errors 200) and a command to not display limit warnings (no\_LIMIT\_warnings).

```
macro diffractometer
{ lam ymin_on_ymax 0.0001
   Lam_recs
   { 0.0159 1.534753 3.6854
      0.5691 1.540596 0.4370
      0.0762 1.541058 0.6000
      0.2517 1.544410 0.5200
      0.0871 1.544721 0.6200 }
LP_Factor(0)
Rp 240
Rs 240
Slit_Width(0.07)
```

Divergence (0.125)
axial\_conv
filament\_length 12
sample\_length 10
receiving\_slit\_length 15
primary\_soller\_angle 2.55
secondary\_soller\_angle 2.55
start\_X 10
finish\_X 75.3 }

The diffractometer setup and the used wavelengths for the fundamental parameter approach are defined. Instead of repeating this part for each temperature-dependent refinement below, it is written as a macro namend 'diffractometer' which is then inserted below.

```
/ _____
'Al203 lattice parameter polynomial coefficients.
'a-axis(T) = a0_al2o3 + ax1_al2o3 *T + ax2_al2o3 *T^2.
/ _____
          4.75792`_0.00021
prm a0_al2o3
prm ax1_al2o3 4590.88876`_84.77322 ' * 100'000'000 too large!
prm ax2_al2o3 -0.48723`_0.07161
                             ' * 100'000'000 too large!
prm c0_al2o3
            13.01946`_0.00067
prm cx1_al2o3 4318.12403`_294.20815 ' * 100'000'000 too large!
prm cx2_al2o3 5.21644`_0.27224 ' * 100'000'000 too large!
ı_____
'Al5B09 lattice parameter polynomial coefficients.
'a-axis(T) = a0_al5bo9 + ax1_al5bo9 *T + ax2_al5bo9 *T^2.
/ _____
prm a0_al5bo9
             5.66852`_0.00005
prm ax1_al5bo9 1754.40050`_25.82802 ' * 100'000'000 too large!
prm ax2_al5bo9 0.61439`_0.02564 ' * 100'000'000 too large!
prm b0_al5bo9 1 5.01305`_0.00020
prm bx1_al5bo9 9300.74475`_62.96127 ' * 100'000'000 too large!
prm bx2_al5bo9 0.22990`_0.05751 ' * 100'000'000 too large!
             7.69477`_0.00006
prm c0 al5bo9
prm cx1_al5bo9 3804.41837 `_27.77816 ' * 100'000'000 too large!
prm cx2_al5bo9 0.36876`_0.02333 ' * 100'000'000 too large!
/ _____
'Height displacement polynomial coefficients.
'h(T) = height_0 + height_x1 *T +height_x2 *T^2.
/ _____
prm height_0 0.01420`_0.00035 ' * 10'000 too large!
prm height_x1 -3.07394`_0.01445 ' * 10'000 too large!
prm height_x2 -0.00144`_0.00001 ' * 10'000 too large!
```

This section contains the second order polynomial coefficients expressing  $Al_5BO_9$  and corundum lattice parameters and specimen displacement of the heating chamber as a function of temperature. Because Topas can handle but does not output very small numbers, coefficients were divided by a factor, which means that the numbers above are 100'000'000 times and 10'000 times too large, respectively.

```
xdd 0025.raw
diffractometer
prm !T25 = 25;
bkg @ 132.785828`_0.41725691 ... ... 2.99189097`_0.25181182
```

This part contains the data file name (xdd 0025.raw), information about the diffractometer setup is input using the macro 'diffractometer' defined above. A fixed parameter T25 defines temperature with a value of 25. Chebyshev back-ground parameters are refined (two shown).

```
hkl_Is
lebail 1
phase name Al2O3 T25
prm !a_al2o3_T25 = a0_al2o3 + ax1_al2o3/10^8 (T25) + ax2_al2o3/10^8 (T25)^2;
prm !c_al2o3_T25 = c0_al2o3 + cx1_al2o3/10^8 (T25) + cx2_al2o3/10^8 (T25)^2;
Trigonal (=a_al2o3_T25;: 4.75906`_0.00019, =c_al2o3_T25;: 13.02057`_0.00060)
prm height_T25 = height_0 + height_x1/10000 (T25) + height_x2/10000 (T25)^2;
Specimen_Displacement (!displ_T25, =height_T25;: 0.00643`_0.00032)
CS_L(@, 10000.00000`_3176.77964)
CS_G(@, 172.58488`_21.82458)
Strain_L(@, 0.12754 \_0.01169)
Strain_G(@, 0.00010`_0.03040)
space_group "R-3c"
load hkl_m_d_th2 I
                3.48231
{ 0 }
    1 2 6
                           25.55942 0.620856079
  ... not all shown }
```

A LeBail refinement (hkl\_Is and lebail 1 keywords) of corundum is set up. Lattice parameters a\_al2o3\_T25 and c\_al2o3\_T25 are calculated from the polynomials. Calculated lattice parameters are then used to describe the corundum unit cell (line starting with 'Trigonal...'). The values are only input but not refined, as corundum lattice parameters are fixed parameters. The only refinable parameters are polynomial coefficients listed further above! Specimen displacement is calculated from the polynomial and the fixed value height\_T25 is used as specimen displacement for this temperature. As above, polynomial parameters are the only refineable parameters.

```
hkl_Is
lebail 1
phase_name Al5BO9_T25
prm !a_al5bo9_T25 = a0_al5bo9 + ax1_al5bo9/100000000 (T25) + ax2_al5bo9/100000000 (T25)^2;
prm !b_al5bo9_T25 = b0_al5bo9 + bx1_al5bo9/100000000 (T25) + bx2_al5bo9/100000000 (T25)^2;
prm !c_al5bo9_T25 = c0_al5bo9 + cx1_al5bo9/100000000 (T25) + cx2_al5bo9/100000000 (T25)^2;
a =a_al5bo9_T25;: 5.66896`_0.00004
b =b_al5bo9_T25;: 15.01537 \_0.00020
c =c_al5bo9_T25;: 7.69572`_0.00006
cell_volume v_al5bo9_T25 655.071844`_0.0108051685
Specimen_Displacement (!displ_T25, =height_T25;: 0.00643`_0.00032)
CS_L(@, 547.78105`_41.65186)
CS_G(@, 352.93544`_992.96429)
Strain_L(0, 0.11538`_0.00807)
Strain_G(@, 0.11127`_0.00777)
space_group "Cmc21"
load hkl_m_d_th2 I
{ 0
    2 0 2 7.50768
                           11.77797 0.0931799498
  ... not all shown }
```

Parameterization of  $Al_5BO_9$  lattice parameters is identical to the one above, however, specimen displacement for this phase is not calculated again but the value is taken from the corundum part.

out "results.txt" append

An output file is saved (results.txt) and values of all parameters defined above are written to the file.  $Out(a_al2o3_T25, \%11.5f'', \%11.5f'')$  means that the value of the parameter  $a_al2o3_T25$  is written to the file as floating number with five digits. When using the bootstrap error model, esds are not written to the output file. The first two lines define the file header.

Similar inputs for the 100°C to 1200°C patterns follow here. Temperature dependent parameters have to be changed to the corresponding temperature in each refinement section (e.g. a\_al2o3\_T25 becomes a\_al2o3\_T100 in the 100°C refinement and prm !T25 = 25; becomes prm !T100 = 100; in the 100°C refinement). Below, the last refinement (1000°C) is shown.

```
xdd 1000.raw
diffractometer
prm !T1000 = 1000;
bkg @ 114.453609`_0.360956049 ... ... 4.42960845`_0.26267103
hkl_Is
lebail 1
phase_name Al2O3_T1000
prm !a_al2o3_T1000 = a0_al2o3 + ax1_al2o3/10^8 (T1000) + ax2_al2o3/10^8 (T1000)^2;
prm !c_al2o3_T1000 = c0_al2o3 + cx1_al2o3/10^8 (T1000) + cx2_al2o3/10^8 (T1000)^2;
Trigonal (=a_al2o3_T1000;: 4.79895`_0.00018, =c_al2o3_T1000;: 13.11481`_0.00100)
prm height_T1000 = height_0 + height_x1/10000 (T1000) + height_x2/10000 (T1000)^2;
Specimen_Displacement (!displ_T1000, =height_T1000;: -0.43681`_0.00040)
CS_L(@, 980.06673 \_2824.56239)
CS_G(@, 378.67059`_2320.60807)
Strain_L(@, 0.06310`_0.02097)
Strain_G(@, 0.00010`_0.03512)
space_group "R-3c"
load hkl_m_d_th2 I
{ 0 1 2 6 3.51036 25.35179 0.590999731
  ... not all shown }
hkl_Is
lebail 1
phase_name A15B09_T1000
prm !a_al5bo9_T1000 = a0_al5bo9 + ax1_al5bo9/10^8 (T1000) + ax2_al5bo9/10^8 (T1000)^2;
prm !b_al5bo9_T1000 = b0_al5bo9 + bx1_al5bo9/10^8 (T1000) + bx2_al5bo9/10^8 (T1000)^2;
prm !c_al5bo9_T1000 = c0_al5bo9 + cx1_al5bo9/10^8 (T1000) + cx2_al5bo9/10^8 (T1000)^2;
a =a_al5bo9_T1000;: 5.69221`_0.00004
b =b_al5bo9_T1000;: 15.10835`_0.00023
c =c_al5bo9_T1000;: 7.73650`_0.00008
cell_volume v_al5bo9_T1000 665.337929'_0.0183982884
```

```
Specimen_Displacement (!displ_T1000, =height_T1000;: -0.43681`_0.00040)
CS_L(@, 600.33718`_34.78527)
CS_G(@, 348.56574`_23.46300)
Strain_L(0, 0.11324`_0.00464)
Strain_G(@, 0.02702`_0.01518)
space_group "Cmc21"
load hkl_m_d_th2 I
{ 0
      2 0 2
                      7.55417
                                     11.70523 0.106817448
   ... not all shown }
out "results.txt" append
Out(T1000, " \n %11.1f")
Out(a_al2o3_T1000, " %11.5f", " %11.5f")
Out(c_al2o3_T1000, " %11.5f", " %11.5f")
Out(Get(r_wp), " %11.2f")
Out(Get(gof), " %11.2f")
Out(height_T1000, "%11.5f", " %11.5f")
Out(a_al5bo9_T1000, "%11.5f", " %11.5f")
Out(b_al5bo9_T1000, "%11.5f", " %11.5f")
Out(c_al5bo9_T1000, "%11.5f", " %11.5f")
Out(v_al5bo9_T1000, "%11.3f", " %11.3f")
                                                                                                             AI5BO9_T25
AI5BO9_T100
AI5BO9_T200
AI5BO9_T300
AI5BO9_T400
AI5BO9_T400
AI5BO9_T600
AI5BC9_T700
AI5BC9_T800
                                                                                                                     0.00 9
                                                                                                                    0.00 %
0.00 %
0.00 %
0.00 %
0.00 %
0.00 %
0.00 %
0.00 %
  110
  100-
   90-
   80-
                                                                                                             AI5BO9_T900
AI5BO9_T1000
                                                                                                                     0.00 %
   70-
   60-
Sqrt (Counts)
   50-
   40-
   30-
   20
   10-
    0
   -10-
   -20-
   -30
                                                          40
                                                                   45
                                                            2-Theta [°]
```

**Fig. B1.43:** Final refinement plot of the parameterized high-temperature refinement of  $Al_5BO_9$ . *hkl* ticks of  $Al_5BO_9$  from 25°C to 1000°C (top to bottom) are shown in black.

## **C** Acknowledgments

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## **D** Publications

#### **Peer Reviewed Publications**

 Charles A. Geiger, Evgeny Alekseev, Biljana Lazic, Martin Fisch, Thomas Armbruster, Ramona Langner, Michael Fechtelkord, Namjun Kin, Thomas Pettke, Werner Weppner: Crystal Chemistry and Stability of "Li<sub>7</sub>La<sub>3</sub>Zr<sub>2</sub>O<sub>12</sub>" Garnet: A Fast Lithium-Ion Conductor.

Inorganic Chemistry 50 (2011) 1089–1097.

- Martin Fisch, Thomas Armbruster, Daniel Rentsch, Eugen Libowitzky, Thomas Pettke: Crystal-chemistry of mullite-type aluminoborates Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub> and Al<sub>5</sub>BO<sub>9</sub>: A stoichiometry puzzle. Journal of Solid-State Chemistry 184 (2011) 70–80.
- Eva Wadoski, Thomas Armbruster, Biljana Lazic, Martin Fisch: Dehydration of the natural zeolite goosecreekite CaAl<sub>2</sub>Si<sub>6</sub>O<sub>16</sub>·5H<sub>2</sub>O upon stepwise heating: A single-crystal and powder X-ray study. *American Mineralogist*, accepted 2011.
- Martin Fisch, Thomas Armbruster: Thermal Expansion of Aluminoborates.
   In: Sergey Krivovichev (ed.) Minerals as Advanced Materials II, Springer-Verlag Berlin, submitted 2010.
- Nicola Döbelin, Tobias J. Brunner, Wendelin J. Stark, Martin Fisch, Egle Conforto, Marc Bohner:

Thermal treatment of flame-synthesized amorphous tricalcium phosphate nanoparticles. Journal of the American Ceramic Society 93(10) (2010) 3455–3463.

- Martin Fisch, Thomas Armbruster, Eugen Libowitzky: Microporous CsAlSi<sub>5</sub>O<sub>12</sub> at non-ambient conditions: Partial Na exchange at 100°C. Proceedings of the 3<sup>rd</sup> International Symposium on Advanced micro- and mesoporous materials 2009, accepted 2010.
- G. Diego Gatta, Nicola Rotiroti, Martin Fisch, Thomas Armbruster: Stability at high pressure, elastic behavior and pressure-induced structural evolution of "Al<sub>5</sub>BO<sub>9</sub>", a mullite-type ceramic material. *Physics and Chemistry of Minerals* 37(4) (2010) 227–236.

• Nicola Döbelin, Tobias J. Brunner, Wendelin J. Stark, Manuel Eggimann, Martin Fisch, Marc Bohner:

Phase evolution of thermally treated amorphous tricalcium phosphate nanoparticles. *Key Engineering Materials Volumes* 396-398 (2009) 595–598.

- Giuseppina Balassone, Angela Mormone, Manuela Rossi, Alessandro Bernardi, Martin Fisch, Thomas Armbruster, Anna Kathrin Malsy, Alfons Berger: Crystal chemical and structural characterization of an Mg-rich osumilite from Vesuvius volcano. European Journal of Mineralogy 20(5) (2008) 713–720.
- G. Diego Gatta, Nicola Rotiroti, Martin Fisch, Milen Kadiyski, Thomas Armbruster: Stability at high-pressure, elastic behaviour and pressure-induced structural evolution of CsAlSi<sub>5</sub>O<sub>12</sub>, a potential host for nuclear waste. *Physics and Chemistry of Minerals* 35(9) (2008) 521–533.
- Martin Fisch, Thomas Armbruster, Boris Kolesov: Temperature-dependent structural study of microporous CsAlSi<sub>5</sub>O<sub>12</sub>. Journal of Solid-State Chemistry 181 (2008) 423–431.

### **Conference Contributions (Poster Presentations)**

- Nicola Döbelin, Marc Bohner, Martin Fisch, Urs Eggenberger: A Novel XRD Sample Preparation Technique Simulating Physiological Conditions.
   23<sup>rd</sup> European Conference on Biomaterials 2010, Tampere, Finland.
- Martin Fisch, Thomas Armbruster, Daniel Rentsch, Eugen Libowitzky, Thomas Pettke: Mullite-type Al<sub>5</sub>BO<sub>9</sub> and Al<sub>18</sub>B<sub>4</sub>O<sub>33</sub>: A crystal-chemical puzzle?
   20<sup>th</sup> Meeting of the International Mineralogical Association 2010, Budapest, Hungary.
- Eva Wadoski, Thomas Armbruster, Biljana Lazic, Martin Fisch: Stepwise dehydration of Goosecreekite: A structural study.
   20<sup>th</sup> Meeting of the International Mineralogical Association 2010, Budapest, Hungary.
- Nicola Döbelin, Marc Bohner, Martin Fisch, Urs Eggenberger: XRD analysis of brushite cement samples under simulated physiological conditions. Swiss Society for Biomaterials 2010, Dübendorf, Switzerland.
- Martin Fisch, Thomas Armbruster: Microporous CsAlSi<sub>5</sub>O<sub>12</sub> at non-ambient conditions.
   3<sup>rd</sup> International Symposium on Advanced micro- and mesoporous materials 2009, Varna, Bulgaria.

## <u>Erklärung</u>

gemäss Art. 28 Abs. 2 RSL 05

Name/Vorname:	Fisch Martin					
Matrikelnummer:	01-123-827					
Studiengang:	Doktorat Philnat. (Erdwissenschaften)					
	Bachelor  Master  Dissertation					
Titel der Arbeit:	Crystal Chemistry of Boron-Bearing Mullite-Type Compounds					

Leiter der Arbeit: Prof. Dr. Thomas Armbruster

Ich erkläre hiermit, dass ich diese Arbeit selbständig verfasst und keine anderen als die angegebenen Quellen benutzt habe. Alle Stellen, die wörtlich oder sinngemäss aus Quellen entnommen wurden, habe ich als solche gekennzeichnet. Mir ist bekannt, dass andernfalls der Senat gemäss Artikel 36 Absatz 1 Buchstabe o des Gesetzes vom 5. September 1996 über die Universität zum Entzug des auf Grund dieser Arbeit verliehenen Titels berechtigt ist.

3012 Bern, 8. März 2011

Martin Fisch

# F Curriculum Vitae

## **Personal Details**

Name:	Fisch
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Citizenship:	Swiss

## Education

1987–1990:	Primary School, Bremgarten BE
1990–1994:	Secondary School, Bremgarten BE
1994–2001:	Grammar School, Freies Gymnasium Bern
05/2001:	Matura Typus C (Natural Sciences)
07/2001-10/2001:	Military service as Medic/Truck Driver (Soldier), Monte Ceneri (TI)
10/2001-07/2003:	Studies of Physics, University of Bern
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08/2005:	Bachelor in Earth Sciences, University of Bern Thesis: <i>CH-Asbest</i> Supervisor: Prof. Dr. T. Armbruster
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10/2007-04/2011:	Doctorate in Mineralogical Crystallography at the Institute of Geolog- ical Sciences, University of Bern Thesis: <i>Crystal Chemistry of Boron-Bearing Mullites</i> Supervisor: Prof. T. Armbruster